

Introduction to Noncommutative Geometry  
Chapter 9:  
Quantized Calculus on Manifolds.  
Noncommutative Residue, and Lower Dimensional  
Volumes

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## Main References

- Slides of my lectures on the noncommutative residue for my 2022 online course on NCG.
- Ponge, R.: *Noncommutative geometry and lower dimensional volumes in Riemannian geometry*. Lett. Math. Phys. **83** (2008), 19–32.

# Schwartz Kernels of $\Psi$ DOs

## Theorem (Schwartz's Kernel Theorem)

Let  $P : C_c^\infty(U) \rightarrow \mathcal{D}'(U)$  be a linear operator, where  $U \subseteq \mathbb{R}^n$  is an open set. TFAE:

- (i)  $P$  is continuous.
- (ii) There is  $k_P(x, y) \in \mathcal{D}'(U \times V)$  such that

$$\langle Pu, v \rangle = \langle k_P(x, y), u(x)v(y) \rangle \quad \forall u, v \in C_c^\infty(U).$$

## Remark

If  $P$  is continuous, then  $k_P(x, y)$  is unique and is called the **Schwartz kernel** of  $P$ .

## Remark

The above result continues to hold if  $U$  is a manifold.

## Lemma

Let  $\sigma(x, \xi) \in S^m(U \times \mathbb{R}^n)$ ,  $m \in \mathbb{R}$ , and set  $P = \sigma(x, D)$ . Then the Schwartz kernel of  $P$  is

$$k_P(x, y) = \check{\sigma}_{\xi \rightarrow y}(x, x - y).$$

Namely,

$$\langle k_P(x, y), u(x)v(y) \rangle = \langle \check{\sigma}_{\xi \rightarrow y}(x, y), u(x)v(x - y) \rangle.$$

## Remark

- If  $m < -n$ , then

$$\check{\sigma}_{\xi \rightarrow y}(x, y) = \int_{\mathbb{R}^n} e^{ix \cdot y} \sigma(x, \xi) d\xi \in C(U \times \mathbb{R}^n).$$

- In particular, on the diagonal  $x = y$  we have

$$k_P(x, x) = \check{\sigma}_{\xi \rightarrow y}(x, 0) = \int_{\mathbb{R}^n} \sigma(x, \xi) d\xi \in C^\infty(U).$$

- In general,  $\check{\sigma}_{\xi \rightarrow y}(x, y)$  makes sense as a distribution,

$$\langle \check{\sigma}_{\xi \rightarrow y}(x, y), v(y) \rangle := \langle \sigma(x, \xi), \check{v}(\xi) \rangle = \int_{\mathbb{R}^n} \sigma(x, \xi) \check{v}(\xi) d\xi.$$

## Lemma

Let  $\sigma(x, \xi) \in S^m(U \times \mathbb{R}^n)$ ,  $m \in \mathbb{R}$ . Then  $\check{\sigma}_{\xi \rightarrow y}(x, y)$  is  $C^\infty$  on  $U \times (\mathbb{R}^n \setminus 0)$ .

## Notation

$\Gamma = \{(x, x); x \in U\}$  (diagonal of  $U \times U$ ).

## Proposition

Let  $P \in \Psi^m(U)$ ,  $m \in \mathbb{C}$ , have Schwartz kernel  $k_P(x, y)$ .

- 1  $k_P(x, y)$  is  $C^\infty$  on  $(U \times U) \setminus \Gamma$ .
- 2 If  $m < -n$ , then

$$k_P(x, y) \in C(U \times U) \quad \text{and} \quad k_P(x, x) \in C^\infty(U).$$

## Setup

- $M$  = smooth manifold of dimension  $n$ .
- $\Gamma = \{(x, x); x \in M\}$ .

## Proposition

Let  $P \in \Psi^m(M)$ ,  $m \in \mathbb{R}$ , have Schwartz kernel  $k_P(x, x)$ .

- ①  $k_P(x, y)$  is  $C^\infty$  on  $(M \times M) \setminus \Gamma$ .
- ② If  $m < -n$ , then  $k_P(x, y)$  is continuous on  $M \times M$ , and  $k_P(x, x)$  is a smooth density on  $M$ .

## Remark

Let  $\phi : M' \rightarrow M$  be a diffeomorphism.

- We have

$$k_{\phi^*P}(x, y) = |\phi'(y)| k_P(\phi(x), \phi(y)), \quad x, y \in M'.$$

- In particular, on the diagonal  $y = x$  we get

$$k_{\phi^*P}(x, x) = |\phi'(x)| k_P(\phi(x), \phi(x)), \quad x \in M'.$$

This is the transformation law for densities.



## Proposition

Assume  $M$  is compact. Every operator  $P \in \Psi^m(M)$ ,  $m < -n$ , is trace-class, and

$$\mathrm{Tr}[P] = \int_M k_P(x, x).$$

# Logarithmic Singularity of $\Psi$ DO Kernels

## Setup

- $U \subseteq \mathbb{R}^n$  open set.
- $P \in \Psi^m(U)$ ,  $m \in \mathbb{Z}$ , with symbol  $\sigma(x, \xi) \sim \sum \sigma_{m-j}(x, \xi)$ .
- $k_P(x, y) =$  Schwartz kernel of  $P$ .

## Reminder

- 1  $k_P(x, y)$  is  $C^\infty$  for  $x \neq y$ .
- 2 If  $m < -n$ , then  $k_P(x, y)$  is continuous on  $U \times U$ .

## Question

If  $m \geq -n$ , then what is the behaviour of  $k_P(x, y)$  near  $x = y$ ?

# Logarithmic Singularity of $\Psi$ DO Kernels

## Proposition

Let  $P \in \Psi^m(U)$ ,  $m \in \mathbb{Z}$ ,  $m \geq -n$ . Near  $x = y$  its Schwartz kernel  $k_P(x, y)$  has a behaviour of the form,

$$k_P(x, y) = \sum_{1 \leq k \leq m+n} a_k(x, \frac{x-y}{|x-y|}) |x-y|^{-k} - c_P(x) \log |x-y| + O(1).$$

Here  $a_k(x, \theta) \in C^\infty(U \times \mathbb{S}^{n-1})$ , and  $c_P(x) \in C^\infty(U)$  is given by

$$c_P(x) = (2\pi)^{-n} \int_{\mathbb{S}^{n-1}} \sigma_{-n}(x, \xi) d\xi,$$

where  $\sigma_{-n}(x, \xi)$  is the symbol of degree  $-n$  of  $P$ .

## Remark

$a_k(x, \theta)$  only depends on the homogeneous symbol  $\sigma_{m-k}(x, \xi)$  of degree  $m - k$  of  $P$ .

# Logarithmic Singularity of $\Psi$ DO Kernels

## Setup

- $U_1 \subseteq \mathbb{R}^n$  open set.
- $\phi : U_1 \rightarrow U$  is a  $C^\infty$ -diffeomorphism.

## Fact

$$\phi^* P := (\phi^{-1})_* \in \Psi^m(U_1).$$

## Proposition

We have

$$c_{\phi^* P}(x) = |\phi'(x)| c_P(\phi(x)) \quad \forall x \in U_1.$$

## Remark

In other words the logarithmic singularities  $c_P(x)$  satisfy the transformation law of densities.

# Logarithmic Singularity of $\Psi$ DO Kernels

## Setup

$M$  is a smooth manifold of dimension  $n$ .

## Reminder

If  $P \in \Psi^m(M)$ , then,  $\kappa_*(P|_U) \in \Psi^m(V)$  for every chart  $\kappa : U \rightarrow V$ .

## Proposition

Let  $P \in \Psi^m(M)$ ,  $m \in \mathbb{Z}$ ,  $m \geq -n$ . There is a unique smooth density  $c_P(x)$  on  $M$  such that, for every chart  $\kappa : U \rightarrow V$ , we have

$$(c_P)_\kappa(x) = c_{\kappa_*(P|_U)}(x) = (2\pi)^{-n} \int_{\mathbb{S}^{n-1}} \sigma_{-n}^\kappa(x, \xi) d\xi,$$

where  $\sigma_{-n}^\kappa(x, \xi)$  is the symbol of degree  $-n$  of  $\kappa_*(P|_U)$ .

# Logarithmic Singularity of $\Psi$ DO Kernels

## Remark

If  $P$  is a differential operator on  $M$ , then  $c_P(x) = 0$  everywhere, since  $P$  has no symbols of negative degree.

## Remark

- In case  $P \in \Psi^{-n}(M)$ , i.e.,  $P$  has order  $m = -n$ , then its symbol of degree  $-n$  agrees with its principal symbol.
- It thus makes sense as a function  $\sigma_{-n}(x, \xi) \in C^\infty(T^*M \setminus 0)$ .
- It then can be shown that, if  $u \in C_c(M)$ , then

$$\int_M u(x)c_P(x) = (2\pi)^{-n} \iint_{S^*M} f(x)\sigma_{-n}(x, \xi)dx d\xi,$$

where  $S^*M = (T^*M \setminus 0)/\mathbb{R}_+^*$  is the cosphere bundle and  $dx d\xi$  is its Liouville measure.

# Noncommutative Residue

## Setup

$M$  is a **compact** manifold of dimension  $n$ .

## Definition (Noncommutative Residue)

If  $P \in \Psi^m(M)$ ,  $m \in \mathbb{Z}$ , then its **noncommutative residue** is

$$\text{Res}(P) := \int_M c_P(x).$$

## Remarks

- 1  $\text{Res}(P) = 0$  if  $P$  is a differential operator or if  $P$  has order  $< -n$ .
- 2 If  $P$  has order  $-n$ , then

$$\text{Res}(P) = (2\pi)^{-n} \iint_{S^*M} \sigma_{-n}(x, \xi) dx d\xi,$$

where  $\sigma_{-n}(x, \xi)$  is the principal symbol of  $P$ .

# Noncommutative Residue

## Notation

$$\Psi^{\mathbb{Z}}(M) := \bigcup_{m \in \mathbb{Z}} \Psi^m(M).$$

## Remark

As  $M$  is compact,  $\Psi^{\mathbb{Z}}(M)$  is an algebra:

- If  $P_1 \in \Psi^{m_1}(M)$  and  $P_2 \in \Psi^{m_2}(M)$  with  $m_1, m_2 \in \mathbb{Z}$ , then  $m_1 + m_2 \in \mathbb{Z}$ , and so

$$P_1 P_2 \in \Psi^{m_1+m_2}(M) \subseteq \Psi^{\mathbb{Z}}(M).$$

- If  $m_2 \geq m_1$ , then  $\Psi^{m_1}(M) \subseteq \Psi^{m_2}(M)$ , and so we have

$$P_1 + P_2 \in \Psi^{m_1}(M) + \Psi^{m_2}(M) \subseteq \Psi^{m_2}(M) \subseteq \Psi^{\mathbb{Z}}(M).$$

## Remark

We then can regard the noncommutative residue as a linear functional  $\text{Res} : \Psi^{\mathbb{Z}}(M) \rightarrow \mathbb{C}$ .



## Proposition

- ① If  $P_1 \in \Psi^{m_1}(M)$  and  $P_2 \in \Psi^{m_2}(M)$  are such that  $m_1 + m_2 \in \mathbb{Z}$ , then

$$\text{Res}(P_1 P_2) = \text{Res}(P_2 P_1).$$

- ② In particular, the noncommutative residue  $\text{Res} : \Psi^{\mathbb{Z}}(M) \rightarrow \mathbb{C}$  is a linear trace on the algebra  $\Psi^{\mathbb{Z}}(M)$ .

## Theorem (Wodzicki, Guillemin)

Every linear trace on  $\Psi^{\mathbb{Z}}(M)$  is a constant multiple of the noncommutative residue.

## Example

$\Delta_g$  = Laplacian on  $M$  associated with some Riemannian metric  $g$ .

- The principal symbol of  $\Delta_g$  is  $\sigma_2(x, \xi) = |\xi|_g^2$ , where  $|\xi|_g^2 = \sum g^{ij} \xi_i \xi_j$  is the Riemannian metric on  $T^*M$ .
- $\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$ , and its principal symbol is  $\sigma_{-n}(x, \xi) = |\xi|_g^{-n}$ .
- It then can be shown that

$$c_{\Delta_g^{-n/2}}(x) = (2\pi)^{-n} \int_{S_x^*M} |\xi|_g^{-n} d\xi = (2\pi)^{-n} |\mathbb{S}^{n-1}| \nu(g)(x),$$

where  $\nu(g)$  is the Riemannian density.

- Thus, if we set  $c(n) := (2\pi)^{-n} |\mathbb{B}^{n-1}| = n^{-1} (2\pi)^{-n} |\mathbb{S}^{n-1}|$ , then

$$\frac{1}{n} \operatorname{Res} (\Delta_g^{-\frac{n}{2}}) = c(n) \int_M \nu(g) = c(n) \operatorname{Vol}_g(M).$$

# Meromorphic Extension of the Trace

## Setup

- $M$  = compact manifold of dimension  $n$ .
- $\mu$  = positive smooth measure on  $M$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $\sigma_m(x, \xi) =$  principal symbol of  $P$ .

## Remark

The positive-elliptic assumption means that

- $\sigma_m(x, \xi) > 0$  for all  $(x, \xi) \in T^*M \setminus 0$ .
- $P^* = P$  and  $\text{Sp}(P) \subseteq [0, \infty)$ .

# Meromorphic Extension of the Trace

## Reminder

- Each operator  $P^{-z}$ ,  $z \in \mathbb{C}$ , is in  $\Psi^{-mz}(M)$ .
- Its principal symbol is  $\sigma_m m(x, \xi)^{-z}$ .

## Facts

Let  $A \in \Psi^a(M)$ ,  $a \in \mathbb{R}$ .

- The operator  $AP^{-z}$  is in  $\Psi^{a-mz}(M)$ .
- In particular  $AP^{-z}$  is trace-class for  $a - m\Re z > -n$ , i.e.,  $\Re z > m^{-1}(n + a)$ .
- Thus,  $\text{Tr}[AP^{-z}]$  is well defined for  $\Re z > m^{-1}(n + a)$ .

# Meromorphic Extension of the Trace

## Proposition (Wodzicki, Guillemin)

Let  $A \in \Psi^a(M)$ ,  $a \in \mathbb{R}$ , and set  $\Sigma := \{m^{-1}(n + a - j); j \in \mathbb{N}_0\}$ .

- 1 The function  $z \rightarrow \text{Tr}[AP^{-z}]$  has a meromorphic extension to  $\mathbb{C}$  with at worst simple pole singularities on  $\Sigma$ .
- 2 If  $\sigma \in \Sigma$ , then

$$m \text{Res}_{z=\sigma} \text{Tr} [AP^{-z}] = \text{Res} [AP^{-\sigma}].$$

- 3 In particular, if  $a \in \mathbb{Z}$ ,  $a \geq -n$ , then for  $\sigma = 0$  we get

$$m \text{Res}_{z=0} \text{Tr} [AP^{-z}] = \text{Res}(A).$$

## Remark

$\mathbb{N}_0 = \{0, 1, 2, 3, \dots\}$ .

# Schatten Properties of $\Psi$ DOs

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $\sigma_m(x, \xi)$  = principal symbol of  $P$ .

## Reminder

The positive-ellipticity of  $P$  ensures the following:

- 1 Its spectrum of  $P$  can be arranged as a sequence,

$$0 \leq \lambda_0(P) \leq \lambda_1(P) \leq \lambda_2(P) \leq \dots,$$

where each eigenvalue is repeated according to multiplicity.

- 2 For every  $z \in \mathbb{C}$ , the power  $P^z$  is in  $\Psi^{mz}(M)$  and has  $\sigma_m(x, \xi)^z$  as principal symbol.

## Reminder (Weyl's Law)

- By Weyl's law, as  $j \rightarrow \infty$  we have

$$\lambda_j(P) \sim \left( \frac{j}{\gamma(P)} \right)^{\frac{m}{n}},$$

where

$$\gamma(P) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} \sigma_m(x, \xi)^{-\frac{n}{m}} dx d\xi.$$

- Note that  $P^{-m/m}$  is an operator in  $\Psi^{-n}(M)$  whose principal symbol is  $\sigma_m(x, \xi)^{-n/m}$ .
- Thus,

$$\text{Res} [P^{-\frac{n}{m}}] = (2\pi)^{-n} \iint_{S^*M} \sigma_m(x, \xi)^{-\frac{n}{m}} dx d\xi.$$

- It then follows that

$$\gamma(P) = \frac{1}{n} \text{Res} [P^{-\frac{n}{m}}].$$

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O(j^{-\frac{m}{n}}).$$

That is,  $P^{-1} \in \mathcal{L}^{m^{-1}n, \infty}$ .

- In particular, if  $g$  is a Riemannian metric, then

$$(1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}^{m^{-1}n, \infty} \quad \forall m > 0.$$



## Proposition

Every  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , is in the weak Schatten class  $\mathcal{L}^{nm^{-1}, \infty}$ .

In particular, for  $m = n$  we obtain:

## Corollary

Every  $P \in \Psi^{-n}(M)$  is in the weak trace-class  $\mathcal{L}^{1, \infty}$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}^{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^m(M)$ , we see that  $P(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .
- This ensures that  $P(1 + \Delta_g)^{m/2}$  is bounded, i.e., it is contained in  $\mathcal{L}(L^2(M))$ .
- As  $\mathcal{L}^{nm^{-1}, \infty}$  is an ideal of  $\mathcal{L}(L^2(M))$ , it follows that

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}^{nm^{-1}, \infty}.$$

The proof is complete. □

## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).

- In this case the Weyl's law gives

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-1},$$

where  $\gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1})$ .

- It then follows that  $P^{-1}$  is strongly measurable, and we have

$$\int P^{-1} = \gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1}).$$

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \text{Res}(P).$$

## Remark

- 1 Connes (CMP '88) established measurability and derived the trace formula.
- 2 Lord-Potapov-Sukochev (Adv. Math. '13) obtained strong measurability.

Proof.

**Step 1:**  $P \in \Psi^{-(n+1)}(M)$ .

- In this case  $\text{Res}(P) = 0$ .
- Moreover,  $P$  is in the weak Schatten class  $\mathcal{L}^{p,\infty}$ , with  $p = n(n+1)^{-1} < 1$ .
- Thus,  $P$  is in  $\mathcal{L}^1 \subseteq \mathcal{L}_0^{1,\infty}$ .
- It follows that  $P$  is strongly measurable, and we have

$$\int P = 0 = \frac{1}{n} \text{Res}(P).$$

□

Proof.

**Step 2:**  $P = Q^{-1}$ , where  $Q \in \Psi^n(M)$  is positive-elliptic.

- In this case we saw that  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \int_{S^*M} \sigma_n(Q)(x, \xi)^{-1} dx d\xi.$$

- Here  $\sigma_n(Q)(x, \xi)^{-1} = \sigma_{-n}(Q^{-1})(x, \xi) = \sigma_{-n}(P)(x, \xi)$ .
- Thus,

$$\int P = \frac{1}{n} \int_{S^*M} \sigma_{-n}(P)(x, \xi) dx d\xi = \frac{1}{n} \text{Res}(P).$$

□

## Proof.

**Step 3:**  $\sigma_{-n}(P)(x, \xi) > 0$ .

- Let  $Q \in \Psi^n(M)$  be pos.-ellipt. w/  $\sigma_n(Q) = \sigma_{-n}(P)(x, \xi)^{-1}$ .
- Here  $Q^{-1}$  is an operator in  $\Psi^{-n}(M)$  with

$$\sigma_{-n}(Q^{-1})(x, \xi) = \sigma_n(Q)(x, \xi)^{-1} = \sigma_{-n}(P)(x, \xi).$$

- This ensures that  $R := P - Q^{-1}$  is in  $\Psi^{-(n+1)}(M)$ .
- In particular  $\text{Res}(P) = \text{Res}(Q^{-1})$ .
- By Step 2  $Q^{-1}$  is strongly meas. and  $\int Q^{-1} = \frac{1}{n} \text{Res}(Q^{-1})$ .
- As  $R \in \Psi^{-(n+1)}(M)$ , by Step 1 it is strongly measurable and  $\int R = 0$ .
- It follows that  $P = Q^{-1} + R$  is strongly measurable, and

$$\int P = \int Q^{-1} + \int R = \frac{1}{n} \text{Res}(Q^{-1}) = \frac{1}{n} \text{Res}(P).$$



Proof.

**Step 4:**  $P^* = P$ .

- Here  $\sigma_{-n}(P)(x, \xi)$  is Hermitian for all  $(x, \xi) \in T^*M \setminus 0$ .
- Bearing in mind that  $S^*M$  is compact, set

$$c := \sup \{ |\sigma_{-n}(P)(x, \xi)|; (x, \xi) \in S^*M \} < \infty.$$

- For  $(x, \xi) \in T^*M \setminus 0$ , we then have

$$\begin{aligned} \sigma_{-n}(P)(x, \xi) &\leq |\sigma_{-n}(P)(x, \xi)| \\ &\leq |\xi|_g^{-n} |\sigma_{-n}(P)(x, |\xi|_g^{-n} \xi)| \\ &\leq c |\xi|_g^{-n}. \end{aligned}$$

□



## Proof.

- Let  $P_1 \in \Psi^{-n}(M)$  have principal symbol

$$\sigma_{-n}(P_1)(x, \xi) = (c + \epsilon)|\xi|_g \operatorname{id}_{\mathcal{E}_x}, \quad \epsilon > 0.$$

- Set  $P_2 = P_1 - P$ . Then  $P_2 \in \Psi^{-n}(M)$ , and

$$\begin{aligned}\sigma_{-n}(P_2)(x, \xi) &= \sigma_{-n}(P_1)(x, \xi) - \sigma_{-n}(P)(x, \xi) \\ &\geq (c + \epsilon)|\xi|_g^{-n} - c|\xi|_g^{-n} \\ &\geq \epsilon|\xi|_g^{-n} > 0.\end{aligned}$$

- As  $\sigma_{-n}(P_j)(x, \xi)$ , Step 3 ensures that each  $P_j$  is strongly measurable and  $\int P_j = \frac{1}{n} \operatorname{Res}(P_j)$ .
- It then follows that  $P = P_1 - P_2$  is strongly measurable, and

$$\int P = \int P_1 - \int P_2 = \frac{1}{n} \operatorname{Res}(P_1) - \frac{1}{n} \operatorname{Res}(P_2) = \frac{1}{n} \operatorname{Res}(P).$$

□

Proof.

**Step 5:** General case  $P \in \Psi^{-n}(M)$ .

- Put  $P = \Re P + i\Im P$ , with

$$P_1 = \Re P = \frac{1}{2}(P + P^*), \quad P_2 = \Im P = \frac{1}{2i}(P - P^*).$$

- Here  $P_1$  and  $P_2$  are selfadjoint operators in  $\Psi^{-n}(M)$ .
- By Step 4 each operator  $P_j$  is strongly measurable and  $\int P_j = \frac{1}{n} \text{Res}(P_j)$ .
- It then follows that  $P = P_1 + iP_2$  is strongly measurable, and

$$\int P = \int P_1 + i \int P_2 = \frac{1}{n} \text{Res}(P_1) + \frac{i}{n} \text{Res}(P_2) = \frac{1}{n} \text{Res}(P).$$

□

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes sense for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.
- This includes  $\Psi$ DOs that are not infinitesimals or even bounded!

## Definition

For any  $P \in \Psi^{\mathbb{Z}}(M)$  we set

$$\int P := \frac{1}{n} \text{Res}(P) = \frac{1}{n} \int_M c_P(x).$$

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.
- $\Delta_g$  = Laplacian associated with  $g$ .

## Theorem (Connes' Integration Formula)

For every  $f \in C^\infty(M)$ , the operator  $f \Delta_g^{-n/2}$  is strongly measurable, and we have

$$\int f \Delta_g^{-\frac{n}{2}} = c(n) \int_M f(x) \nu(g)(x),$$

where we have set  $c(n) := (2\pi)^{-n} |\mathbb{B}^n|$ .

# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order  $0$ .
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable, and we have

$$\begin{aligned}\int f\Delta_g^{-\frac{n}{2}} &= \frac{1}{n} \operatorname{Res} (f\Delta_g^{-\frac{n}{2}}) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} f(x)|\xi|_g^{-n} dx d\xi \\ &= \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| \int_M f(x) \nu(g)(x).\end{aligned}$$

This proves the result. □

## Remark

- The integration formula actually holds for all  $f \in C(M)$ .
- It can be rewritten as

$$\int \Delta_g^{-\frac{n}{4}} f \Delta_g^{-\frac{n}{4}} = c(n) \int_M f(x) \nu(g)(x) \quad \forall f \in C(M).$$

## Definition (Zygmund)

$L\log L(M)$  consists of measurable functions  $f : M \rightarrow \mathbb{C}$  such that

$$\int_M |f(x)| \log(1 + |f(x)|) \nu(g)(x) < \infty.$$

## Proposition

$L\log L(M)$  is a Banach space with respect to the norm,

$$\|f\|_{L\log L} := \inf \left\{ \lambda > 0; \int_M |\lambda^{-1} f(x)| \log(1 + \lambda^{-1} |f(x)|) \nu(g)(x) < 1 \right\}.$$

## Theorem (Solomyak, Sukochev-Zanin)

- 1 If  $f \in L\log L(M)$ , then  $(1 + \Delta)^{-n/4} f (1 + \Delta)^{-n/4} \in \mathcal{L}^{1, \infty}$ .
- 2 There is  $C > 0$  such that

$$\|(1 + \Delta)^{-\frac{n}{4}} f (1 + \Delta)^{-\frac{n}{4}}\|_{1, \infty} \leq C \|f\|_{L\log L} \quad \forall f \in L\log L(M).$$



## Theorem

If  $f \in L \log L(M)$ , then  $(1 + \Delta)^{-n/4} f (1 + \Delta)^{-n/4}$  is strongly measurable, and

$$\int (1 + \Delta)^{-\frac{n}{4}} f (1 + \Delta)^{-\frac{n}{4}} = c(n) \int_M f(x) \nu(g)(x).$$

## Consequence

- Connes' integration formula shows that the NC integral recaptures the **Riemannian volume density**. Namely,

$$\int_M f(x) \nu(g)(x) = c(n)^{-1} \int f \Delta_g^{-\frac{n}{2}} \quad \forall f \in C(M).$$

- We regard  $c(n)^{-1} \Delta_g^{-n/2}$  as the **NC volume element**.
- The volume element is  $ds^n$ , where  $ds$  is the length element.
- Thus,  $ds$  is the  $n$ -th root of the volume element.

## Definition

The **NC length element** of  $(M^n, g)$  is the operator,

$$ds := \left( c(n)^{-1} \Delta_g^{-\frac{n}{2}} \right)^{\frac{1}{n}} = c(n)^{-\frac{1}{n}} \Delta_g^{-\frac{1}{2}}.$$

## Remark

$ds$  is a  $\Psi$ DO of order  $-1$ .

## Facts

- For  $k = 1, \dots, n$  the  $k$ -th dimensional volume is meant to be the integral of  $ds^k$ .
- Here  $ds^k$  is a  $\Psi$ DO of order  $-k$ .
- The NC integral has been extended to all  $\Psi$ DOs.
- This enables us to define  $k$ -dimensional volumes for all  $k = 1, \dots, n - 1$ .

## Definition

For  $k = 1, \dots, n$ , the  $k$ -th dimensional volume of  $(M^n, g)$  is

$$\text{Vol}_g^{(k)}(M) := \int ds^k = c(n)^{-\frac{k}{n}} \int \Delta_g^{-\frac{k}{2}}.$$

In particular, the length and area of  $(M^n, g)$  are

$$\text{Length}_g(M) := \int ds = c(n)^{-\frac{1}{n}} \int \Delta_g^{-\frac{1}{2}},$$

$$\text{Area}_g(M) := \int ds^2 = c(n)^{-\frac{2}{n}} \int \Delta_g^{-1}.$$

## Proposition

① If  $k$  and  $n$  have opposite parities (i.e.,  $n - k$  is odd), then  $\text{Vol}_g^{(k)}(M) = 0$ .

② If  $k = n - 2$ , then

$$\text{Vol}_g^{(n-2)}(M) = c(n, 2) \int_M \kappa_g(x) \nu(g)(x),$$

where  $\kappa_g(x)$  is the scalar curvature of  $(M, g)$ .

③ In general, we have

$$\text{Vol}_g^{(n-k)}(M) = c(n, k) \int_M I_g^{(k)}(x) \nu(g)(x),$$

where  $I_g^{(k)}(x)$  is a universal polynomial in the curvature tensor and its covariant derivatives.

## Remark

- The definition of the  $k$ -th dimensional volumes involved noncommutative geometry.
- However, the formulas in the previous slide provide purely differential-geometric expressions for the  $k$ -th dimensional volumes.

## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the **Einstein-Hilbert action**.
- It accounts for the contribution of **gravity forces** in the **Standard Model** from theoretical physics.
- We have

$$\begin{aligned}\int_M \kappa_g \nu(g) &= c(n, 2)^{-1} \int \Delta_g^{-n+2} \\ &= \frac{1}{n} c(n, 2)^{-1} \operatorname{Res} (\Delta_g^{-n+2}) \\ &= \frac{2}{n} c(n, 2)^{-1} \operatorname{Res}_{s=\frac{n}{2}-1} \operatorname{Tr} [\Delta_g^{-s}].\end{aligned}$$

- This yields a **spectral theoretic interpretation** of the Einstein-Hilbert action.
- This an important ingredient in the spectral action formalism of Connes-Chamseddine-Marcolli.