

# Noncommutative Geometry

## Chapter 8: Connes' Trace Theorem and Lower Dimensional Volumes

Sichuan University, Spring 2025

## Additional References

- Slides of my 2022 online course.

## Definition (Weak Schatten Classes $\mathcal{L}_{p,\infty}$ )

Let  $p > 0$ .

- 1 The weak Schatten class  $\mathcal{L}_{p,\infty}$  consists of all  $T \in \mathcal{L}(\mathcal{H})$  such that

$$\mu_j(T) = O\left(j^{-\frac{1}{p}}\right) \quad \text{as } j \rightarrow \infty.$$

## Definition (Weak Schatten Classes $\mathcal{L}_{p,\infty}$ )

Let  $p > 0$ .

- 1 The weak Schatten class  $\mathcal{L}_{p,\infty}$  consists of all  $T \in \mathcal{L}(\mathcal{H})$  such that

$$\mu_j(T) = O\left(j^{-\frac{1}{p}}\right) \quad \text{as } j \rightarrow \infty.$$

- 2 For  $T \in \mathcal{L}(\mathcal{H})$ , we set

$$\|T\|_{p,\infty} := \sup_{j \geq 0} (n+1)^{\frac{1}{p}} \mu_j(T).$$

## Proposition

- ①  $\mathcal{L}_{p,\infty}$  is a two-sided ideal of  $\mathcal{L}(\mathcal{H})$ .
- ②  $\|\cdot\|_{p,\infty}$  is a quasi-norm which respect to which  $\mathcal{L}_{p,\infty}$  is a quasi-Banach ideal.

## Proposition

- ①  $\mathcal{L}_{p,\infty}$  is a two-sided ideal of  $\mathcal{L}(\mathcal{H})$ .
- ②  $\|\cdot\|_{p,\infty}$  is a quasi-norm which respect to which  $\mathcal{L}_{p,\infty}$  is a quasi-Banach ideal.

## Remark

- ① If  $p > 1$ , then  $\|\cdot\|_{p,\infty}$  is equivalent to the norm,

$$\|T\|'_{p,\infty} = \sup_{N \geq 1} \left\{ N^{-1+\frac{1}{p}} \sum_{j < N} \mu_j(T) \right\}, \quad T \in \mathcal{L}_{p,\infty}.$$

- ② In this case,  $\mathcal{L}_{p,\infty}$  is a Banach ideal (w.r.t. that norm).

# Reminder: Weak Schatten Classes

## Notation

$(\mathcal{L}_{p,\infty})_0$  is the closure in  $\mathcal{L}_{p,\infty}$  of finite-rank operators.

# Reminder: Weak Schatten Classes

## Notation

$(\mathcal{L}_{p,\infty})_0$  is the closure in  $\mathcal{L}_{p,\infty}$  of finite-rank operators.

## Proposition

*We have*

$$(\mathcal{L}_{p,\infty})_0 = \left\{ T \in \mathcal{L}(\mathcal{H}); \mu_j(T) = o\left(j^{-\frac{1}{p}}\right) \right\}.$$



# Reminder: Weak Schatten Classes

## Notation

$(\mathcal{L}_{p,\infty})_0$  is the closure in  $\mathcal{L}_{p,\infty}$  of finite-rank operators.

## Proposition

*We have*

$$(\mathcal{L}_{p,\infty})_0 = \left\{ T \in \mathcal{L}(\mathcal{H}); \mu_j(T) = o\left(j^{-\frac{1}{p}}\right) \right\}.$$

## Remark

For  $0 < p < q$  we have strict inclusions,

$$\mathcal{L}_p \subsetneq (\mathcal{L}_{p,\infty})_0 \subsetneq \mathcal{L}_{p,\infty} \subsetneq \mathcal{L}_q.$$

# Reminder: Quantized Calculus

Classical	Quantum (Connes)
Complex variable	Operator on Hilbert space $\mathcal{H}$
Real variable	Selfadjoint operator on $\mathcal{H}$
Infinitesimal variable	Compact operator on $\mathcal{H}$
Infinitesimal of order $\alpha$	Compact operator s.t. $\mu_j(T) = O(j^{-\alpha})$
Integral $\int f(x)dx$	NC integral $\oint T$

# Reminder: Quantized Calculus

Classical	Quantum (Connes)
Complex variable	Operator on Hilbert space $\mathcal{H}$
Real variable	Selfadjoint operator on $\mathcal{H}$
Infinitesimal variable	Compact operator on $\mathcal{H}$
Infinitesimal of order $\alpha$	Compact operator s.t. $\mu_j(T) = O(j^{-\alpha})$
Integral $\int f(x)dx$	NC integral $\oint T$

## Remarks

- $\mu_0(T) \geq \mu_1(T) \geq \dots$  are the singular values of  $T$ .
- $T$  is infinitesimal of order  $\alpha$  iff  $T \in \mathcal{L}_{p,\infty}$ , with  $p = \alpha^{-1}$ .

## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

- 1 It is defined on infinitesimals of order 1, i.e., on the weak trace class  $\mathcal{L}_{1,\infty}$ .

## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

- ① It is defined on infinitesimals of order 1, i.e., on the weak trace class  $\mathcal{L}_{1,\infty}$ .
- ② It should take non-negative values on positive operators.

## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

- ① It is defined on infinitesimals of order 1, i.e., on the weak trace class  $\mathcal{L}_{1,\infty}$ .
- ② It should take non-negative values on positive operators.
- ③ It vanishes on infinitesimals of order  $> 1$ .

## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

- ① It is defined on infinitesimals of order 1, i.e., on the weak trace class  $\mathcal{L}_{1,\infty}$ .
- ② It should take non-negative values on positive operators.
- ③ It vanishes on infinitesimals of order  $> 1$ .
- ④ It vanishes on the commutator space,

$$\text{Com}(\mathcal{L}_{1,\infty}) := \text{Span} \{ [A, T]; A \in \mathcal{L}(\mathcal{H}), T \in \mathcal{L}_{1,\infty} \}.$$



## Ansatz for a NC Integral (Connes)

The NC integral should have the following properties:

- ① It is defined on infinitesimals of order 1, i.e., on the weak trace class  $\mathcal{L}_{1,\infty}$ .
- ② It should take non-negative values on positive operators.
- ③ It vanishes on infinitesimals of order  $> 1$ .
- ④ It vanishes on the commutator space,

$$\text{Com}(\mathcal{L}_{1,\infty}) := \text{Span} \{[A, T]; A \in \mathcal{L}(\mathcal{H}), T \in \mathcal{L}_{1,\infty}\}.$$

In other words, it should be a positive trace on  $\mathcal{L}_{1,\infty}$  vanishing on infinitesimals of order  $> 1$ .

## Definition

If  $T$  is a compact operator, then by an eigenvalue sequence  $\lambda(T) = \{\lambda_j(T)\}_{j \geq 0}$  it is meant any sequence such that:

- $\lambda_j(T)$  is an eigenvalue of  $T$  and is repeated according to (algebraic) multiplicity.
- $|\lambda_0(T)| \geq |\lambda_1(T)| \geq \dots$ .

# Reminder: NC Integral

## Definition

If  $T$  is a compact operator, then by an eigenvalue sequence  $\lambda(T) = \{\lambda_j(T)\}_{j \geq 0}$  it is meant any sequence such that:

- $\lambda_j(T)$  is an eigenvalue of  $T$  and is repeated according to (algebraic) multiplicity.
- $|\lambda_0(T)| \geq |\lambda_1(T)| \geq \dots$ .

## Facts

- ① If  $T \in \mathcal{L}_{1,\infty}$ , and  $\{\lambda_j(T)\}_{j \geq 0}$  is any eigenvalue sequence, then

$$\frac{1}{\log N} \sum_{j < N} \lambda_j(T) = O(1).$$

# Reminder: NC Integral

## Definition

If  $T$  is a compact operator, then by an eigenvalue sequence  $\lambda(T) = \{\lambda_j(T)\}_{j \geq 0}$  it is meant any sequence such that:

- $\lambda_j(T)$  is an eigenvalue of  $T$  and is repeated according to (algebraic) multiplicity.
- $|\lambda_0(T)| \geq |\lambda_1(T)| \geq \dots$ .

## Facts

- ① If  $T \in \mathcal{L}_{1,\infty}$ , and  $\{\lambda_j(T)\}_{j \geq 0}$  is any eigenvalue sequence, then

$$\frac{1}{\log N} \sum_{j < N} \lambda_j(T) = O(1).$$

- ② If  $T \in (\mathcal{L}_{1,\infty})_0$ , then we further have

$$\frac{1}{\log N} \sum_{j < N} \lambda_j(T) \longrightarrow 0.$$

# Reminder: NC Integral

## Notation

- $l_\infty = C^*$ -algebra of **bounded** complex-valued sequences.
- $c_0$  = closed ideal of sequences converging to **0**.

# Reminder: NC Integral

## Notation

- $\ell_\infty = C^*$ -algebra of **bounded** complex-valued sequences.
- $\mathfrak{c}_0$  = closed ideal of sequences converging to **0**.

## Definition

An extended limit is any positive linear map  $\lim_\omega : \ell_\infty \rightarrow \mathbb{C}$  s.t.:

- (i)  $\lim_\omega 1 = 1$ .
- (ii)  $\lim_\omega a_j = 0$  if  $(a_j) \in \mathfrak{c}_0$ .

# Reminder: NC Integral

## Notation

- $\ell_\infty = C^*$ -algebra of **bounded** complex-valued sequences.
- $\mathfrak{c}_0$  = closed ideal of sequences converging to **0**.

## Definition

An extended limit is any positive linear map  $\lim_\omega : \ell_\infty \rightarrow \mathbb{C}$  s.t.:

- (i)  $\lim_\omega 1 = 1$ .
- (ii)  $\lim_\omega a_j = 0$  if  $(a_j) \in \mathfrak{c}_0$ .

## Remark

If  $(a_j) \in \ell_\infty$ , then

$$(\lim a_j = L) \iff (\lim_\omega a_j = L \quad \forall \lim_\omega).$$

## Definition

If  $\lim_\omega$  is an extended limit, then  $\text{Tr}_\omega : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is given by

$$\text{Tr}_\omega(T) := \lim_\omega \left\{ \frac{1}{\log N} \sum_{j < N} \lambda_j(T) \right\}, \quad T \in \mathcal{L}_{1,\infty}.$$



# Reminder: NC Integral

## Definition

If  $\lim_\omega$  is an extended limit, then  $\mathrm{Tr}_\omega : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is given by

$$\mathrm{Tr}_\omega(T) := \lim_\omega \left\{ \frac{1}{\log N} \sum_{j < N} \lambda_j(T) \right\}, \quad T \in \mathcal{L}_{1,\infty}.$$

## Proposition (Dixmier)

- 1  $\mathrm{Tr}_\omega$  is a positive linear trace on  $\mathcal{L}_{1,\infty}$ .
- 2 It is annihilated by  $(\mathcal{L}_{1,\infty})_0$ , and hence it vanishes on infinitesimals of order  $> 1$ .

# Reminder: NC Integral

## Definition

If  $\lim_\omega$  is an extended limit, then  $\text{Tr}_\omega : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is given by

$$\text{Tr}_\omega(T) := \lim_\omega \left\{ \frac{1}{\log N} \sum_{j < N} \lambda_j(T) \right\}, \quad T \in \mathcal{L}_{1,\infty}.$$

## Proposition (Dixmier)

- ①  $\text{Tr}_\omega$  is a positive linear trace on  $\mathcal{L}_{1,\infty}$ .
- ② It is annihilated by  $(\mathcal{L}_{1,\infty})_0$ , and hence it vanishes on infinitesimals of order  $> 1$ .

## Definition

$\text{Tr}_\omega$  is called the Dixmier trace associated with  $\lim_\omega$ .

## Definition (Connes)

- 1 An operator  $T \in \mathcal{L}_{1,\infty}$  is called measurable if the value of  $\mathrm{Tr}_\omega(T)$  does not depend on the extended limit.

## Definition (Connes)

- 1 An operator  $T \in \mathcal{L}_{1,\infty}$  is called measurable if the value of  $\mathrm{Tr}_\omega(T)$  does not depend on the extended limit.
- 2 We denote by  $\mathcal{M}$  the class of measurable operators.

# Reminder: NC Integral

## Definition (Connes)

- 1 An operator  $T \in \mathcal{L}_{1,\infty}$  is called measurable if the value of  $\text{Tr}_\omega(T)$  does not depend on the extended limit.
- 2 We denote by  $\mathcal{M}$  the class of measurable operators.
- 3 If  $T \in \mathcal{M}$ , we define its NC integral by

$$\oint T := \text{Tr}_\omega(A),$$

where  $\text{Tr}_\omega$  is any Dixmier trace.

# Reminder: NC Integral

## Definition (Connes)

- ① An operator  $T \in \mathcal{L}_{1,\infty}$  is called measurable if the value of  $\text{Tr}_\omega(T)$  does not depend on the extended limit.
- ② We denote by  $\mathcal{M}$  the class of measurable operators.
- ③ If  $T \in \mathcal{M}$ , we define its NC integral by

$$\oint T := \text{Tr}_\omega(A),$$

where  $\text{Tr}_\omega$  is any Dixmier trace.

## Proposition (Connes, Lord-Sukochev-Zanin)

Let  $T \in \mathcal{L}_{1,\infty}$ . Then TFAE:

- ①  $T$  is measurable and  $\oint T = L$ .
- ② We have

$$\lim_{N \rightarrow \infty} \frac{1}{\log N} \sum_{j < N} \lambda_j(T) = L.$$

## Proposition

- ①  $\mathcal{M}$  is a closed subspace of  $\mathcal{L}_{1,\infty}$  that contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ .

## Proposition

- ①  $\mathcal{M}$  is a closed subspace of  $\mathcal{L}_{1,\infty}$  that contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ .
- ②  $f: \mathcal{M} \rightarrow \mathbb{C}$  is a positive linear functional that vanishes on  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ .



## Proposition

- ①  $\mathcal{M}$  is a closed subspace of  $\mathcal{L}_{1,\infty}$  that contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ .
- ②  $f: \mathcal{M} \rightarrow \mathbb{C}$  is a positive linear functional that vanishes on  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ .
- ③ In particular, this is a positive trace that annihilates infinitesimals of order  $> 1$ .

## Definition

A trace  $\varphi : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is called normalized if

$$(T \geq 0 \text{ and } \lambda_j(T) = (j+1)^{-1}) \implies \varphi(T) = 1.$$

## Definition

A trace  $\varphi : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is called normalized if

$$(T \geq 0 \text{ and } \lambda_j(T) = (j+1)^{-1}) \implies \varphi(T) = 1.$$

## Remarks

- 1 Every Dixmier trace  $\text{Tr}_\omega$  is a normalized trace.

## Definition

A trace  $\varphi : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is called normalized if

$$(T \geq 0 \text{ and } \lambda_j(T) = (j+1)^{-1}) \implies \varphi(T) = 1.$$

## Remarks

- ① Every Dixmier trace  $\text{Tr}_\omega$  is a normalized trace.
- ② There are (uncountably) many normalized positive traces on  $\mathcal{L}_{1,\infty}$  that are not Dixmier traces.

## Definition

A trace  $\varphi : \mathcal{L}_{1,\infty} \rightarrow \mathbb{C}$  is called normalized if

$$(T \geq 0 \text{ and } \lambda_j(T) = (j+1)^{-1}) \implies \varphi(T) = 1.$$

## Remarks

- ① Every Dixmier trace  $\text{Tr}_\omega$  is a normalized trace.
- ② There are (uncountably) many normalized positive traces on  $\mathcal{L}_{1,\infty}$  that are not Dixmier traces.

## Definition

An operator  $T \in \mathcal{L}_{1,\infty}$  is called strongly measurable (or positively measurable) if  $\varphi(T)$  takes the same value as  $\varphi$  ranges over all normalized positive traces.

# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

## Proposition

- 1  $\mathcal{M}_s$  is a closed subspace of  $\mathcal{L}_{1,\infty}$ .

# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

## Proposition

- 1  $\mathcal{M}_s$  is a closed subspace of  $\mathcal{L}_{1,\infty}$ .
- 2 It contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ . In particular, it contains all infinitesimals of order  $> 1$ .



# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

## Proposition

- ①  $\mathcal{M}_s$  is a closed subspace of  $\mathcal{L}_{1,\infty}$ .
- ② It contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ . In particular, it contains all infinitesimals of order  $> 1$ .
- ③ It does not depend on the inner product of  $\mathcal{L}(\mathcal{H})$ .

# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

## Proposition

- ①  $\mathcal{M}_s$  is a closed subspace of  $\mathcal{L}_{1,\infty}$ .
- ② It contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ . In particular, it contains all infinitesimals of order  $> 1$ .
- ③ It does not depend on the inner product of  $\mathcal{L}(\mathcal{H})$ .

## Proposition

Let  $T \in \mathcal{L}_{1,\infty}$ ,  $T \geq 0$ , be such that

$$\lim_{j \rightarrow \infty} j \lambda_j(T) = L.$$

# Strong Measurability

## Notation

$\mathcal{M}_s$  = class of strongly measurable operators.

## Proposition

- ①  $\mathcal{M}_s$  is a closed subspace of  $\mathcal{L}_{1,\infty}$ .
- ② It contains  $\text{Com}(\mathcal{L}_{1,\infty})$  and  $(\mathcal{L}_{1,\infty})_0$ . In particular, it contains all infinitesimals of order  $> 1$ .
- ③ It does not depend on the inner product of  $\mathcal{L}(\mathcal{H})$ .

## Proposition

Let  $T \in \mathcal{L}_{1,\infty}$ ,  $T \geq 0$ , be such that

$$\lim_{j \rightarrow \infty} j \lambda_j(T) = L.$$

Then  $T$  is strongly measurable and  $\int T = L$ .

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $p_m(x, \xi)$  = principal symbol of  $P$ .

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $p_m(x, \xi)$  = principal symbol of  $P$ .

## Remark

Positive-ellipticity means that  $p_m(x, \xi) > 0$  &  $P$  is selfadjoint with non-negative spectrum.

## Reminder

- The spectrum of  $P$  can be arranged as a sequence,

$$0 \leq \lambda_0(P) \leq \lambda_1(P) \leq \lambda_2(P) \leq \cdots,$$

where each eigenvalue is repeated according to multiplicity.



## Reminder

- The spectrum of  $P$  can be arranged as a sequence,

$$0 \leq \lambda_0(P) \leq \lambda_1(P) \leq \lambda_2(P) \leq \cdots,$$

where each eigenvalue is repeated according to multiplicity.

- By Weyl's law, as  $j \rightarrow \infty$  we have

$$\lambda_j(P) \sim \left( \frac{j}{\gamma(P)} \right)^{\frac{m}{n}},$$

where

$$\gamma(P) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} p_m(x, \xi)^{-\frac{n}{m}} dx d\xi = \frac{1}{n} \operatorname{Res} \left[ P^{-\frac{n}{m}} \right].$$

## Facts

Assume that  $\ker P = \{0\}$ .

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O\left(j^{-\frac{m}{n}}\right).$$

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O\left(j^{-\frac{m}{n}}\right).$$

That is,  $P^{-1} \in \mathcal{L}_{m^{-1}n, \infty}$ .

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O(j^{-\frac{m}{n}}).$$

That is,  $P^{-1} \in \mathcal{L}_{m^{-1}n, \infty}$ .

- In particular, if  $g$  is a Riemannian metric, then

$$(1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}_{m^{-1}n, \infty} \quad \forall m > 0.$$

## Proposition

If  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , then  $P$  is in the weak Schatten class  $\mathcal{L}_{nm^{-1}, \infty}$ .

## Proposition

If  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , then  $P$  is in the weak Schatten class  $\mathcal{L}_{nm^{-1}, \infty}$ . That is,  $P$  is an infinitesimal of order  $mn^{-1}$ .



## Proposition

If  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , then  $P$  is in the weak Schatten class  $\mathcal{L}_{nm^{-1},\infty}$ . That is,  $P$  is an infinitesimal of order  $mn^{-1}$ .

In particular, for  $m = n$  we obtain:

## Corollary

If  $P \in \Psi^{-n}(M)$ , then  $P \in \mathcal{L}_{1,\infty}$ , i.e.,  $P$  is an infinitesimal of order 1.

Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .
- This ensures that  $P(1 + \Delta_g)^{m/2}$  is bounded, i.e., it is contained in  $\mathcal{L}(L^2(M))$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .
- This ensures that  $P(1 + \Delta_g)^{m/2}$  is bounded, i.e., it is contained in  $\mathcal{L}(L^2(M))$ .
- As  $\mathcal{L}_{nm^{-1}, \infty}$  is an ideal of  $\mathcal{L}(L^2(M))$ , it follows that

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}_{nm^{-1}, \infty}.$$

The proof is complete. □

## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).



## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).

- In this case the Weyl's law gives

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-1},$$

where  $\gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1})$ .

## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).

- In this case the Weyl's law gives

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-1},$$

where  $\gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1})$ .

- It then follows that  $P^{-1}$  is strongly measurable, and we have

$$\int P^{-1} = \gamma(P) = \frac{1}{n} \operatorname{Res}(P).$$

More generally, we have:

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \operatorname{Res}(P).$$

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \text{Res}(P).$$

## Remark

- 1 Connes (CMP '88) established measurability and derived the trace formula.

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \text{Res}(P).$$

## Remark

- ❶ Connes (CMP '88) established measurability and derived the trace formula.
- ❷ Lord-Potapov-Sukochev (Adv. Math. '13) obtained strong measurability.

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes sense for all integer order  $\Psi$ DOs.



## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes sense for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes sense for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.
- This includes  $\Psi$ DOs that are not infinitesimals or even bounded!

# Connes' Trace Theorem

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes sense for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.
- This includes  $\Psi$ DOs that are not infinitesimals or even bounded!

## Definition

For any  $P \in \Psi^{\mathbb{Z}}(M)$  we set

$$\oint P := \frac{1}{n} \text{Res}(P) = \frac{1}{n} \int_M c_P(x).$$

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.
- $\Delta_g$  = Laplacian associated with  $g$ ,

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.
- $\Delta_g$  = Laplacian associated with  $g$ ,

## Theorem (Connes' Integration Formula)

For every  $f \in C^\infty(M)$ , the operator  $f \Delta_g^{-n/2}$  is strongly measurable, and we have

$$\int f \Delta_g^{-\frac{n}{2}} = c(n) \int_M f(x) \nu(g)(x),$$

where we have set  $c(n) := \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| = (2\pi)^{-n} |\mathbb{B}^n|$ .

# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .



# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.

# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.

# Connes' Integration Formula

## Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable,

# Connes' Integration Formula

## Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable, and we have

$$\begin{aligned}\int f \Delta_g^{-\frac{n}{2}} &= \frac{1}{n} \operatorname{Res} (f \Delta_g^{-\frac{n}{2}}) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} f(x) |\xi|_g^{-n} dx d\xi \\ &= \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| \int_M f(x) \nu(g)(x).\end{aligned}$$

# Connes' Integration Formula

## Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable, and we have

$$\begin{aligned}\int f\Delta_g^{-\frac{n}{2}} &= \frac{1}{n} \operatorname{Res} (f\Delta_g^{-\frac{n}{2}}) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} f(x)|\xi|_g^{-n} dx d\xi \\ &= \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| \int_M f(x) \nu(g)(x).\end{aligned}$$

This proves the result. □

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.

## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $p_m(x, \xi)$  = principal symbol of  $P$ .



## Setup

- $M$  = compact manifold with positive smooth density  $\mu$ .
- $P \in \Psi^m(M)$ ,  $m > 0$ , is positive-elliptic.
- $p_m(x, \xi)$  = principal symbol of  $P$ .

## Remark

Positive-ellipticity means that  $p_m(x, \xi) > 0$  and  $P^* = P \geq 0$ .

## Reminder

- The spectrum of  $P$  can be arranged as a sequence,

$$0 \leq \lambda_0(P) \leq \lambda_1(P) \leq \lambda_2(P) \leq \cdots,$$

where each eigenvalue is repeated according to multiplicity.

## Reminder

- The spectrum of  $P$  can be arranged as a sequence,

$$0 \leq \lambda_0(P) \leq \lambda_1(P) \leq \lambda_2(P) \leq \cdots,$$

where each eigenvalue is repeated according to multiplicity.

- By Weyl's law, as  $j \rightarrow \infty$  we have

$$\lambda_j(P) \sim \left( \frac{j}{\gamma(P)} \right)^{\frac{m}{n}},$$

where

$$\gamma(P) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} p_m(x, \xi)^{-\frac{n}{m}} dx d\xi = \frac{1}{n} \operatorname{Res} \left[ P^{-\frac{n}{m}} \right].$$

## Facts

Assume that  $\ker P = \{0\}$ .

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O\left(j^{-\frac{m}{n}}\right).$$

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O\left(j^{-\frac{m}{n}}\right).$$

That is,  $P^{-1} \in \mathcal{L}_{m^{-1}n, \infty}$ .

## Facts

Assume that  $\ker P = \{0\}$ .

- We then have

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-\frac{m}{n}}.$$

- Thus,

$$\mu_j(P^{-1}) = \lambda_j(P^{-1}) = O(j^{-\frac{m}{n}}).$$

That is,  $P^{-1} \in \mathcal{L}_{m^{-1}n, \infty}$ .

- In particular, if  $g$  is a Riemannian metric, then

$$(1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}_{m^{-1}n, \infty} \quad \forall m > 0.$$



## Proposition

If  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , then  $P$  is in the weak Schatten class  $\mathcal{L}_{nm^{-1}, \infty}$ . That is,  $P$  is an infinitesimal of order  $mn^{-1}$ .

## Proposition

If  $P \in \Psi^{-m}(M)$ ,  $m > 0$ , then  $P$  is in the weak Schatten class  $\mathcal{L}_{nm^{-1},\infty}$ . That is,  $P$  is an infinitesimal of order  $mn^{-1}$ .

In particular, for  $m = n$  we obtain:

## Corollary

If  $P \in \Psi^{-n}(M)$ , then  $P \in \mathcal{L}_{1,\infty}$ , i.e.,  $P$  is an infinitesimal of order 1.

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .
- This ensures that  $P(1 + \Delta_g)^{m/2}$  is bounded, i.e., it is contained in  $\mathcal{L}(L^2(M))$ .

## Proof of the Proposition.

Let  $m > 0$  and  $P \in \Psi^{-m}(M)$ .

- Pick a Riemannian metric  $g$  on  $M$ , and write

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}}$$

- We already know that  $(1 + \Delta_g)^{-m/2} \in \mathcal{L}_{nm^{-1}, \infty}$ .
- As  $P \in \Psi^{-m}(M)$  and  $(1 + \Delta_g)^{m/2} \in \Psi^{-m}(M)$ , we see that  $A(1 + \Delta_g)^{m/2} \in \Psi^0(M)$ .
- This ensures that  $P(1 + \Delta_g)^{m/2}$  is bounded, i.e., it is contained in  $\mathcal{L}(L^2(M))$ .
- As  $\mathcal{L}_{nm^{-1}, \infty}$  is an ideal of  $\mathcal{L}(L^2(M))$ , it follows that

$$P = P(1 + \Delta_g)^{\frac{m}{2}} \cdot (1 + \Delta_g)^{-\frac{m}{2}} \in \mathcal{L}_{nm^{-1}, \infty}.$$

The proof is complete. □

## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).

- In this case the Weyl's law gives

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-1},$$

where  $\gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1})$ .



## Facts

Let  $P \in \Psi^n(M)$  be positive-elliptic with  $\ker P = \{0\}$  (e.g.,  $P = (1 + \Delta_g)^{n/2}$ ).

- In this case the Weyl's law gives

$$\lambda_j(P^{-1}) = \lambda_j(P)^{-1} \sim \left( \frac{j}{\gamma(P)} \right)^{-1},$$

where  $\gamma(P) = \frac{1}{n} \operatorname{Res}(P^{-1})$ .

- It then follows that  $P^{-1}$  is strongly measurable, and we have

$$\int P^{-1} = \gamma(P) = \frac{1}{n} \operatorname{Res}(P).$$

More generally, we have:

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \operatorname{Res}(P).$$

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \operatorname{Res}(P).$$

## Remark

- 1 Connes (CMP '88) established measurability and derived the trace formula.

# Connes' Trace Theorem

More generally, we have:

## Theorem (Connes' Trace Theorem)

If  $P \in \Psi^{-n}(M)$ , then  $P$  is strongly measurable, and we have

$$\int P = \frac{1}{n} \operatorname{Res}(P).$$

## Remark

- ❶ Connes (CMP '88) established measurability and derived the trace formula.
- ❷ Lord-Potapov-Sukochev (Adv. Math. '13) obtained strong measurability.

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes for all integer order  $\Psi$ DOs.

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.



## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.
- This includes  $\Psi$ DOs that are not infinitesimals or are not even bounded.

## Consequence

- The NC integral makes sense for  $\Psi$ DOs of order  $\leq -n$
- The NC residue, however, makes for all integer order  $\Psi$ DOs.
- Therefore, the equality between the NC integral and the NC residue enables us to extend the NC integral to all  $\Psi$ DOs.
- This includes  $\Psi$ DOs that are not infinitesimals or are not even bounded.

## Definition

For any  $P \in \Psi^{\mathbb{Z}}(M)$  we set

$$\oint P := \frac{1}{n} \text{Res}(P) = \frac{1}{n} \int_M c_P(x).$$

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.
- $\Delta_g$  = Laplacian associated with  $g$ ,

# Connes' Integration Formula

## Setup

- $(M^n, g)$  = compact Riemannian manifold.
- $\nu(g)$  = Riemannian density.
- $\Delta_g$  = Laplacian associated with  $g$ ,

## Theorem (Connes' Integration Formula)

For every  $f \in C^\infty(M)$ , the operator  $f \Delta_g^{-n/2}$  is strongly measurable, and we have

$$\int f \Delta_g^{-\frac{n}{2}} = c(n) \int_M f(x) \nu(g)(x),$$

where we have set  $c(n) := \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| = (2\pi)^{-n} |\mathbb{B}^n|$ .

# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .

# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.



# Connes' Integration Formula

Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order  $0$ .
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.

# Connes' Integration Formula

## Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable, and we have

$$\begin{aligned}\int f \Delta_g^{-\frac{n}{2}} &= \frac{1}{n} \operatorname{Res} (f \Delta_g^{-\frac{n}{2}}) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} f(x) |\xi|_g^{-n} dx d\xi \\ &= \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| \int_M f(x) \nu(g)(x).\end{aligned}$$

# Connes' Integration Formula

## Proof.

- $\Delta_g^{-n/2}$  is  $\Psi$ DOs of order  $-n$  whose principal symbol is  $|\xi|_g^{-n}$ .
- As  $f \in C^\infty(M)$ , this is a  $\Psi$ DO of order 0.
- Thus  $f\Delta_g^{-n/2}$  is in  $\Psi^{-n}(M)$  and it has  $f(x)|\xi|_g^{-n}$  as principal symbol.
- Thus, by Connes' trace theorem  $f\Delta_g^{-n/2}$  is strongly measurable, and we have

$$\begin{aligned}\int f\Delta_g^{-\frac{n}{2}} &= \frac{1}{n} \operatorname{Res} (f\Delta_g^{-\frac{n}{2}}) = \frac{(2\pi)^{-n}}{n} \iint_{S^*M} f(x)|\xi|_g^{-n} dx d\xi \\ &= \frac{1}{n} (2\pi)^{-n} |\mathbb{S}^{n-1}| \int_M f(x) \nu(g)(x).\end{aligned}$$

This proves the result. □

## Consequence

- Connes' integration formula shows that the NC integral recaptures the Riemannian volume density. Namely,

$$\int_M f \nu(g) = c(n)^{-1} \int f \Delta_g^{-\frac{n}{2}} \quad \forall f \in C^\infty(M).$$

## Consequence

- Connes' integration formula shows that the NC integral recaptures the Riemannian volume density. Namely,

$$\int_M f \nu(g) = c(n)^{-1} \oint f \Delta_g^{-\frac{n}{2}} \quad \forall f \in C^\infty(M).$$

- We regard  $c(n)^{-1} \Delta_g^{-n/2}$  as the NC volume element.

## Consequence

- Connes' integration formula shows that the NC integral recaptures the Riemannian volume density. Namely,

$$\int_M f \nu(g) = c(n)^{-1} \oint f \Delta_g^{-\frac{n}{2}} \quad \forall f \in C^\infty(M).$$

- We regard  $c(n)^{-1} \Delta_g^{-n/2}$  as the NC volume element.
- The volume element is  $ds^n$ , where  $ds$  is the length element.

## Consequence

- Connes' integration formula shows that the NC integral recaptures the Riemannian volume density. Namely,

$$\int_M f \nu(g) = c(n)^{-1} \int f \Delta_g^{-\frac{n}{2}} \quad \forall f \in C^\infty(M).$$

- We regard  $c(n)^{-1} \Delta_g^{-n/2}$  as the NC volume element.
- The volume element is  $ds^n$ , where  $ds$  is the length element.
- Thus,  $ds$  is the  $n$ -th root of the volume element.

## Definition

The NC length element of  $(M^n, g)$  is the operator,

$$ds := \left( c(n)^{-1} \Delta_g^{-\frac{n}{2}} \right)^{\frac{1}{n}} = c(n)^{-\frac{1}{n}} \Delta_g^{-\frac{1}{2}}.$$



## Definition

The NC length element of  $(M^n, g)$  is the operator,

$$ds := \left( c(n)^{-1} \Delta_g^{-\frac{n}{2}} \right)^{\frac{1}{n}} = c(n)^{-\frac{1}{n}} \Delta_g^{-\frac{1}{2}}.$$

## Remark

$ds$  is a  $\Psi$ DO of order  $-1$ .

## Facts

- For  $k = 1, \dots, n$  the  $k$ -th dimensional volume is meant to be the integral of  $ds^k$ .

## Facts

- For  $k = 1, \dots, n$  the  $k$ -th dimensional volume is meant to be the integral of  $ds^k$ .
- Here  $ds^k$  is a  $\Psi$ DO of order  $-k$ .

## Facts

- For  $k = 1, \dots, n$  the  $k$ -th dimensional volume is meant to be the integral of  $ds^k$ .
- Here  $ds^k$  is a  $\Psi$ DO of order  $-k$ .
- The NC integral has been extended to all  $\Psi$ DOs.

## Facts

- For  $k = 1, \dots, n$  the  $k$ -th dimensional volume is meant to be the integral of  $ds^k$ .
- Here  $ds^k$  is a  $\Psi$ DO of order  $-k$ .
- The NC integral has been extended to all  $\Psi$ DOs.
- This enables us to define  $k$ -dimensional volumes for all  $k = 1, \dots, n - 1$ .

## Definition

For  $k = 1, \dots, n$ , the  $k$ -th dimensional volume of  $(M^n, g)$  is

$$\text{Vol}_g^{(k)}(M) := \int ds^k = c(n)^{-\frac{k}{n}} \int \Delta_g^{-\frac{k}{2}}.$$

## Definition

For  $k = 1, \dots, n$ , the  $k$ -th dimensional volume of  $(M^n, g)$  is

$$\text{Vol}_g^{(k)}(M) := \int ds^k = c(n)^{-\frac{k}{n}} \int \Delta_g^{-\frac{k}{2}}.$$

In particular, the length and area of  $(M^n, g)$  are

$$\text{Length}_g(M) := \int ds = c(n)^{-\frac{1}{n}} \int \Delta_g^{-\frac{1}{2}},$$

$$\text{Area}_g(M) := \int ds^2 = c(n)^{-\frac{2}{n}} \int \Delta_g^{-1}.$$

## Proposition

- ① If  $k$  and  $n$  have opposite parities (i.e.,  $n - k$  is odd), then  $\text{Vol}_g^{(k)}(M) = 0$ .



## Proposition

- 1 If  $k$  and  $n$  have opposite parities (i.e.,  $n - k$  is odd), then  $\text{Vol}_g^{(k)}(M) = 0$ .
- 2 If  $k = n - 2$ , then

$$\text{Vol}_g^{(n-2)}(M) = c(n, 2) \int_M \kappa_g(x) \nu(g)(x),$$

where  $\kappa_g(x)$  is the scalar curvature of  $(M, g)$ .

## Proposition

① If  $k$  and  $n$  have opposite parities (i.e.,  $n - k$  is odd), then  $\text{Vol}_g^{(k)}(M) = 0$ .

② If  $k = n - 2$ , then

$$\text{Vol}_g^{(n-2)}(M) = c(n, 2) \int_M \kappa_g(x) \nu(g)(x),$$

where  $\kappa_g(x)$  is the scalar curvature of  $(M, g)$ .

③ In general, we have

$$\text{Vol}_g^{(n-k)}(M) = c(n, k) \int_M I_g^{(k)}(x) \nu(g)(x),$$

where  $I_g^{(k)}(x)$  is a universal polynomial in the curvature tensor and its covariant derivatives.

## Remark

- The definition of the  $k$ -th dimensional volumes involved noncommutative geometry.

## Remark

- The definition of the  $k$ -th dimensional volumes involved noncommutative geometry.
- However, the formulas in the previous slide provide purely differential-geometric expressions for the  $k$ -th dimensional volumes.

## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the Einstein-Hilbert action.

## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the Einstein-Hilbert action.
- It accounts for the contribution of gravity forces in the Standard Model from Theoretical Physics.

## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the Einstein-Hilbert action.
- It accounts for the contribution of gravity forces in the Standard Model from Theoretical Physics.
- We have

$$\begin{aligned} \int_M \kappa_g \nu(g) &= c(n, 2)^{-1} \int \Delta_g^{-n+2} \\ &= \frac{1}{n} c(n, 2)^{-1} \operatorname{Res} (\Delta_g^{-n+2}) \\ &= \frac{2}{n} c(n, 2)^{-1} \operatorname{Res}_{s=\frac{n}{2}-1} \operatorname{Tr} [\Delta_g^{-s}]. \end{aligned}$$

## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the Einstein-Hilbert action.
- It accounts for the contribution of gravity forces in the Standard Model from Theoretical Physics.
- We have

$$\begin{aligned} \int_M \kappa_g \nu(g) &= c(n, 2)^{-1} \int \Delta_g^{-n+2} \\ &= \frac{1}{n} c(n, 2)^{-1} \operatorname{Res} (\Delta_g^{-n+2}) \\ &= \frac{2}{n} c(n, 2)^{-1} \operatorname{Res}_{s=\frac{n}{2}-1} \operatorname{Tr} [\Delta_g^{-s}]. \end{aligned}$$

- This yields a spectral theoretic interpretation of the Einstein-Hilbert action.



## Remark

- The functional  $g \rightarrow \int_M \kappa_g \nu(g)$  is called the Einstein-Hilbert action.
- It accounts for the contribution of gravity forces in the Standard Model from Theoretical Physics.
- We have

$$\begin{aligned} \int_M \kappa_g \nu(g) &= c(n, 2)^{-1} \int \Delta_g^{-n+2} \\ &= \frac{1}{n} c(n, 2)^{-1} \operatorname{Res} (\Delta_g^{-n+2}) \\ &= \frac{2}{n} c(n, 2)^{-1} \operatorname{Res}_{s=\frac{n}{2}-1} \operatorname{Tr} [\Delta_g^{-s}]. \end{aligned}$$

- This yields a spectral theoretic interpretation of the Einstein-Hilbert action.
- This an important ingredient in the spectral action formalism of Connes-Chamseddine.