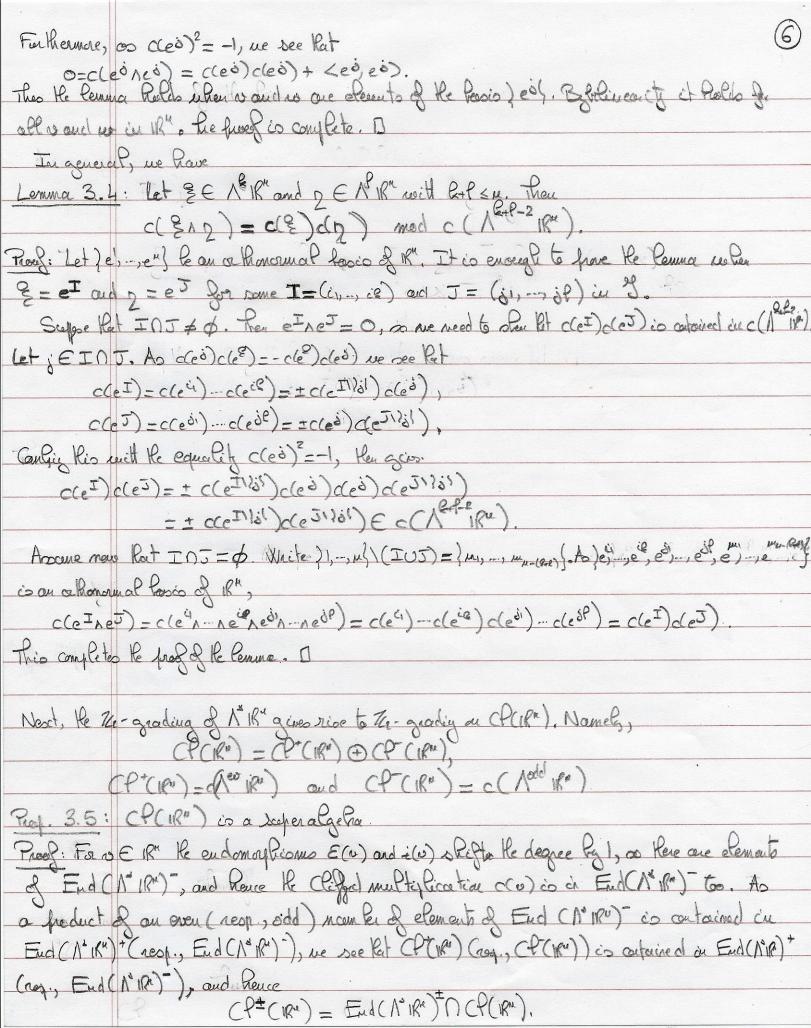


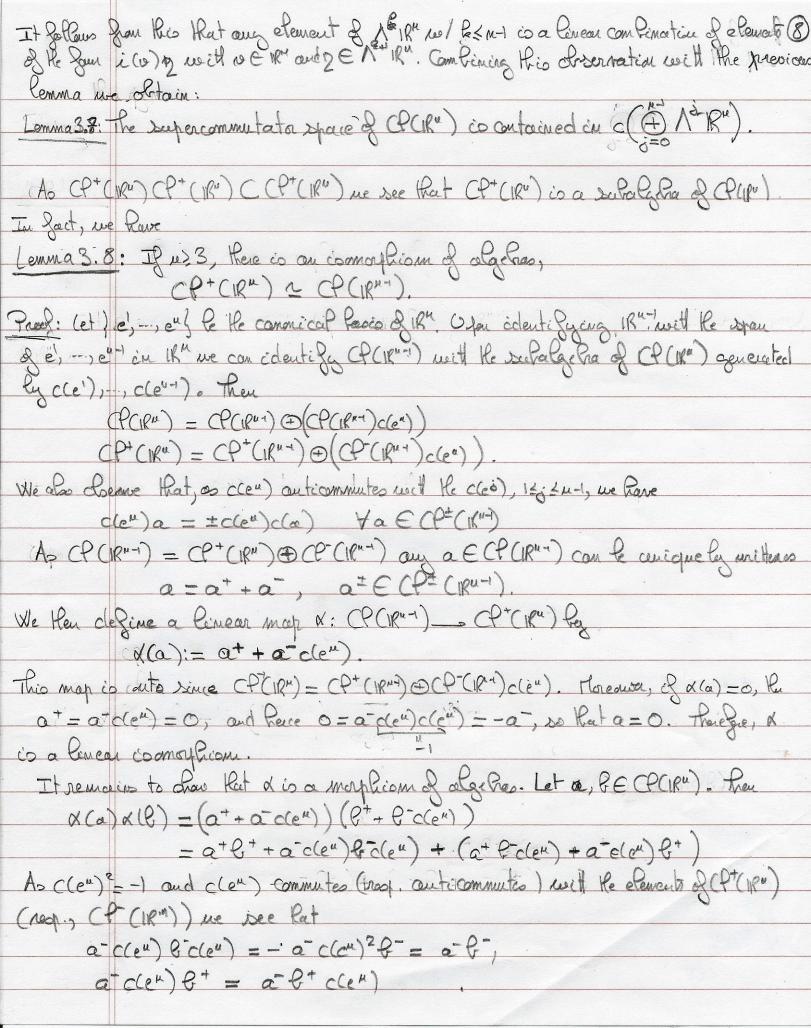
There gre, by lilamenity we obtain: Lemma 22: For all no, we  $\mathbb{R}^{n}$ ,  $\mathcal{E}(\omega) \mathcal{E}(\omega) + \mathcal{E}(\omega) \mathcal{E}(\omega) = i(\omega) i(\omega) + i(\omega) i(\omega) = 0$ ,  $\mathcal{E}(\omega)i(\omega) + i(\omega)\mathcal{E}(\omega) = \langle \omega, \omega \rangle.$ Observe also that the formulas of Lemma 2.1 show that, with respect to the or Romonnal fasio / e Titog, le matrices of le oudemaghions E(e) and i(e) are transper of each other. This,  $\mathcal{E}(e\delta)^* = i(e\delta).$ By linearity this gives: Lemma 23: For all ve 18", E(v) = i(v) and  $i(v)^* = E(v)$ . 3. The Clifford algebra CPCIRM). Def.: The Clifford algebra of 18th, denoted (PCIR"), is the subalgebra of End of (N'IR") generated the operators, -c(0):=E(0)-i(0), vEIR". Rule: The aperator (w) is called the Clifford multiplication by w. Oping Lemma 2.2 and Lemma 2.3 we get Lemma 3.1:(2) For allo, us E 1Ru, c(v)c(w) +c(w)c(v) = -2 (v) w). This implies that, gitt, (ii) For off ve IR",  $c(v)^* = -c(v)$ . (ce) c(e) = - de) de), It follows from (i) that for any orthonormal hosis je", -s e" we have and  $c(e\dot{\delta})(e^{\xi}) + c(e^{\xi})(e^{\dot{\delta}}) = -28\dot{\delta}^{\xi}$   $c(e\dot{\delta})^2 = -1$ Threamer, (ii) implies Rat CPCRa) is a Bed conder taking adjuicts, i.e., CPCRa) is an involutive sabalgelia of Endre (1 18°) (and so Riscio a C'-algeria over 1R). Let /e', e" le ou a thonormal lossis of IR". As le I Tex is an orthonormal lessis of NR" the linear map 18" 9 0 - sc(v) E CPCIR") is an iquely extanded into the linear maps c: 1'18" \_\_ PCIR")

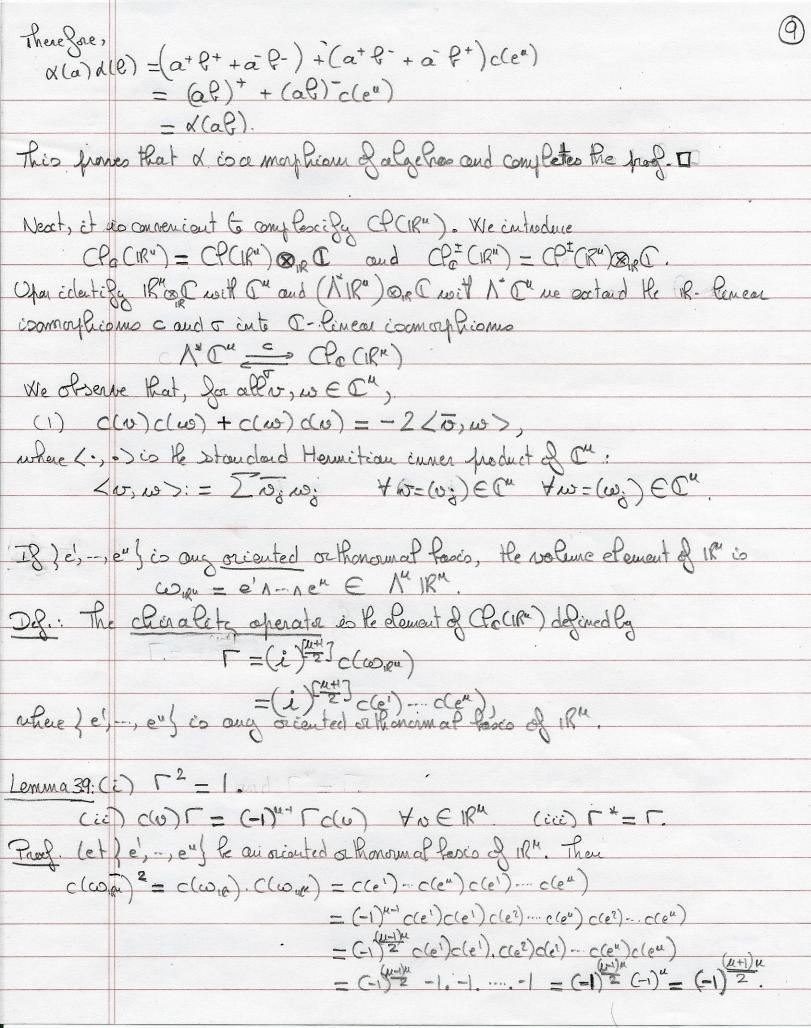
such that  $C(e^{I}) = \begin{cases} c(e^{i}) & c(e^{i}) \end{cases} \quad i \begin{cases} I = (c_{1}, \dots, c_{n}), \\ i \end{cases}$ P On the other, as CP (IRM) is an algebra of endamorphisms of 1° IRM use have a natural linear map o: CPCIRU) -> NIRU defined by T(a) = a.1 / Ha E CP(1Ra). Proop. 3.2: The linear map could to are inverses of each other. Proof: It is clear that c is onto fecause 1, c(e'), ..., c(e") generate (P(1R") Furthermore, let IES. If I= ob, the c(e) = | gud Ruce c(e), 1 = 1. Supe let I=(i), (ie) E 31 p. As c(ei) = E(ei) - i(ei) and i(ei) Cowers le degree, cle] = cleu) - c(ele) = (E(ei)-i(ei)) - (E(eie) - i(eie)) = E(e4) --- E(e42) + T, where T is an operator which Gwess le degree. In peticolo, T.1 = 0, and so we has C(e]) 1 = E(ei) ... E(eie) 1 = ei 1 ... A eie = e I This Shows that voc(eI) = eI Y IES, and Rence This pass lat c is injective and z is a left-inverse forc. As me already beause that a company, it follows that a linear isomorphism with inverse z. The proof is outlier I Ruse: An immediate consequence of the fact that to is the inverse of c, is the fact that the map c does not depend on the choic of the orthonorum at boars of eig. In fact, for any or thogonal boois } go fine have c(g'in-ng'e) = c(g'i) - c(g'e), \(\forall (i, -, ce) \in 3/\phi\).

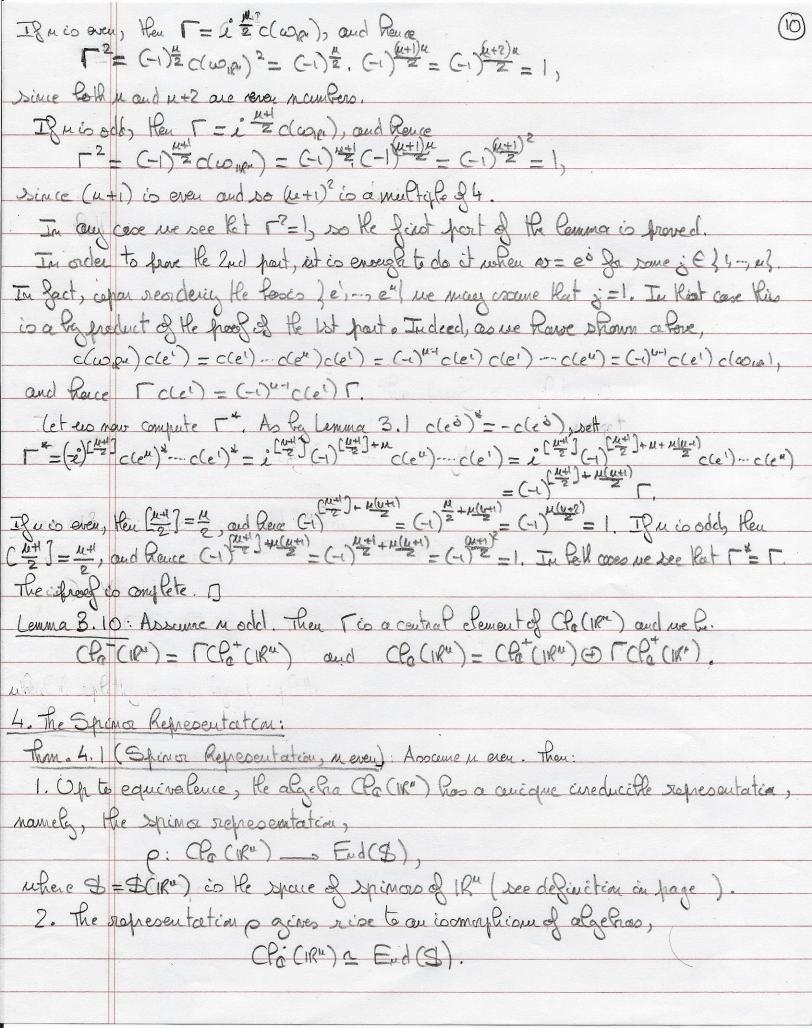
It should be stressed that could trave not isomorphisms of algebras, because we have Lemma B. 3: For all w, wo EIR", ((onw) = ((o) ((w) + (o, w). Proof: (et je',..., e" ) le an orthonormal fosis of IR". Ig j 7 R, Ken  $c(e\delta_1 c^2) = c(e\delta)c(e^2) = c(e\delta)c(e^2) + \langle e\delta, e^2 \rangle$ .



Thus, CP+CR") CP+CR") C (End(N'R")+, End(N'R")+) nCPCR") C CP+(R"), CP-CR") C (End(N'R")+, End(N'R")+) nCPCR") C CP+CR") This fraves that CPCIR") is a superalgebra. I As CPCIR") is a superalge from me can bok at superammentators. The following result, although seemingly simple, will has many important anaequences in the sequel. Cemma 36: Let & EIRM and DE 18th. Then  $[c(\sigma), c(\rho)] = 2c(i(\sigma))$ . Proof: (et) e', -, e" | le au athonomial foois of IR". Then it is enough to frame le Comma when 2 = e and = = e game I=(i, -, ie) E'S and i E ?1, -, uf. Fortermore, apar readery le losis jedjue may osocume Pat j=1 and cile I=(1,.., &) of jEI a I=(2,.., R+1) o Penocia. Super Pat j & I. Then electo = 0. Movemen, as clet )cleq) = -cleq)clet) for \$1 \$9, we have  $C(e^{\frac{1}{2}})c(e^{\frac{1}{2}}) = c(e^{\frac{1}{2}})c(e^{\frac{1}{2}})\cdots c(e^{\frac{1}{2}})\cdots c(e^{\frac{1}{2}})\cdots c(e^{\frac{1}{2}})c(e^{$  $(c(e^{\frac{1}{2}}), c(e^{\frac{1}{2}})^{\frac{1}{2}} = c(e^{\frac{1}{2}}) c(e^{\frac{1}{2}}) - (-1)^{\frac{1}{2}} c(e^{\frac{1}{2}}) c(e^{\frac{1}{2}}) = 0 = -2c(e(e^{\frac{1}{2}})e^{\frac{1}{2}}).$ Assume you lat jEI. There (e)eI = i(e)(e'n-nel) = e'a-nel. Faithermore, cos cle )2 = -1, me Rave - - 1  $c(e^{2})c(e^{2}) = c(e^{2})c(e^{2})c(e^{2}) - c(e^{2}) = -c(e^{2}) - c(e^{2}),$ (e) c(e) = c(e) -c(e) c(e) = (-1) -c(e) 2 c(e) -c(e) = (-1) -c(e). [c(e), c(e])]'=c(e)c(e])-(-)ec(e])c(e) = - c(e2) -- c(e2) - (-1)2 (-1)2 c(e2) -- c(e2) = -2c(e2)--c(e2)  $=-2c(\dot{\epsilon}(\dot{\epsilon}\dot{\epsilon})e^{\pm}).$ This completes the proof. D. Obour let if wind, , not one of Roganal muit mertas, Hen 10/1-10 = i(v)(v/10/1/1-100)







	(n)
Rmk: The ineducte	anicity up to equiner force means that if II: Clo(IK") - End (E) is ano the supresentation then & and E are isomorphic Clo(IK")-modules, i.e., there invertisomorphism (P: D - E 2.t.  CP (P(a) 5) = II(a) (P15) Y 5 E \$ Y a E CPa (IK").
eacon at	CP(P(a) 5) = Tica) P(5) Y 5 E \$ 7 a E CPa (180).
Rm. R.: Ao	we shall & has dimension 24/2, so we get the fallowing isomorphism of
0	CR (IR") 2 End (\$) 2 H2 (IR).  (IR) is a simple algebre, it is clear fact (Pa(IR") Rosa unique inveducible ion.
Since Mar	(1K) is a simple algebre, it is dear Bat (Po(1KM) Rosa unique ineducible
sepresentat	ion.
0 1	
The meat 10	plages are devoted to proving Theorem 4.1. PLEASE TORN OVER.

Proof of The	orem 4.1:
Aconstruction	of Respinor representation.
Set $w = \frac{10}{2}$	and let )c;, end le an or thonorma Person of 18".
Tree } 30, 30	and let )c;, out le aux or thonormal Peoco of 18.  30 = 1 (20 + ie 20 +1) and 35 = 30 = 1 (20 - ie 20 +1).  15 = 10 and leave mal fessio & the and we get the orthogonal splitting,  The = 1.00 + 1.00 + 1.00 for the splitting.
11111010	
This xelling	Por sine rice to the or Pengnal de confectione  N'C" = A AM Cha  fig=0
Nece	18,9 Cm - (V1,0 Cm /5 (V0,1 Cm )9
	= Span } 34 1 134 18 1 1 50 } 1 50 1 50 1 50 1 50 1 50 1 50
For Reamone	, we see that the algebra (Pa(IR") is governated by the digs bound the digs ) b. Hereaver,
ty coming	le Samula (1) n page 9 ne see that (30) c(30) + c(30) + c(30) = 0,
1)	$c(g^{\delta})dg^{\Xi}) + dg^{\Xi})dg^{\delta}) = -28^{\delta^{\Xi}},  1 \leq j, \ell \leq m \text{ (See Remon & page 5)}$ or famille is defined to $f^{\Xi}$ $f_{\mu} := \int_{0}^{0} f^{\alpha} dx = \int_{0}^{0} \int_{0}^{0} f^{\alpha} dx$
We define o	Du:= No Cu = D No Cu Pinear map o from c(N'Cu) = GN'Cu) & (No Cu) & Jallers. IS 3E No Cu, lu
£	$\rho(d\bar{z})\sigma = \sqrt{2}E(\bar{z});$ ) is the exterior product by $\bar{z}$ , i.e.,
If ze N'	E(3) c: = 3 r Y c E N'C'.
where i(z)	e(c(8)) = - 12 c(8), ) to the interior freduct for 3, that is,
į	(3) (win nwi) = = Gi) + (3, w> win nwin nwi Ywo E Note.

```
By wing Lemma 2.2 ct is not difficult to check that, for in=1,-, m, i(gò) (gè) + i(gè) i(gò) = E(gō) E(gē) + E(gē) E(gō) = 0
  and Rence i(80) E(85) + E(85) i(80) = 800,
        e(d8))(p(d8))+p(d8))p(d8))=0
         e(c(82))e(c(82))+c(c(82))e(c(82))=0

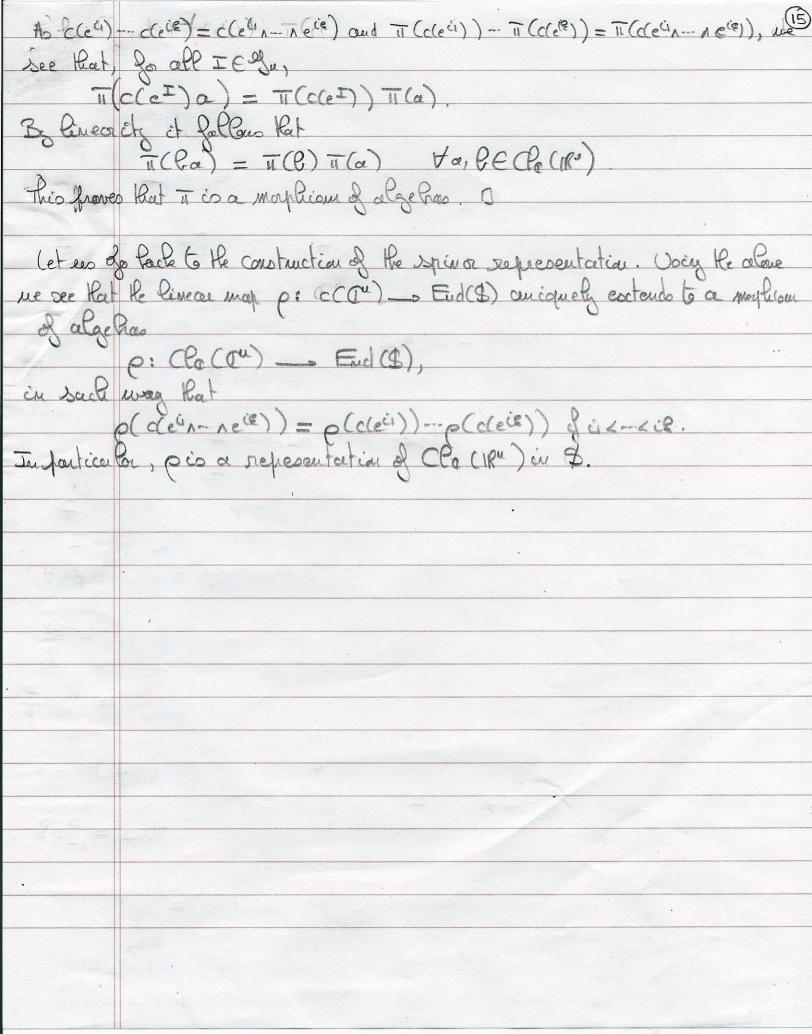
e(c(82))e(c(82))+c(c(82))e(c(82))=-286
 By felimeonity it Hen gottans that, for all vous ECu,
      p(c(v))p(c(w)) + p(c(v))p(c(v)) = -2/5, 10>
Lomma 42:let A le a (conital) algebra over C and let II: c(Cu) - A le a C-lincon maps.
 Then it canciquely extends to as mary him of algeries,
           c: CR (IR") -> A.
 Inaddition, for any I= (4, -, is) E TH
              TI (c(e'1) - TI (c(e'1)) - TI (c(e'1)).
Proof: The writing gottom for the fact that the cas) o, w E C", generate Cha (R"). The fire, me only have to show that the map II: c(C") - A can be extanded into a may him.
  of atgelias II: Cho (IR") - A.
    We excloud II outo let-lancar Mape II: Cla(IRM) o A s.t., for all I = (in), ie) EIn,
             T(C(cin-neic)) = T(C(ei)) - T(C(eic)) + I - (cin-ce) E 31 $.
  As cleir n. neie) = cleir) ... cleie), ne see but, by définition,
            Tr(c(ei) ... c(eie)) = Tr(c(ei)) ... Tr(c(eie)) 128 1644 ... 61861
  Oberne also Rat (401) implies that
          \pi(c(e^2))\pi(c(e^2))+\pi(c(e^2))\pi(c(e^3))=-286^2,
  and hence
     (4.2) TT(c(e)) TT (c(e)) = - TT (c(e)) TT (c(e)) if if e
    (4.3) = -1
```

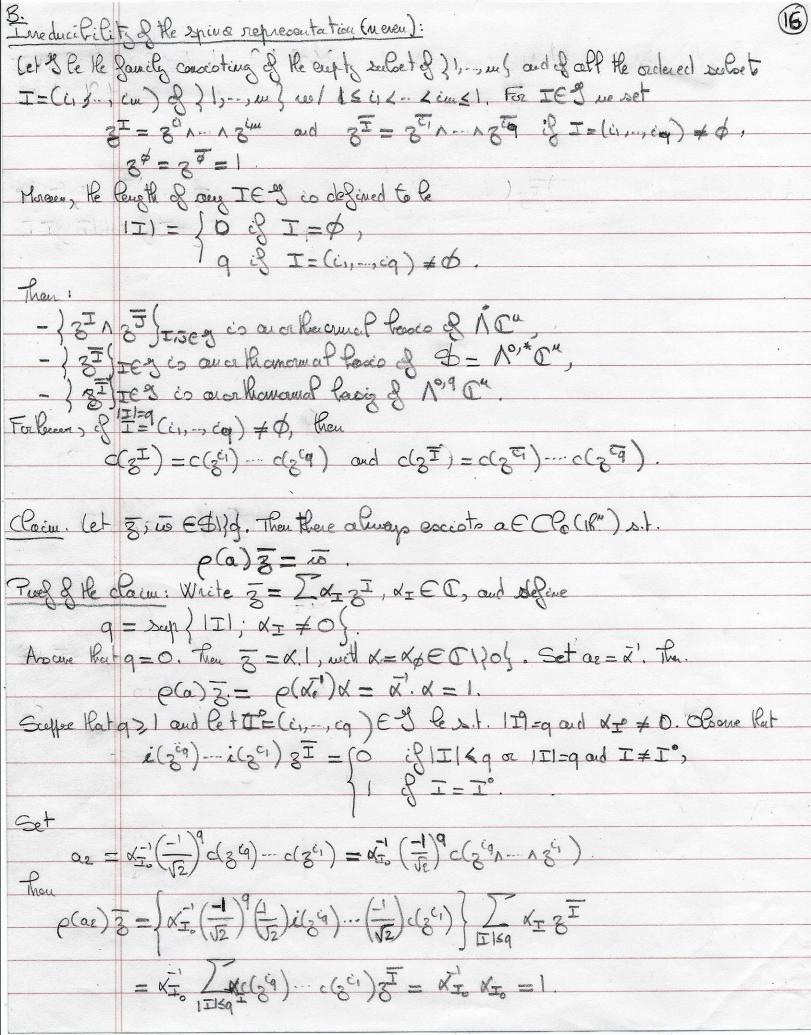
```
Claim: Let v∈ C" and a∈ CPa (IRu). Then
               TI (c(u)a) = TI (c(u)) Ti (a)
Proof of the Caim: It is enough to promo the claim when we e' and a = c(e I) for
 some; E ? 1, --, uf and I = (i, --, ie) E Ju.
 Apoceme Rat & I I I & Las Pen
          T(c(e) c(e)) = T(c(e) c(e)) - c(e))
                                = T(c(e3)) T(c(e4)) -- T(c(c2))
                                 = TI (c(e)) TI (c(e))
 Scoppe Reat i> is and let p = sep ) q; iq < is. Then
      Tocles) cle I) = cles) clein) -- cleie)
                               = (-1) cleir) - cleir) cleir) (lei) moleir)
     TI (cleb) cle I) = CI) TI (cleu) ) - TI (cleu) ) TI (cleu) ) TI (cleu)
 Therene, crises (4.2) we get
   Ti (cle) cle I) = Ti (cles).) Ti (cei) - Ti (cles) = Ti (cles) Ti (cles)

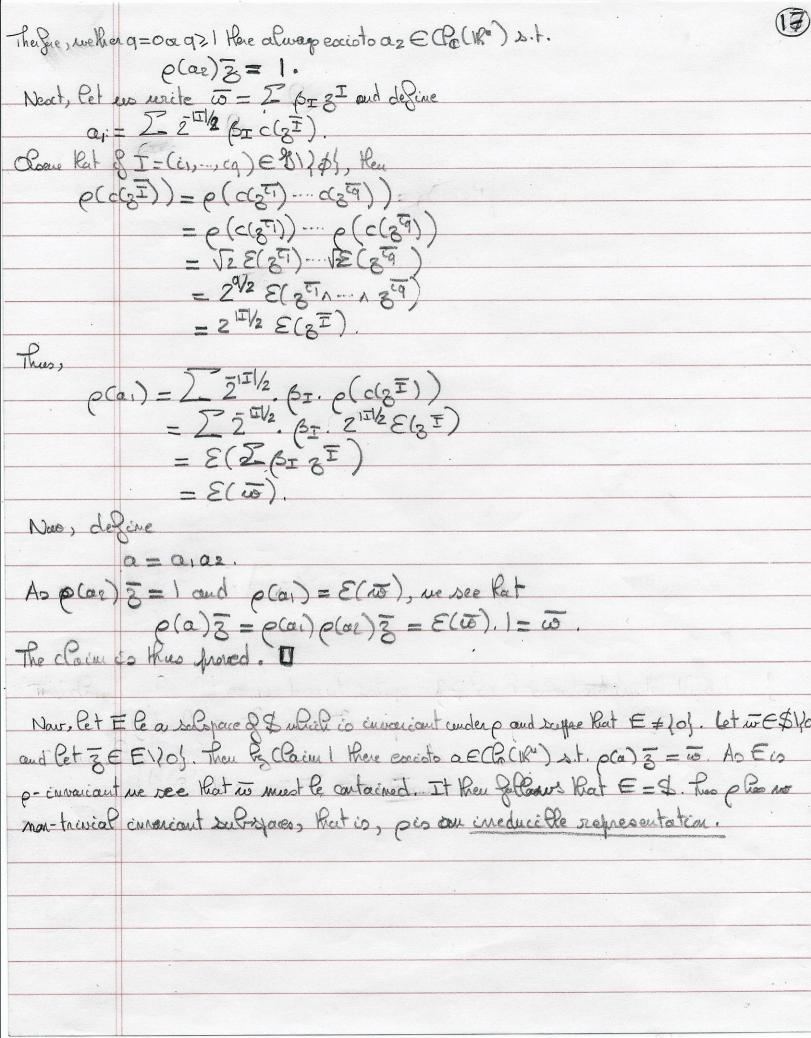
Aroceme man that i E I and let pe the index s.t. in=i. Then
        c(e) c(e) = c(ei) c(ei) ~ c(eil)
                           = (-) pt cleir) -- cleir) deix) -- cleir)
                            = C-DR o(ei) - o(eil) - o(eil)
Cikewia, by coing (4.2)-(4.8) we get
       The (clear) The (clear) The (clear) The (clear)
                         = G-) TT (c(ei)) -- TT (c(ei)) +00 TT (c(eis)).
 Ao Ti(c(e')...c(e'f)...c(e'f)) = Ti(c(e'f))...Ti(c(e'f))...Ti(c(e'f)), me deduce Pet
               \pi(c(e^{\delta})c(e^{I})) = \pi(c(e^{\delta}))\pi(e(e^{I}))
 This completes the front of the claim. I
  (et es gans faite to the proof of Lemma 4.2. (eta E (Pale"). The combining the above claim with an easy induction shows that

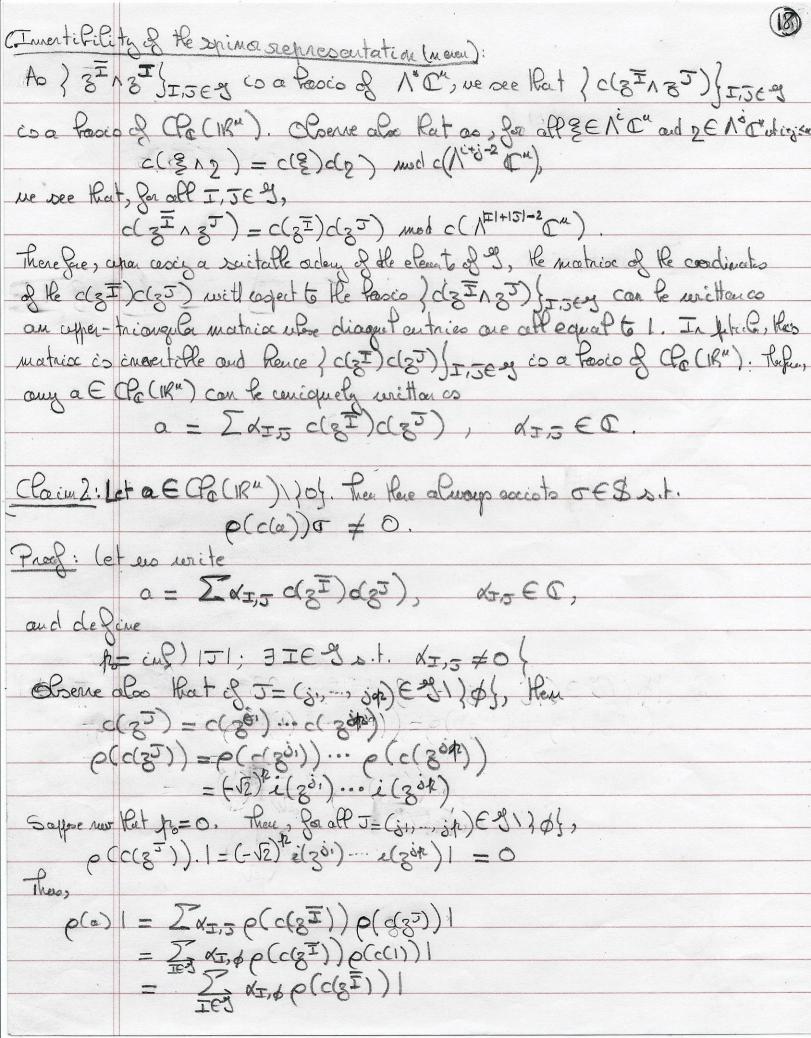
There is no induction shows that

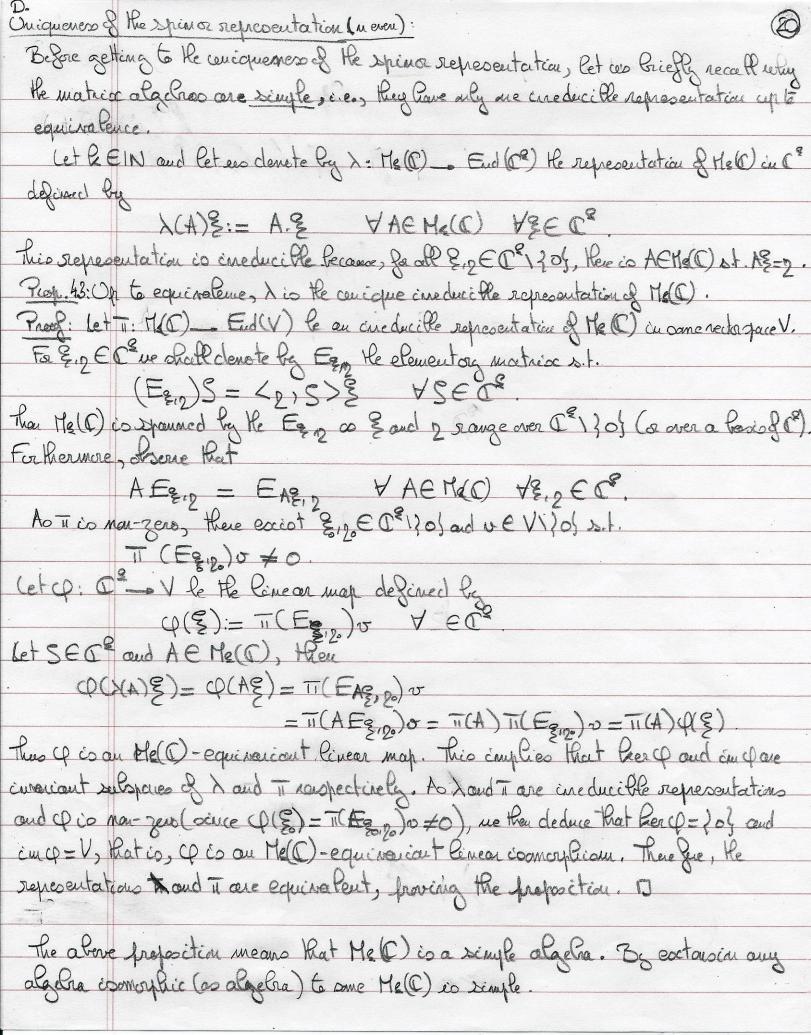
There is no induction of the continuity of it is not in the continuity of its notation.
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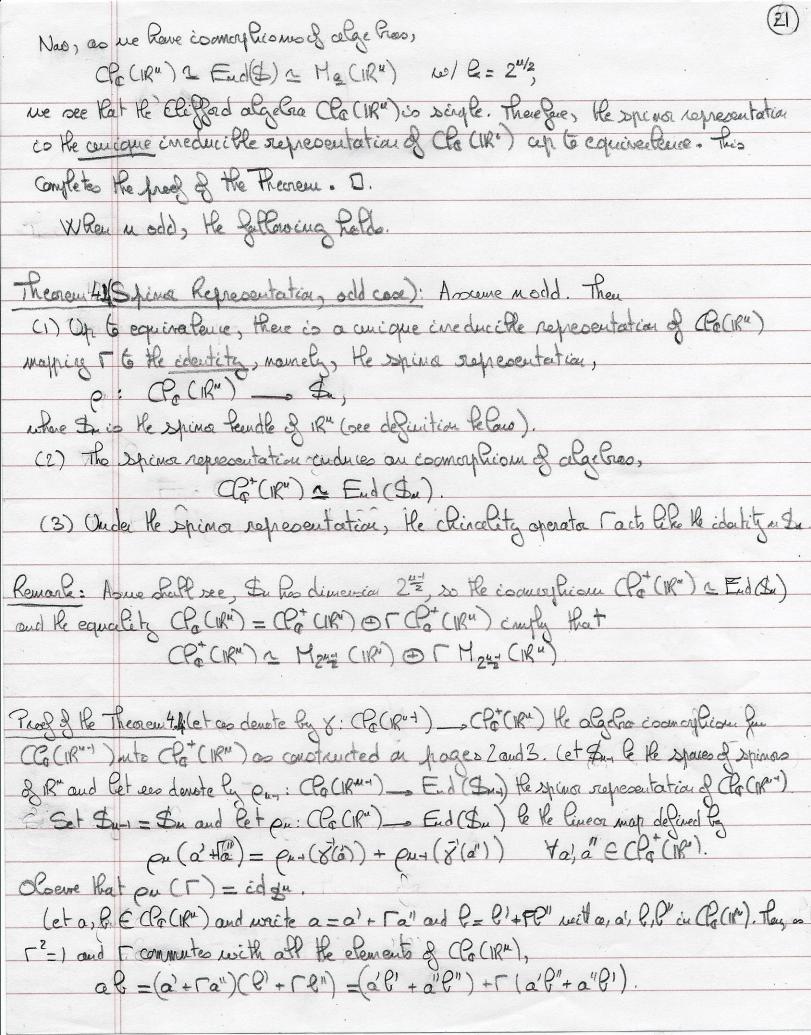


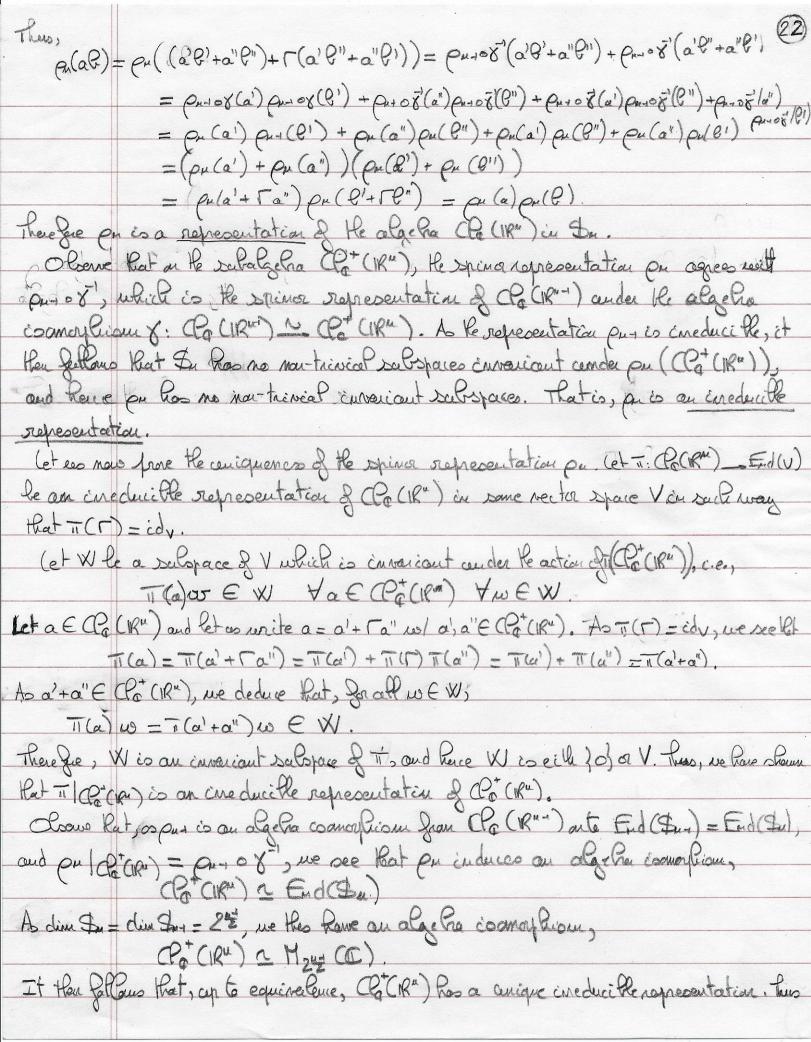


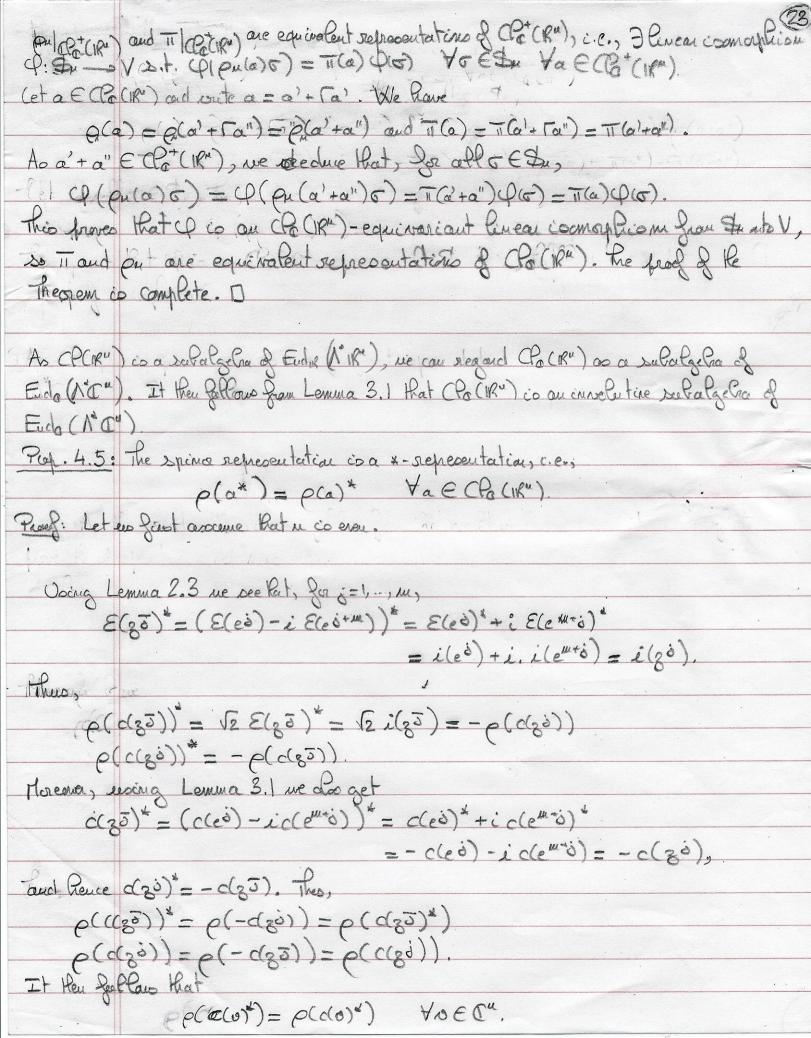


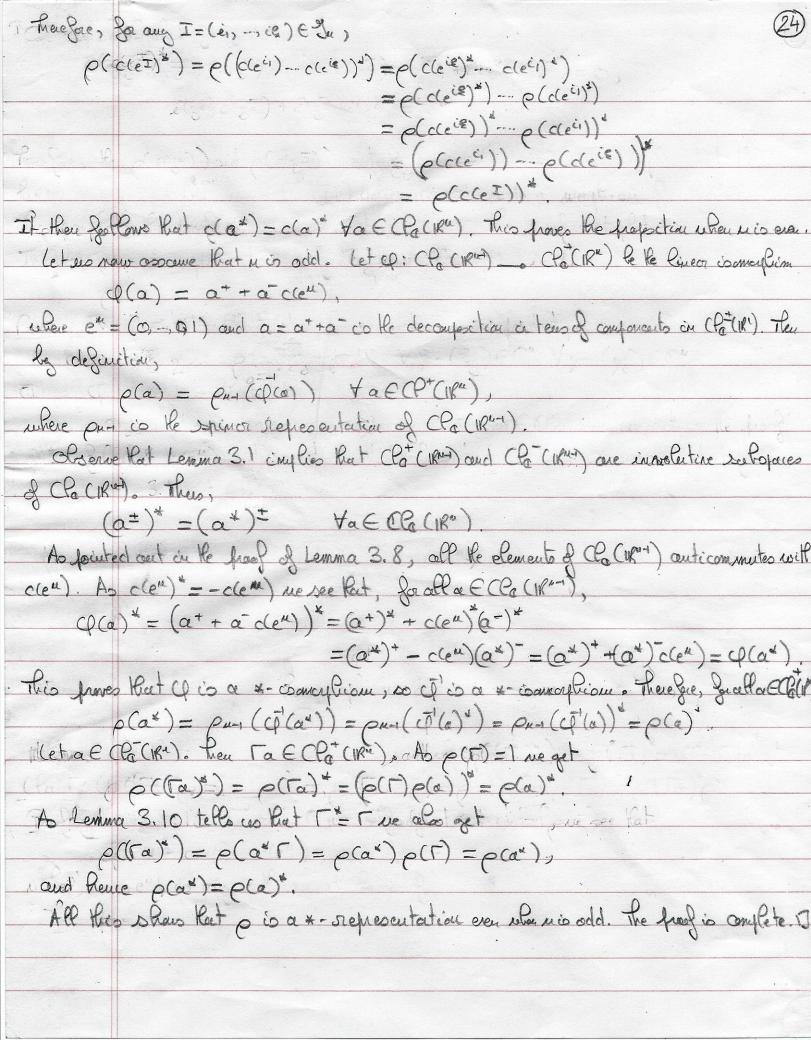


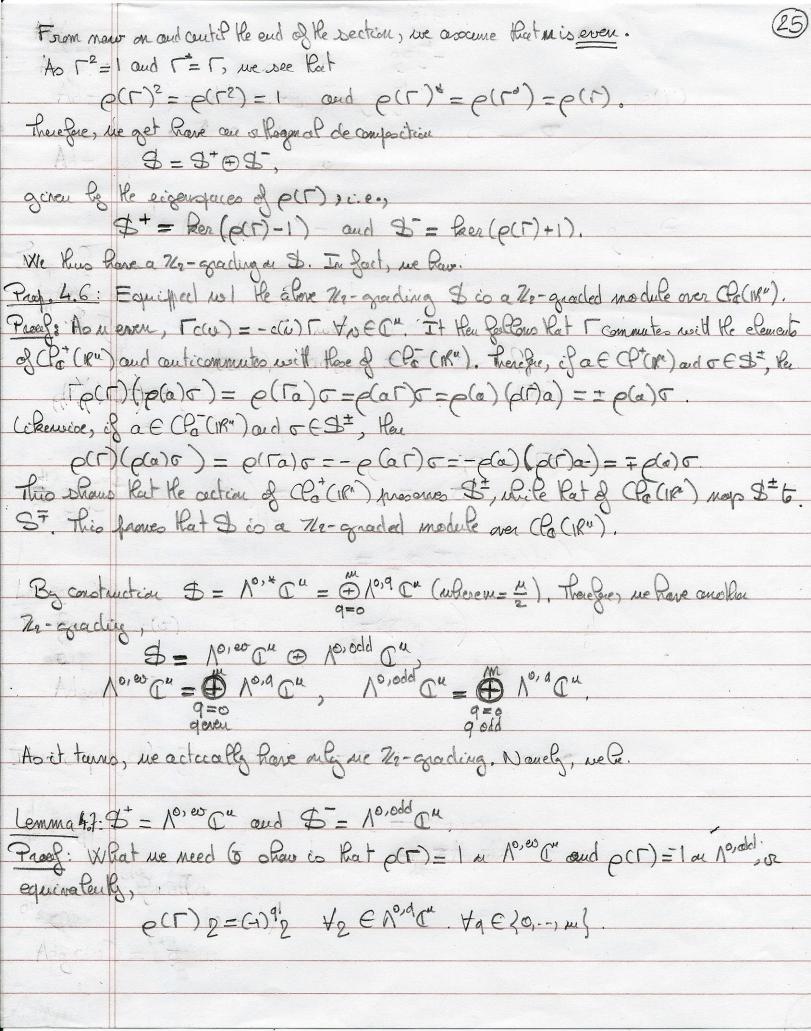






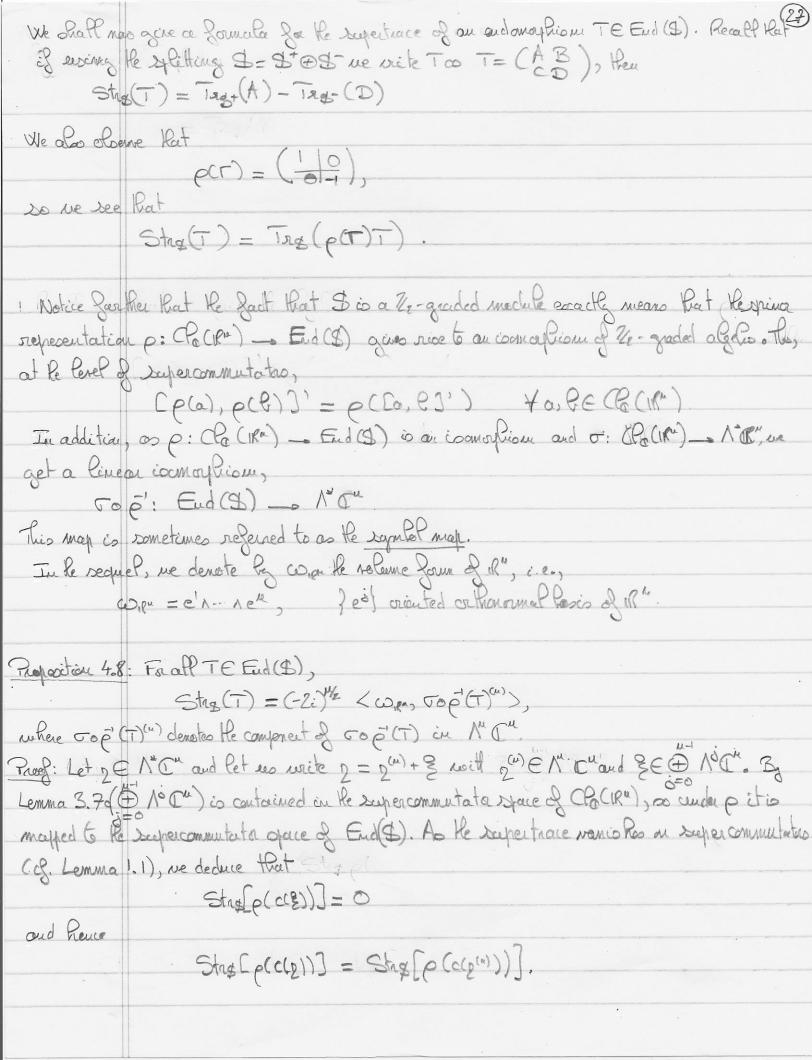






In E(35) io Spanned by 3 = 5 ID; we see But Ber Tij io Spanned by 3 = 5 ID; In eddition, de gine (1-211) 8 = ] - 8 = if if EI, [3] is id I. Therefore, FZI = C-DIZE, that is, for all q = 0, --, me F1 10,90m = G179 Observe also Pat if je I, Ken 2(23)E(85) 3= =0, but if i & I, Haon 1(30) E(30) 3I = 1(30) (30 V 3I) = 8I Therefore, we see that e(30) E(35) is the orthogonal projection who span 3 = (10) = Paer II; , that is , e(30) E(30) = 1-II; Theis, -211; = 2 i(3) E(3) -1 = - e(c(3)) e(c(3)) -1 = -e(c(3))c(3)+1). Observe garther that  $\frac{c(2\delta)c(2\delta) = \frac{1}{2} \left( c(e^{2\delta}) + i c(e^{2\delta+1}) \right) \left( c(e^{2\delta}) - i c(e^{2\delta+1}) \right)}{= \frac{1}{2} \left( c(e^{2\delta})^2 + c(e^{2\delta+1})^2 - i c(e^{2\delta}) c(e^{2\delta+1}) + i c(e^{2\delta+1}) c(e^{2\delta}) \right)}$ = - ( | + i c(e20) c(e20+1)) p(1-211) = (c(80)c(80)+1) = i c(e20)c(e20+1), 0 (F) = 11 i c(e26) d(e26+1) = e c(e1) c(e2) ... c(e24+1) c(e24+1) = r.

and Rence It then Julies or Keet facilia Elpining (CT) / NOIG = F/ NOIG = C-1)9 This completes He poof. D



and Ruce  $C(2^{(u)}) = \langle \omega_{i} | u, 2^{(u)} \rangle$  and  $\Gamma = i^{\frac{u}{2}} c(\omega_{i} | u)$ , we see that  $c(2^{(u)}) = \langle \omega_{i} | u, 2^{(u)} \rangle c(\omega_{i} | u, u) = (i)^{\frac{u}{2}} \langle \omega, 2^{(u)} \rangle \Gamma$ , Stay p(c(2)) = (-i) 4/2 < w, 2 (m) Stay (p(1)). Ao pCT) = (0-1), me see Pat Strig p(r) = Try+idy+ - Try(-idy-) = Try+idy+ + Tridy- = dim \$ + dim \$ = dim \$ = 24/2. There gie, we see Roct Str  $\mathfrak{sl}_{\mathcal{O}}(C(2))$ ] =  $(-2i)^{n/2} < \omega$ ,  $2^{(n)} > \forall 2 \in \Lambda^{*} \mathbb{C}^{n}$ Applying this to 2 = TOE' (T) with TE End (S) Hengines He proposition. D 5. Modules over (P(Ru): In the section, we assume a even and we love at 722-goded modules over Charm) Thanks to Prop. 4.6 we know that I is The graded module over Cla (18"). In the requer, we shall simply denote by (a, o) - at the artist of Cla (18") on I, i.e., a o = p(a) o. Moreovers & Prop. 4.5 the spinor representation is a +- sepresentation. Ho mentioned contin this implies but the garding operator pcr) is selfadyait and its eigenspaces 5 = tem (pcr) =1) one or hoagat complements of each other. Def.: let E le a Co (IRM) - module, i.e., there is a sepresentation IT: Co (IRM) - Endo (E) Then we say that the module E is selfadjoint when it is endowed with a Hermitian innet product with respect to which it is a x- representation, i.e., Tie) = Tie) + 4a Cla(18). Del: let E le a 7/2-graded Pa(18")-module associated & Re(21-graded) supresentation II: Palis)-Endle We say that E is a selfadjoint Zo graded module over Cack ) when E is equified soit a Hermitian fredect with respect to which - The The goding E = E+ DE is an or Ragnal de composition. - The representation II is a \* representation.

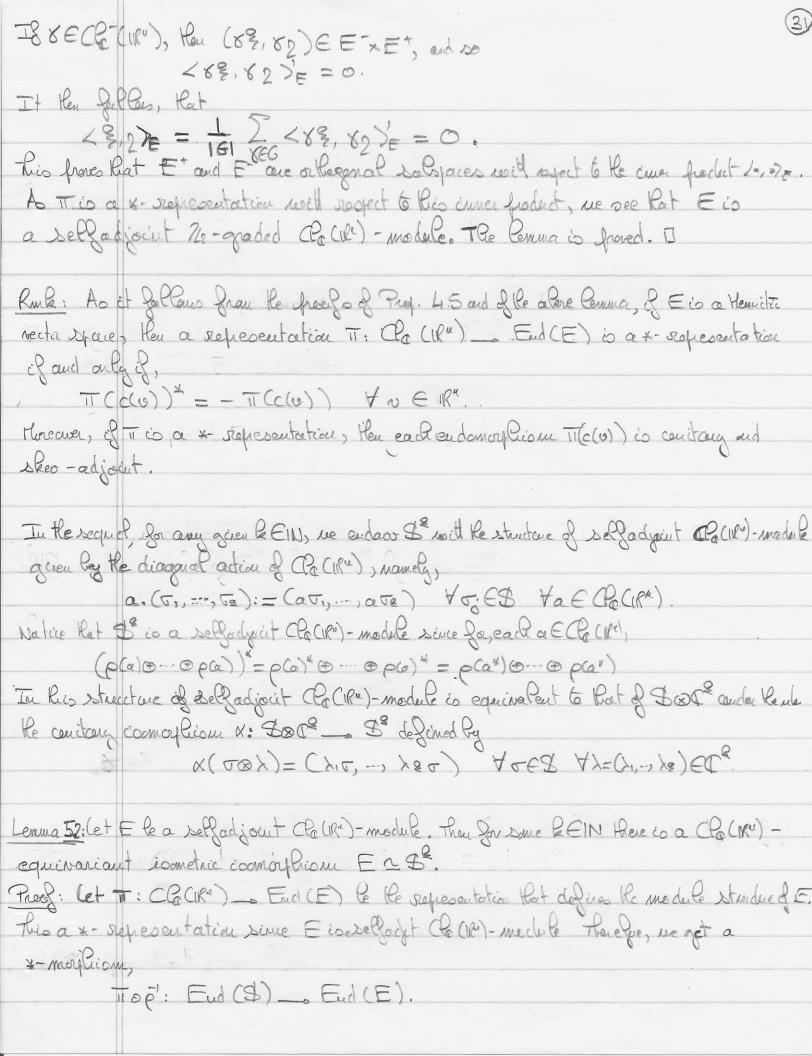
As explained above, the space of spinos & = & Di is ou example of anset godgoint Ze-graded (Pa(180) - module. Other examples are obtained as for Pais.

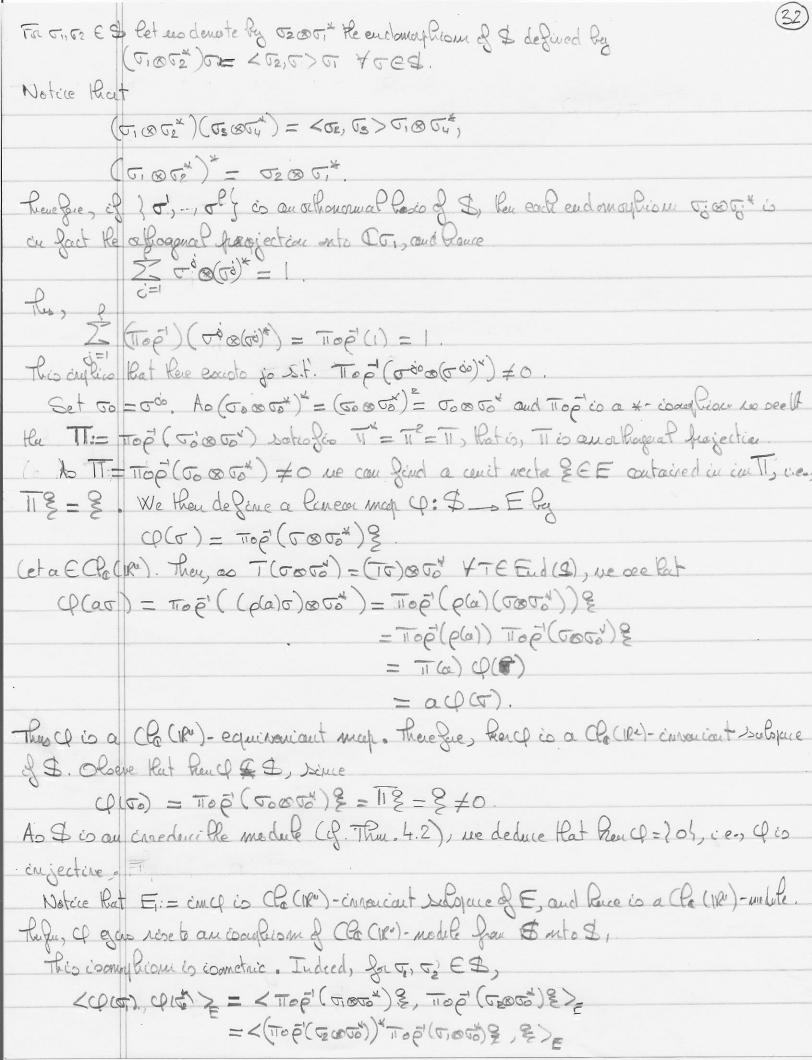
Cet W le a Hermitian space. Then Do Wio a Cla (184)-module with sepect to le

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action of Of (IR") defined by
That is, He corresponding super sociation is PBW: ERCIR") - End (SBW).

The denote by Lo, 2 and Lo, 2 where respective Hermitian consequents of D and W, Rec
ne con endow Do W wet the Hermitian inner froduct defind ly
          4 5,000, 5,000, 200 = <5, 52> < 101, 102> +5, E& Tig EW.
 let a E Cla (ik). Then, for all of ES and no; EW, j=1,2, we have
     < (PQ (0) (TQ WI), TEQ WE) = < 518 W, (PQ &) (TEQ WE) FOW
                                   = < 0 8001, (06) (2) 8002 gow
                                    = < 51, p(0) 52 & < 1001 2 NDE 7 W
                                    = < p(a) * T1, T2 } < 101, 102 } W
                                    = < p(ax) 5, 52 } < 101, 102 Zy
                                   = <(pols)(0)((Tiow), Teow) (0)
This frames that
       (poss)(a) = (poss)(a) Ya e CPc(18).
That is, pols is a x- Depresentation and SBW is a selfadjoint Clo(12)-modele.
 Supper man that W has an orthogonal Zo-grading W = W & W. Then & & W is
equipped with a Tr-grading given on Section 1, i.e., SOW - (SOW) + (SOW), with
         ($\otimes W) = ($^{\dagger} \otimes W) \oplus ($^{\dagger} \otimes W^{-}),
         ($@W) = ($t@W-) @($t@W+).
Then (Do W) and (Do W) are or hogoral salopares. Hereway,
        CROTURN (SOW) = (CROTURN) S) ON W,
 and Rence
        CP+(IRe)(Sow)= CP+(IRe)(StoW+OBFOW-)
                              C (CP+(IN) $=) & W+ (D) (CP+(IN) $=) & W-9
                              C ($ = & W +) = ($ = & W -) = ($ = & W) = .
Similarly,
        CPO (R°) (S&W) = C (S&W) =.
There fore, $00 W is a selfadjoint Za-graded Co (184) - module.
  We shall be module Soo W a trycoted Clalled-module. Howe shall now see any
Clo (IRª) - module (of finete din cuoia) is isomorphic to a twisted module.
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In the sequel by Cla(Me) - module we will exclusively mean a finite dimensional module
Lawwo 51.6 + Floo Polypul-module.
· (i) We celway can endow Excit a Hernitian inner fractact will sofert to whice E Roma
a seladjoint Co (We)- module.
(ii) IS E io 7/2 - goded, Ren use even con choose the inner federt so Pat Eio a salfodyit
1/2-graded Cla (IR)-module.
Proof: (et ) e',, e" , le au alhornounal loois of 18th and define;
G = ) ±19 0 ) to(ei) c(eie); 15 0 < < c25 mg.
Then G spans (Ca CIR"). Horanon, the Chifford relations,
$c(e^{i})^{2}=-1$ , $c(e^{i})c(e^{i})=-c(e^{i})c(e^{i})$ , $i\neq i$ ,
comple that Gir actually a Sinite Dolorage & Pa (180).
(et <., e) le a Hermitian cuner product on E, and define a new Hermitian curren froduct
Lo, DEME Reg
Lo, DEME By  Le, DE = 161 866 < 82, 82 E.  Thun Lo, DE is a G-invariant meaning But TIG: G = End(E) is a curitary
There Lo, . > = io a G-invariant meaning But TIG: G_ End(E) is a cuitary
sepresentation. In particular, lon j=1,-, u.
$\pi(c(c\delta))^* = \pi(c(c\delta))^*$
As (ceò) = -1 and c(eò) = -c(eò) (cf. Lemma 3.1) we see fut
$(\pi(c(e^{i}))^{2} = \pi(c(e^{i})^{2}) = \pi(-1) = -1,$
and Rence
$\pi(c(e\circ))' = -\pi(c(e\circ)) = \pi(c(e\circ)) = \pi(c(e\circ))'$
Rus,
TI (c(ceò)) = T(c(ceò)) Y; E } 1,-, u/.
We then can argue as in the freely of Prop. 4.5 & frame that
$\pi(\alpha)^* = \pi(\alpha^*)  \forall \alpha \in Cl_{\alpha}(\mathbb{R}^n).$
· Patio, Tio a +- representation and E is a selfadjoint (Pa(IR)-madule.
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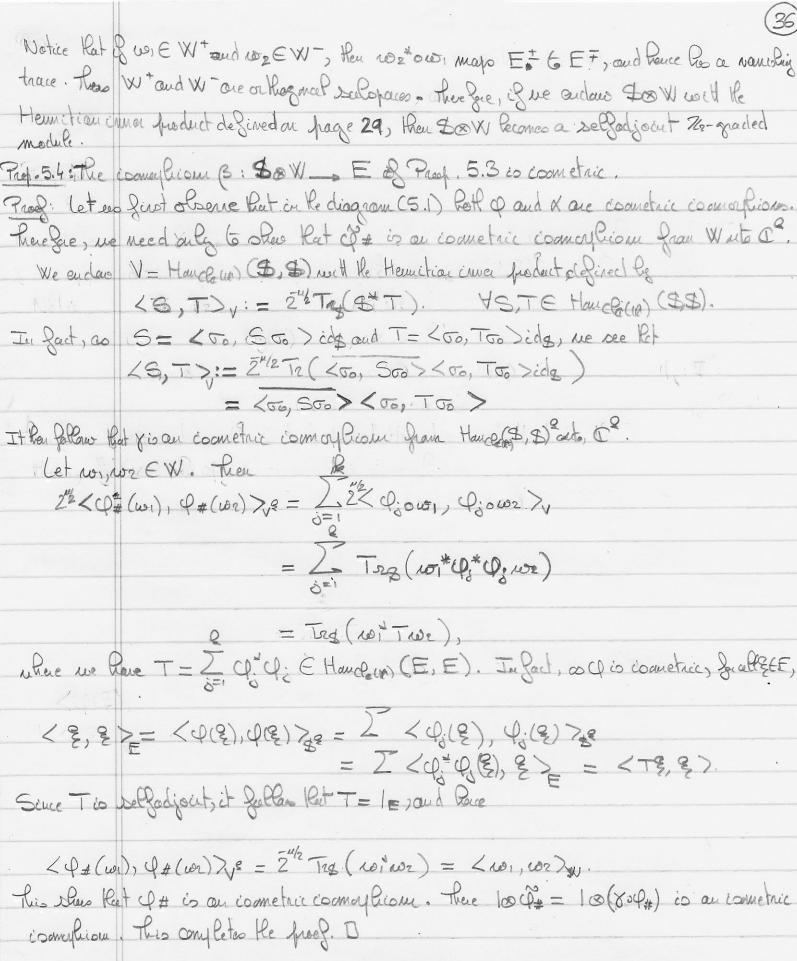
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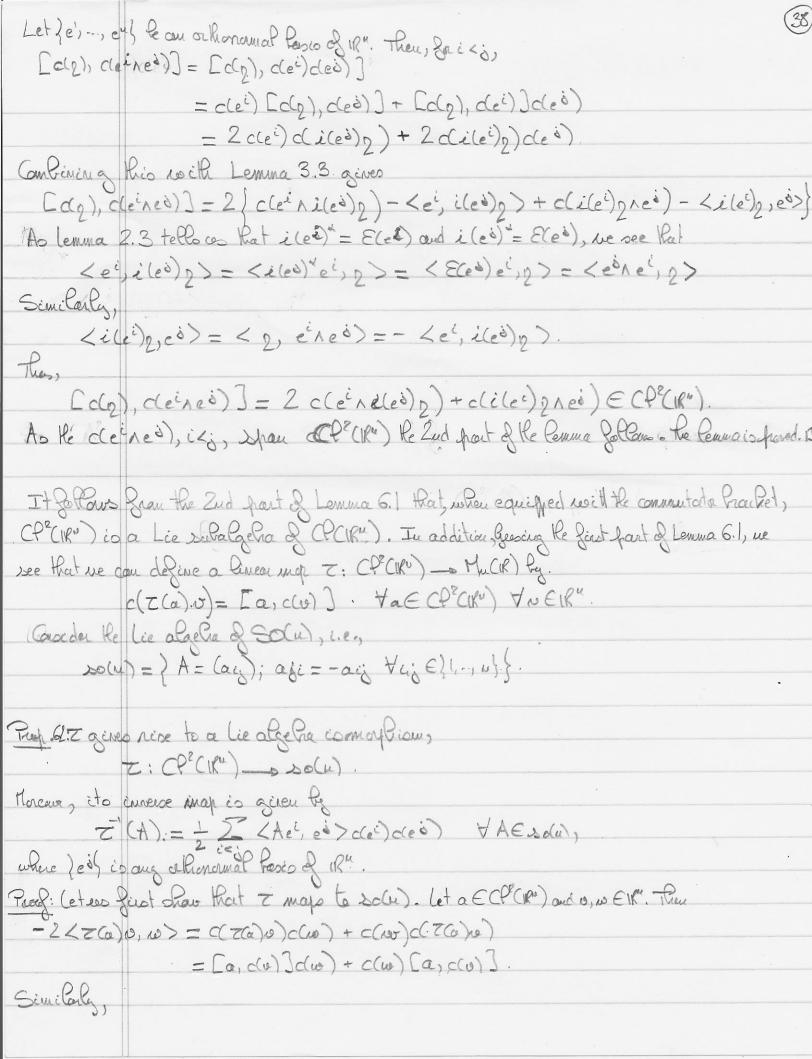
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Finally, let E= SoW le a twisted 7/2- gracled moderle and Cla (180), where WEWOW is a superspace. Let TE End (W). Then Sta Bow (TT) = Stag (p(r)) Staw (T) As shown in the proof of Prop. 4.8, Stryper) = 21/2, so we see But Strywer = 21/2 Stry (T), that is, Strw(T) = 2 Strgow (TT) This famula motivato le gallaving definitia. Def. (Relative Superstrace): Let E le a 7/2-opaded Ca(18) - modele. Par le relative superstrace en E is le lineau James Streps: Hampern (E, E) - O defined by Str E/2 (T) = 2"2 Str = (TT) YTE Houngelier) (E, E). Ruh letS, TE HomoPolar (E, E). As Formules resit SandT, FCS, TJ' = CFS, TJ', and hauce Streng [S,T]' = 2 the Streng [ES,T]' = 2 the [TS,T]' = 0.
Thus Streng is a supertrace on Homogener (E, E), in he souse Rutit vanishes on super-commutatoro. 6. The Spin Group Spin (u): Define CP'(IR") = c(IR") and CP2(IR") = c(12 IR"). (commaGici) [a,dw) ] ECP'(IR") Ya ECPE(IR") Yo EIR" (ii) [a, P] E CP2 (IR") Ya, PECPE(IR"). Proof: (et ) E 12(1Ru) and v EIR" Then By Lemma 3.6, [d()),c(v)] = [d()),c(v)] = 2c(i(v)) E(P(K"). This frames He first part of the Cemma.

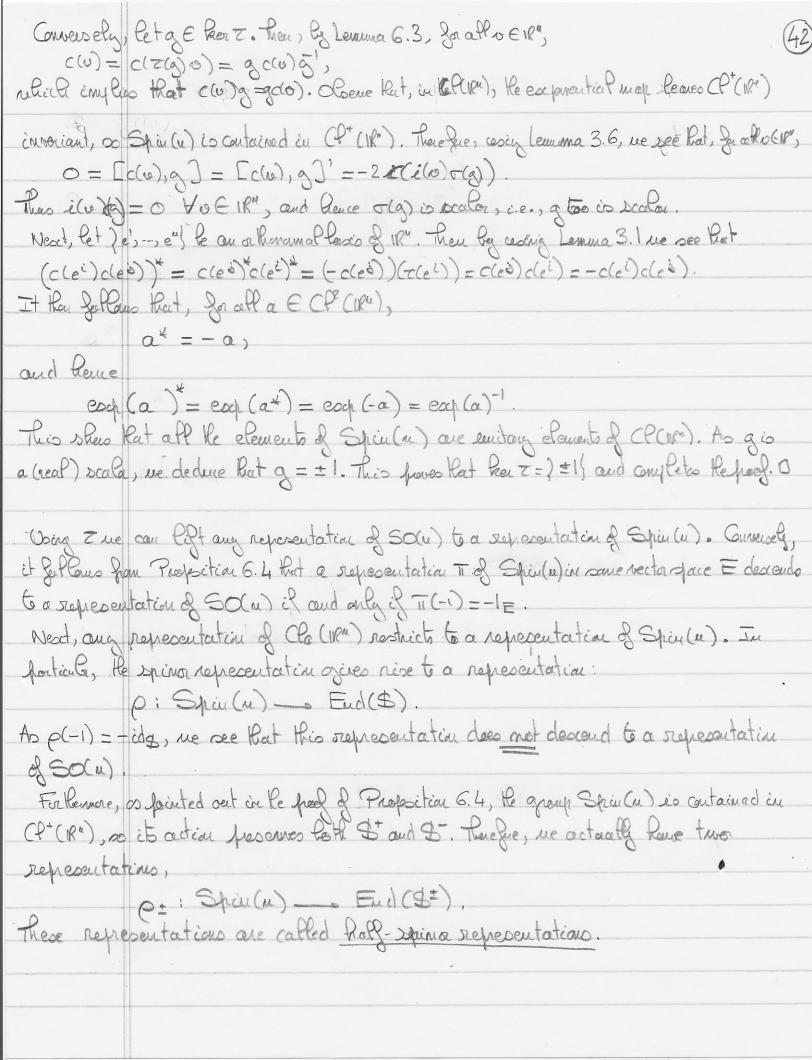


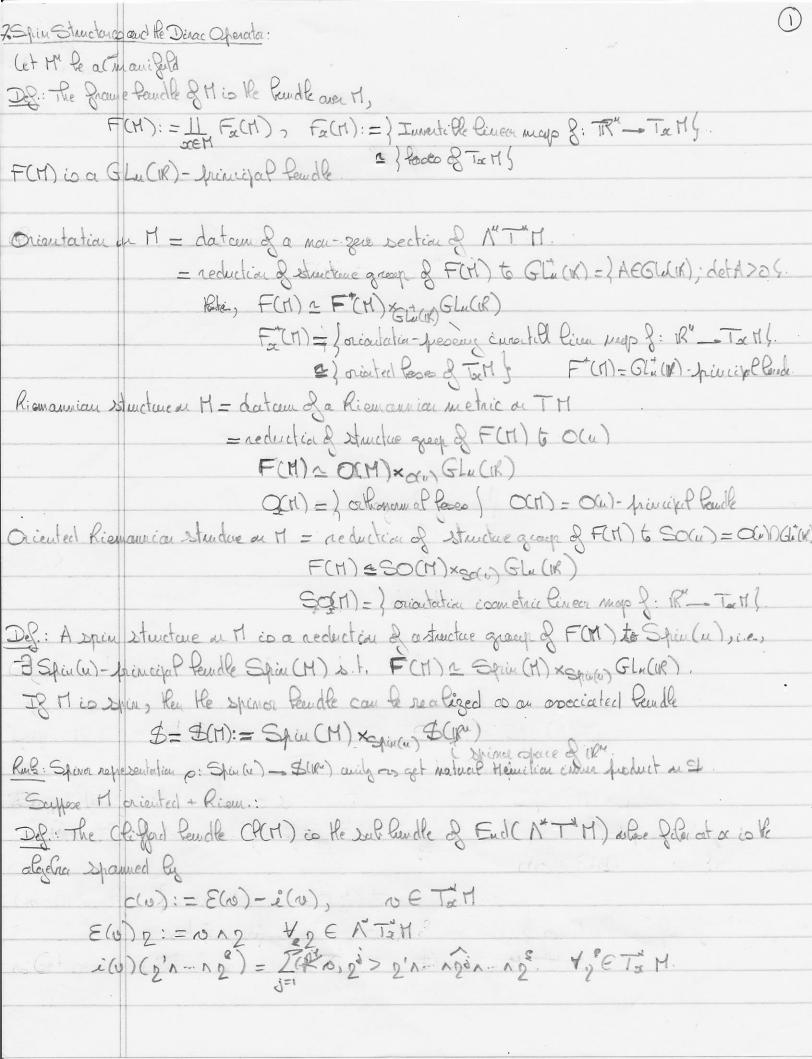
-2 (c) = (a) w> = ((a) c(t(a) w) + c(t(a) w) c(v) = c(w) [a, dw) ] + [a, dw) ]cho). - 2 / < Z(a) 10, 10 > + < 0, Z(a) 10 > = [a, ((v)) ] d(w) + ((v)) [a, d(w)] + ((w)) [a, d(v)] + (a, d(v)) d(v) = [a, c(v)c(w)] + [a,c(w)dv)] = [a, c(v)((v)) + ((v))(v)] - [a, -2<10,w>] There gie, for all 10, 10 E 18", < Z(a)0, w>+ < 10, Z(a) 10>=0. This shows fat Z(a) @ 20(u). Next, let a, le CPE(IR) and of IKM. Then c(T(a, e]) = [[a, e], c(u)] = [ab, c(u)] - [ba, c(u)] = a [B, d(v)] + Ca, d(v)] P - P [a, d(v)] - [P, d(v)] a = [a, [B, du)]] - [B, [a, du)]] = [a, c(z(P)o)] - [P, c(z(a)o)] =c(76).(78))-c(78).(76))=c([76], 7(B)]v). Z([a, e]) = [Z(a), Z(e)] ta, e e cpe(pu), that is, I is a marphism of Lie algebras. Let ) e', ... , en J le au orthonormal Perio of IR" and for eight, in let un dente le Eig the clementary matrices s.t., graff &= 1, , us Eijel=Silei=(ei,ei)el Then so (u) Ras Pasis ) En - En; i < j. Indeed, gray A E solul, A = \( \langle \langle Ae^i, e\) \( \in \) \( [ ((e') ((e)), ((e))]= (c(e'ne)), e]=-2c(i(e)(e'ne)) =- Ke, e) c(e) +2 Ke, e) c(e). Thus, Z(dei)dei)e2)= 5([dei)dei), de)])=-2/ei, 2/ei, 2/ei, 2/ei = 2 Eije - 2 Ejie.

Therefre se see that (6.1) Z(c(e)) = 2(Eij-Eii). It hen follows that I mappinto so(u). As facei)alei); is found ) Ein - Ein; is one respectives Passes & CPCIR") and solve) and Rave some containably, we see Rat Solve) and so (u) have some dimension. The fore, he fact that I map ato solu) actally inglis let I is an isomorphism from Pran (PECIRE) anto solu). Finally, (6.1) where hat, for itis, Z (Eció - Fire) = = c(ec) c(e) Tus, for all AE solut,

Z(A) = Z < Aei, eò > Z'(Euj-Euj) = 12 < Act, ei 7 clet) cles). The pool is complete. I Def. The sip in group of 18", denoted Spin (M), is the Lie group oftained by superentiating the Lie algebra CP2(18") inside CPC18"). Ex: (et vand up le vilhagenal auct vectes in  $\mathbb{R}^n$ . Then  $C(u)^2 = C(u)^2 = -1 \quad \text{and} \quad C(u) C(u) = -C(u) C(v).$ Muss;  $(c(u)c(w))^2 = c(u)c(w)c(w)c(w) = -c(w)c(w)^2c(w) = c(w)^2 = -1.$ An easy induction then show that, graft REINO,  $(c(u)c(w))^{22} = (-i)^2 \quad \text{and} \quad (c(u)c(w))^{22+1} = (-i)^2 c(u)c(w).$ Therefore, for all  $t \in \mathbb{R}$ ,  $\infty$   $\operatorname{ear}(t \, dv)c(w)) = \sum_{\delta=0}^{\infty} \frac{t^{\delta}}{\delta!} (c(v)c(w))^{\delta}$  $= \frac{1}{2} \frac{1}{2} \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \right) \left( \frac{1}{2} \left( \frac{1}{2} \left( \frac{1}{2} \right) \left( \frac{1}{2} \right) \right) \right) \right) \right) \right)}{1} \right) \right) \right) \right) \right) \right)$ 

Recognizing the Maclamin series of cost and sint we see Rat each (tale)alus) = coot + sint ale alus. In forticelle, for t= Ti ue see Pat -1 is contained in Spin (a). The Lie algebra comorphism Z: (PCIP) - solve) descendo to a maphism of Lie graps, Z: Spin (n) \_ 50(n). such that z(esch(a)) = esch (t(a)) ta E(P(iP)). Lomma 6.3: (et g E Spin (u) and let as regard g and g'os efemants of (P(p)). Than g (1) g = c( Z(g) or ) YNE 18". Proof: Let us write a = each (a) with a E (Pe(18"), Thou d) exp (ta) c(u) exp (-ta) ( = exp (ta) [a, c(u)] exp (-ta) = exp (ta) c(z(a)o) exp (-ta). A simple induction then show that, for all & EIN, and hence exp(ta)c(v)eap(-ta) = eap(ta)c(Z(a)v))eap(-ta), of lead (40) c(n) and (-40) = c(2(0) 0) = q c(600th (42(0)) 0) | t=0  $=\frac{d}{dt}\,c\left(\left.z\left(exp(ta)\right)o\right)\right|_{t=0}$  As foll eafta)clu) eaf(ta) and  $c\left(\left.z\left(exp(ta)\right)o\right)$  one analytic functions of t me deduce that they agree everywhere. In ferticity, get = 1 me get e exp(a) c(v) eap (-a) = c(Z(exp(a))v), a c(0) = c(z(g)0). The Cemma to these fraved. D Prop. 64: The Lie group maphion Z: Spirlu) -> Solu) is a double overing of Solu), that is, Z io onte and her Z= = = 1 !! Proof: Ao Z: (Pe(1k) \_ solu) is a isomafism and esp: solu) \_ solu) we see that zmaps
Spin(u) onto Solu). Mocowa, a shown in the passions example, -1 is outained in Spin(u) and by Lemma 6.3, ga all v ∈ 12", c(z(-1)10) = (-1) c(u) (-1) = c(0), i.e., ZC-1) = 1. Two } = 1) is autained in then Z.







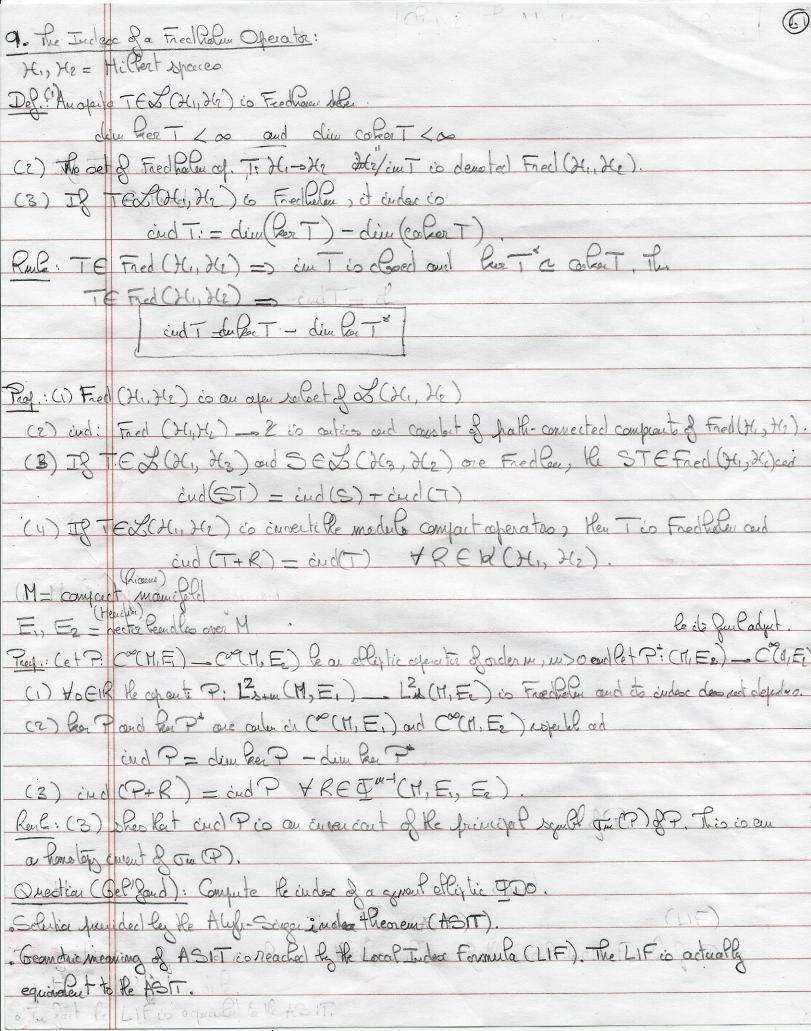
Ex.2: E = 1 M w/ same action on in Ex.1, But w/ grading ATTH = RONC(FI), ハーナーニハナナーのハナーリ、 Γ = Chirality of = i = c(ω) con = roleme form = vdeta(a) doin - nda"

ΓΕ C (Π, End (Λα Τ Μ)) = e<sup>4</sup>n - ne<sup>4</sup> ) e<sup>i</sup>/ other ted of l. frame of

Rul: Γ = i = \*\*, \* = Hodge star operator \*: N T M N T M \* X NB = < X, B> WATE X, BENT THE Ex. 3: M spin, a conted. W= 7/2-goded oriented Hauition necta Rendle E=\$0W (\$0W)= (\$+0W+) (\$+0W-) 1 a. ( TOW ) = (OG) T) OW p: CPO(M) - End (B) spiror repertation E= \$100 W = touisted & Classical-mediale bendle. Rue: Cocally, Ma OCIR" R'as Spin so Coally spins leadle & aluep exist Prof. Cocally any Cliffed-modele female is of the four \$00 W Dig.: Let E le a Chifford-module Rendle. A Chifford connection on E io a connection VEC(ME)\_...

CONTINO E)-5 which presence its Clifford - model Pendle structure, Portio, graff XECTH, TH), (i)  $\nabla_{x}(C^{*}CH, E^{\pm})) \subset C^{*}CH, E^{\pm})$ (ii)  $\{\nabla_{x}\}, \{2\} + \{3\}, \nabla_{x}\} = X(\{2\}, \{2\})$   $\{7\}, \{2\} \in C^{\infty}(M, E)$ (iii)  $\{\nabla_{x}, \pi(\omega)\} = \pi(\nabla_{x}\alpha)$   $\{\alpha \in C^{\infty}(M, Q(M))\}$   $\{\nabla^{Q(M)}, Q(M)\}$ Ex.1: E = Notif, Vita (evi Coota consection lifted & Notifica Cliffed connection Ex2: M spin + originated W= 7/2 - gooded Here. V connection on W satisfying (c)+ (ic). NSOW: = V®1+1®V is a Clifford connection on \$0 W. Prof. Locathy an Chiffed connection is & the few Thow Let Ele a Chilord-module bende w/ Chilord connection VE. RECOM, NITH & End (TM)) = Riemannian curvature We can left RIGE as REECOCH, NoTHORENDED ENTH RE = 1 2 < R(Di, D) 22, ep /(doindal) (c(e)) TI (c(e))) Ex: E=B, 28 = convertice of spin connection Vs.

Prof. The converture FEE COCH, NoTITO End (E) & FF is of the form where FFBE COCH, 12To Mas Home Am (E, E) is called the twisting connetone. Rul: Homoger) (E, E) - ) TEEnd (E); (TT, T(a)) = 0 ta E(Pa(H)). Ex: E = Sow ~ FEIS = 100 FW, FW = connoted TW. VE - V 1+ 180 W De The Dirac aperata associated & (E, VE) is the lost ada differential aperata DE: CM, E) - CM, E) giren by the composition, COCH, E) VE COCH, T'HOE) TOC, COCH, E) 000 3 - 1((0))3 Rule: (ocally, DE = 2 in (c(dxi)) VE Ex.1: E = NoTT to DE = d+d" (de Rham operata) d= exclusive diffustal d: Col, ATT) - Colo, ASTAI Ex 2: E = SOW VE - VOI+ 10 V Diar of Th = (~(1, \$) \( \frac{1}{2} \) (~(1, \frac{1}{1} \odd) \( \frac{1}{2} \) (~(1, \frac{1}{2} \odd) \( \frac{1}{2} \) (~(1, \frac{1}{2} \odd) \( \frac{1}{2} \odd) \) (~(1, \frac{1}{2} \odd) \( \frac{1}{2} \odd) \\ \frac{1}{2} \odd) \( \frac{1}{2} \odd) \( \frac{1}{2} \odd) \\ \frac{1}{2} \odd) \( \frac{1}{2} \odd) \\ \frac{1}{2} \odd) \( \frac{1}{2} \odd) \\ \frac{1}{2} \odd) \\ \frac{1}{2} \odd) \( \frac{1}{2} \odd) \\ 100000 - (E((b))2)00 m as Dw = Drol + (col) (10.7") = twisted Dirac of. Rul: Cocally De is & le Dr Prof. (1) De is selfant and effific. (2) DE = (ODE), DE: (C(1, E=), OF (DE) Prop. (Lickmorowicz Forme): DE 2 = (VE) VE + 1 x + JE/2 JE/4 E COO(H, End (E)) A FIB = 1 FFF (ei, e; ) T (c(e) c(e)) ) ei loc. ork. from of TT 201 dual afram of TOM



#### CHAPTER 8

# The Atiyah-Singer Index Theorem

The main reference for this chapter is the monograph [**BGV**]. We also refer to [**Sh**, §8] for the section on Fredholm operators. The proof of the Atiyah-Singer's local index formula given in Section 8.12 is taken from [**Po**].

We refer to the scans on the course's website for the notes of Sections 1–9 about Clifford algebras, spin structures, Dirac operators and Fredholm operators, and for the notes of the appendix on characteristic classes.

## 8.10. The Local Index Formula of Atiyah-Singer

Let  $(M^n,g)$  be an even dimensional compact oriented Riemannian manifold and let  $E=E^+\oplus E^-$  be a Clifford-module bundle equipped with a Clifford connection  $\nabla^E$ . We shall denote by  $D_E:C^\infty(M,E)\to C^\infty(M,E)$  the associated Dirac operator as defined in Section 8.

As we saw in Section 8, the operator  $D_E$  is elliptic and takes the form,

$$D_E = \begin{pmatrix} 0 & D_E^- \\ D_E^+ & 0 \end{pmatrix}, \qquad D_E^{\pm} : C^{\infty}(M, E^{\pm}) \to C^{\infty}(M, E^{\pm}).$$

Moreover, the selfadjointness of  $D_E$  implies that  $(D_E^{\pm})^* = D_E^{\mp}$ .

As  $D_E$  is elliptic, the results of Section 9 show that  $D_E$  is Fredholm. However, as  $D_E$  is selfadjoint its index must be zero. However, the ellipticity of  $D_E$  implies that  $D_E^{\pm}$  is elliptic, and hence is Fredholm. We then define the index of  $D_E$  to be

$$\operatorname{ind} D_E := \operatorname{ind} D_E^+ = \dim \ker D_E^+ - \dim \ker D_E^-,$$

where we have used the fact that  $D_E^-$  is the adjoint of  $D_E^+$ .

Let us denote by  $\operatorname{Str}_{E/\$}$  the relative supertrace on the fibers of  $\operatorname{Hom}_{\operatorname{Cl}_{\mathbb{C}}(M)}(E,E)$ . Recall that

$$\operatorname{Str}_{E/\$}(T) = 2^{-\frac{n}{2}} \operatorname{Str}_{E}(\Gamma T) \qquad \forall T \in \mathbb{C}^{\infty}(M, \operatorname{Hom}_{\operatorname{Cl}_{\mathbb{C}}(M)}(E, E)),$$

where we have denoted by  $\Gamma$  the chirality operator (or, more precisely, the section End E given by the pointwise action of chirality operator on the fibers of E).

In the sequel, we denote by  $F^{E/\$}$  the twisted curvature of the Clifford connection  $\nabla^E$  (see Section 8 for its precise definition).

Definition 8.10.1. The relative Chern form of the twisted curvature  $F^{E/\$}$  is

$$\operatorname{Ch}(F^{E/\$}) = \operatorname{Str}_{E/\$}\left[e^{-F^{E/\$}}\right] \in C^{\infty}(M, \Lambda^{ev}T^*M).$$

For instance, suppose that M is spin and let  $W = W^+ \otimes W^-$  be a  $\mathbb{Z}_2$ -graded Hermitian vector bundle with Hermitian connection  $\nabla^W$  preserving its  $\mathbb{Z}_2$ -grading.

1

Then we can we form the twisted Clifford-bundle module  $E = \mathcal{S} \otimes W$  and equipped with the twisted Clifford connection,

$$\nabla^E = \nabla^{\mathcal{S}} \otimes 1_W + 1_{\mathcal{S}} \otimes \nabla^W.$$

As alluded to in Section 8, the twisted curvature of this twisted connection is just  $1_{\mathcal{S}} \otimes F^W$ . Therefore, its relative Chern form is equal to

$$\operatorname{Ch}(F^{E/\$}) = \operatorname{Str}_{E/\$} \left[ 1_{\$} \otimes e^{-F^{W}} \right] = \operatorname{Str}_{W} \left[ e^{-F^{W}} \right]$$
$$= \operatorname{Tr}_{W^{+}} \left[ e^{-F^{W^{+}}} \right] - \operatorname{Tr}_{W^{-}} \left[ e^{-F^{W^{-}}} \right]$$
$$= \operatorname{Ch}(F^{W^{+}}) - \operatorname{Ch}(F^{W^{-}}),$$

where  $F^{W^{\pm}}$  is the curvature of the connection on  $W^{\pm}$  induced by  $\nabla^{W}$  and  $\operatorname{Ch}(F^{W^{\pm}})$  is its usual Chern form.

In addition, as explained in the appendix on characteristic classes, the  $\hat{A}$ -form of the Riemann curvature  $R^M$  of M (i.e., the curvature of the Levi-Civita connection on TM) is

$$\hat{A}(R^M) := \det^{\frac{1}{2}} \left\lceil \frac{R^M/2}{\sinh(R^M/2)} \right\rceil.$$

We are now in a position to state the index formula of Atiyah-Singer.

THEOREM 8.10.2 (Atiyah-Singer). We have

(8.1) 
$$\operatorname{ind} D_E = (2i\pi)^{-\frac{n}{2}} \int_M \left[ \hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$}) \right]^{(n)},$$

where  $\left[\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$}\right]^{(n)}$  denotes the n-th degree component of the even form  $\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$})$ .

Remark 8.10.3. The local index formula of Atiyah- Singer continue to hold even M is non-orientable. Notice that the integrand  $\left[\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\mathscr{F}}\right]^{(n)}$  defines a density, namely,

(8.2) 
$$\langle \left[ \hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$} \right]^{(n)}, \omega_M \rangle v_g(x),$$

where  $\omega_M$  is the volume form and  $v_g(x)$  is the Riemannian density. The integral  $\int_M \left[\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$}]^{(n)}\right]$  then is the integral of this density. The above process allows us to identify n-the degree forms and densities, but it

The above process allows us to identify n-the degree forms and densities, but it depends on a choice of orientation. Observe that the definition of the relative Chern form  $Ch(F^{E/\$})$  depends on the definition of relative supertrace which involves the chirality operator, hence depends on the choice of the orientation (see Eq. (8.6). However, if at a point  $x \in M$ , we reverse the orientation of  $T_xM$  then it only affects by a change of sign the values at x of the chirality operator  $\Gamma$  and relative Chern form  $Ch(F^{E/\$})$ . Since this similarly affects the value at x of the volume form, we see that the value at x of the density (8.2) is actually independent of the orientation of  $T_xM$ . Therefore (8.2) defines a density even M is non-orientable. Then the index of  $D_E$  continue to be given by the formula (8.1), where the r.h.s. is interpreted as a multiple of the integral of the density (8.2).

The above index formula is often referred to as the *local index formula* of Atiyah-Singer. It is not a mere special case of the *full index theorem* of Atiyah-Singer [**AS**] for general elliptic  $\Psi$ DOs. In many respects, the local index formula is equally important, if not even more important, than the full index theorem.

As it turns out, K-theoretic arguments show that the local index formula is actually equivalent to the full index theorem (see, e.g.,  $[\mathbf{ABP}]$ ). More importantly, its only the case of Dirac operators on Clifford-module bundles that the Atiyah-Singer's index theorem reaches its true geometric contents. Therefore, the local index formula for Dirac operators is often confused with the full index theorem of Atiyah-Singer.

We shall describe some of the geometric applications of the Atiyah-Singer's local index formula in the next section. Before doing this let us briefly outline of the proof this formula.

As  $D_E^2$  is a selfadjoint operator with non-negative spectrum, the Borel functional calculus allows us to define the heat semi-group  $e^{-tD_E^2}$ ,  $t\geq 0$ , as a family of bounded operators on  $L^2(M,E)$  (cf. Chapter II). Indeed, if  $(\xi_k)_{k\geq 0}$  is an orthonormal basis of eigenvectors such that  $D_E^2\xi_k=\lambda_k(D_E^2)\xi_k$  for all  $k\geq 0$ , then

$$e^{-tD_E^2}\xi_k = e^{-t\lambda_k(D_E^2)}\xi_k \qquad \forall k \ge 0.$$

In addition, we denote by Str the supertrace on the  $\mathbb{Z}_2$ -graded Hilbert space  $L^2(M,E)=L^2(M,E^+)\oplus L^2(M,E^-)$ , that is,

$$\operatorname{Str}\left(\begin{array}{cc} A & B \\ C & D \end{array}\right) = \operatorname{Trace}(A) - \operatorname{Tr}(B)$$

for any trace-class operator 
$$T=\left( \begin{array}{cc} A & B \\ C & D \end{array} \right)$$
 on  $L^2(M,E)=L^2(M,E^+)\oplus L^2(M,E^-).$ 

Let t > 0. The ellipticity of  $D_E^2$  ensures us that, for t > 0, the operator  $e^{-tD_E^2}$  is smoothing, and hence has a smooth Schwartz kernel, i.e., a section  $k_t(x,y)$  in  $C^{\infty}(M \times M, E \boxtimes (E^* \otimes |\Lambda|(M)))$ , which is called the *heat kernel* of  $D_E^2$  (see Section 8.12 on this point). It then follows that  $e^{-tD_E^2}$  is trace-class and its supertrace is given by

$$\operatorname{Str} e^{-tD_E^2} = \int_M \operatorname{Str}_E[k_t(x,x)],$$

where  $Str_E[k_t(x,x)]$  is defined as a density on M.

Theorem 8.10.4 (McKean-Singer Formula). For all t > 0,

ind 
$$D_E = \operatorname{Str} e^{-tD_E^2} = \int_M \operatorname{Str}_E[k_t(x,x)].$$

PROOF. We know that the Dirac operator  $D_E$  is of the form,

$$D_E = \left( \begin{array}{cc} 0 & D_E^- \\ D_E^+ & 0 \end{array} \right),$$

where  $D_E^{\pm}: C^{\infty}(M, E^{\pm}) \to C^{\infty}(M, E^{\pm})$  and  $(D_E^{\pm})^* = D_E^{\mp}$ . Thus,

$$D_E^2 = \left( \begin{array}{cc} D_E^- D_E^+ & 0 \\ 0 & D_E^+ D_E^- \end{array} \right).$$

Let t > 0. Then

$$e^{-tD_E^2} = \begin{pmatrix} e^{-tD_E^- D_E^+} & 0\\ 0 & e^{-tD_E^+ D_E^-} \end{pmatrix}.$$

Therefore,

$$\begin{aligned} \operatorname{Str} e^{-tD_E^2} &= \operatorname{Trace} \left[ e^{-tD_E^- D_E^+} \right] - \operatorname{Trace} \left[ e^{-tD_E^+ D_E^-} \right] \\ &= \sum_{\lambda \in \mathbb{C}} e^{-t\lambda} \dim E_{\lambda}^+ - \sum_{\lambda \in \mathbb{C}} e^{-t\lambda} \dim E_{\lambda}^- \\ &= \dim E_0^+ - \dim E_0^- + \sum_{\lambda \in \mathbb{C} \setminus 0} e^{-t\lambda} (\dim E_{\lambda}^+ - \dim E_{\lambda}^-), \end{aligned}$$

where, for all  $\lambda \in \mathbb{C}$ , we have set

$$E_{\lambda}^+ := \ker(D_E^- D_E^+ - \lambda) \quad \text{and} \quad E_{\lambda}^- := \ker(D_E^+ D_E^- - \lambda).$$

Notice that, as  $(D_E^{\pm})^* = D_E^{\mp}$ , we have

$$E_0^{\pm} = \ker D_E^{\mp} D_E^{\pm} = \ker (D_E^{\pm})^* D_E^{\pm} = \ker D_E^{\pm}.$$

Thus,

$$\dim E_0^+ - \dim E_0^- = \dim \ker D_E^+ - \dim \ker D_E^- = \operatorname{ind} D_E,$$

and hence

(8.3) 
$$\operatorname{Str} e^{-tD_E^2} = \operatorname{ind} D_E + \sum_{\lambda \in \Gamma \setminus 0} e^{-t\lambda} (\dim E_{\lambda}^+ - \dim E_{\lambda}^-).$$

Next, let  $\lambda \in \mathbb{C} \setminus 0$  and let  $\xi$  be in  $E_{\lambda}^+ = \ker(D_E^- D_E^+ - \lambda)$ . Then

$$D_E^+ D_E^- (D_E^+ \xi) = D_E^+ (D_E^- D_E^+ \xi) = \lambda D_E^+ \xi,$$

that is,  $D_E^+\xi$  is contained in  $E_\lambda^- = \ker(D_E^+D_E^- - \lambda)$ . Thus  $D_E^+$  induces a linear map from  $E_\lambda^+$  to  $E_\lambda^-$ .

Similarly, the operator  $D_E^-$  induces a linear map from  $E_{\lambda}^-$  and  $E_{\lambda}^+$ . As by the very definitions of  $E_{\lambda}^+$  and  $E_{\lambda}^-$  we have

$$\left(D_E^- D_E^+\right)_{|E_\lambda^+} = \lambda \operatorname{id}_{E^+} \quad \text{and} \quad \left(D_E^+ D_E^-\right)_{|E_\lambda^-} = \lambda \operatorname{id}_{E^-},$$

we see that  $D_E^+$  induces a linear isomorphism from  $E_{\lambda}^+$  onto  $E_{\lambda}^-$ . Thus,

$$\dim E_{\lambda}^{+} = \dim E_{\lambda}^{-} \quad \forall \lambda \in \mathbb{C} \setminus 0.$$

Combining this with (8.3) yields the McKean-Singer formula.

Now, the local index formula of Atiyah-Singer is proved by combining the McKean-Singer formula with the following:

THEOREM 8.10.5 (Patodi, Gilkey, Atiyah-Bott-Patodi). In  $C^{\infty}(M, |\Lambda|(M))$ ,

$$\operatorname{Str}_{E}[k_{t}(x,y)] \longrightarrow \left[\hat{A}(R^{M}) \wedge \operatorname{Ch}(F^{E/\$})\right]^{(n)}$$
 as  $t \to 0^{+}$ .

Theorem 8.10.5 is often referred to as the *local index theorem*. In the above general form, it was first proved by Atiyah-Bott-Patodi [**ABP**] by making use of Riemannian invariant theory (see also [**Gi**]). The first purely analytical proofs were produced by Getzler [**Ge1**] and Bismut [**Bi**] by means of different approaches. These proofs were further simplified by the celebrated "short proof" of Getzler [**Ge2**],

where was introduced the use of the so-called "Getzler's rescaling". This proof paved the way for many generalizations of the local index formula of Atiyah-Singer'.

In Section 8.12, we reproduce the proof of the local index theorem given in [**Po**]. This proof is especially relevant for applications in noncommutative geometry (cf. Chapter 11).

Notice, that it can be shown that in  $C^{\infty}(M, |\Lambda|(M) \otimes \operatorname{End} E)$ ,

$$k_t(x,x) \sim t^{-\frac{n}{2}} \sum_{j \ge 0} t^j a_j(D_E^2)(x)$$
 as  $t \to 0^+$ ,

where  $a_0(D_E^2)(x)$  is the positive multiple of the Riemannian density (cf. Section 8.12). Therefore, it is a very striking fact that under the supertrace all the divergent terms in the above asymptotics vanish and we obtain a convergent quantity.

It should also be stressed out that this phenomenon is really specific to Dirac operators coming from Clifford connections; in general the local index theorem does not occur for other operators.

### 8.11. Geometric Applications

Let  $(M^n,g)$  be an even-dimensional oriented compact Riemannian manifold. Depending on the existence of an additional structure (e.g. spin structure, complex structure) the Atiyah-Singer index formula has various striking consequences. We give here a brief overview of them, referring to the book of Berline-Getzler-Vergne for a more detailed treatment, including proofs of the results stated here.

**8.11.1.** Dirac Operators on Spin Manifolds. Assume that M is spin and let us denote by  $\mathcal{S}$  its spinor bundle. In addition, let  $W = W^+ \otimes W^-$  be a Hermitian  $\mathbb{Z}_2$ -graded vector bundle equipped with a Hermitian connection  $\nabla^W$  preserving the  $\mathbb{Z}_2$ -grading of W.

We denote by  $\mathcal{D}_W: C^\infty(M, \mathfrak{F} \otimes W) \to C^\infty(M, \mathfrak{F} \otimes W)$  the associated twisted Dirac operator, that is,  $\mathcal{D}_W$  is the Dirac operator associated to the twisted Clifford connection  $\nabla^{\mathfrak{F} \otimes W} := \nabla^{\mathfrak{F}} \otimes 1 + 1 \otimes \nabla^W$ , where  $\nabla^{\mathfrak{F}}$  is the spin connection. As shown before, the twisted curvature of  $\nabla^{\mathfrak{F} \otimes W}$  is  $1 \otimes F^W$ , where  $F^W$  is the

As shown before, the twisted curvature of  $\nabla^{\mathcal{S} \otimes W}$  is  $1 \otimes F^W$ , where  $F^W$  is the curvature of the connection  $\nabla^W$ , and its relative Chern form agrees with the Chern form  $Ch(F^W)$  of  $F^W$ . Therefore, we the local index formula of Atiyah-Singer gives

THEOREM 8.11.1 (Atiyah-Singer). We have

$$\operatorname{ind} {\not \!\! D}_W = (2i\pi)^{-\frac{n}{2}} \int_M \left[ \hat{A}(R^M) \wedge \operatorname{Ch}(F^W) \right]^{(n)}.$$

In particular, in the case of the Dirac operator  $\not \!\! D_M$  of M acting on spinors (i.e., when W is the flat trivial line bundle over M) we get

THEOREM 8.11.2 (Atiyah-Singer). We have

$$\operatorname{ind} \mathcal{D}_{M} = (2i\pi)^{-\frac{n}{2}} \int_{M} \hat{A}(R^{M})^{(n)}.$$

As an immediate corollary we obtain:

Corollary 8.11.3. If the  $\hat{A}$ -number of M, i.e.,

$$\hat{A}(M) := \int_M \left[ \hat{A}(R^M) \right]^{(n)},$$

is not an integer, then M does not admit a spin structure.

In the sequel, we denote by  $\kappa_M$  the scalar curvature of M.

THEOREM 8.11.4 (Lichnerowicz). Suppose that  $\kappa_M(x) > 0$  for all  $x \in M$ . Then the null space of  $\mathcal{D}_M$  is reduced to  $\{0\}$ .

PROOF. This is a direct consequence of the Lichnerowicz's formula,

$$\mathcal{D}_M^2 = (\nabla^{\mathcal{S}})^* \nabla^{\mathcal{S}} + \frac{1}{4} \kappa_M.$$

It implies that, for all  $u \in C^{\infty}(M, \mathcal{S})$ ,

$$\langle \not\!\!D_M u, u \rangle = \langle \nabla^{\sharp} u, \nabla^{\sharp} u \rangle + \frac{1}{4} \langle \kappa_M u, u \rangle \ge \frac{1}{4} \int \kappa_M(x) \|u(x)\|_{\mathscr{F}_x}^2 v_g(x) \ge c \|u\|_{L^2(M, \mathfrak{F})}^2,$$

where  $v_q(x)$  is the Riemannian density and we have set  $c := \inf_{x \in M} \kappa_M(x)$ .

As M is compact and  $\kappa_M$  is positive, the constant c is positive. Therefore, no non-zero smooth section of  $\mathcal{S}$  can be contained in  $\ker \mathcal{D}_M^2$ . As  $\mathcal{D}_M$  is elliptic and selfadjoint,  $\ker \mathcal{D}_M^2 = \ker \mathcal{D}_M \subset C^\infty(M, \mathcal{S})$ , so we see that  $\ker \mathcal{D}_M$  is trivial.  $\square$ 

If  $\not \!\! D_M$  has a trivial null space, then, obviously, ind  $\not \!\! D_M=0$ . Therefore, combining Theorem 8.11.2 and Theorem 8.11.4 gives

Theorem 8.11.5. If M carries a metric of positive scalar curvature, then its  $\hat{A}$ -number  $\hat{A}(M)$  vanishes.

**8.11.2. The Chern-Gauss-Bonnet Theorem.** The (complexified) exterioralgebra bundle  $\Lambda_{\mathbb{C}}^*T^*M$  carries the  $\mathbb{Z}_2$ -grading,

$$\Lambda^* T^* - C^{\infty}(M) = \Lambda^{\operatorname{ev}} T^* M \oplus \Lambda^{\operatorname{odd}} T^* M.$$

This is a Clifford-module bundle with respect to the natural action of  $\mathrm{Cl}_{\mathbb{C}}(M)$  seen as a sub-bundle of  $\mathrm{End}(\Lambda_{\mathbb{C}}^*T^*M)$ . Because we can define another  $\mathbb{Z}_2$ -grading by means of the chirality operator (see next subsection), we shall denote by  $\Lambda^{\mathrm{ev}/\mathrm{odd}}T_{\mathbb{C}}^*M$ , or simply  $\Lambda^{\mathrm{ev}/\mathrm{odd}}$  the Clifford-module bundle defined by means of the above  $\mathbb{Z}_2$ -grading.

As it turns out, the Levi-Civita connection  $\nabla^{\Lambda_{\mathbb{C}}^*T^*M}$  on  $\Lambda_{\mathbb{C}}^{\mathrm{ev}/odd}T^*M$  is a Clifford connection and the associated Dirac operator agrees with the de Rham operator,

$$d + d^* : C^{\infty}(M, \Lambda^*T^*M) \longrightarrow C^{\infty}(M, \Lambda^*T^*M),$$

where d is the usual exterior differential and  $d^*$  is its adjoint (see [BGV]).

The index of  $d + d^*$  has an important topological interpretation as follows. For j = 0, ..., n let us denote by  $H^k(M)$  the (k+1)-th de Rham cohomology group of M, that is,

$$H^k(M) = \ker d_k / \operatorname{im} d_{k-1},$$

where  $d_k: C^{\infty}(M, \Lambda^k T^*M) \to C^{\infty}(M, \Lambda^{k+1} T^*M)$  is the de Rham differential on forms of degree k. The Euler characteristic of M then is

$$\chi(M) := \sum_{k=0}^{n} (-1)^k \dim H^k(M).$$

In addition, let  $\Delta := (d+d^*)^2 = dd^* + d^*d$  be the Laplace-Beltrami operator and let us denote by  $\Delta_k := d_{k-1}d_{k+1}^* + d_k^*d_k$  its restriction to forms of degree k.

Proposition 8.11.6 (Hodge; see [BGV]). We have an orthogonal splitting,

$$\ker d_k = \operatorname{im} d_{k-1} \oplus \ker \Delta_k,$$

and hence

$$H^k(M) \simeq \ker \Delta_k$$
.

That is, any class in  $H^k(M)$  can be represented by a unique harmornic form of degree k.

It follows from this that dim  $H^k(M) = \dim \ker \Delta_k$ , and hence

(8.4) 
$$\chi(M) = \sum_{k=0}^{n} (-1)^k \dim \ker \Delta_k$$

$$= \dim \ker \Delta_{|\Lambda^{\text{ev}}T^*M} - \dim \ker \Delta_{|\Lambda^{\text{odd}}T^*M}$$

$$= \dim \ker (d + d^*)_{|\Lambda^{\text{ev}}T^*M} - \dim \ker (d + d^*)_{|\Lambda^{\text{odd}}T^*M}$$

$$= \operatorname{ind}(d + d^*).$$

In the sequel, we denote by  $\text{Eul}(R^M)$  the Euler form of the curvature  $R^M$  of M. Regarding  $R^M$  as a section of  $(\Lambda^2 T^*M) \otimes \operatorname{End}(TM)$ , the Euler form is the form of degree of n on M such that, for any local orthonormal frame  $\{e_i\}$  of TM with dual coframe  $\{e^i\}$ , we have

$$\frac{1}{(n/2)!} \left( -\frac{1}{4} \sum_{i,j,k,l} R^M_{ijkl} (dx^i \wedge dx^j) \otimes (e^k \wedge e^l) \right)^{\frac{n}{2}} = \operatorname{Eul}(R^M) \otimes (e^1 \wedge \ldots \wedge e^n),$$

where we have set  $R_{ijkl}^M := \langle R^M(\partial_i, \partial_j) e_k, e_l \rangle$ . LEMMA 8.11.7 (See [**BGV**]).

(1) The twisted curvature  $F^{(\Lambda^{ev/odd})/\$}$  of the Levi-Civita connection on  $\Lambda^{ev/odd}_{\mathbb{C}}T^*M$ is such that, for any local orthonormal frame  $\{e_i\}$  of TM with dual coframe  $\{e^i\}$ , we have

$$F^{(\Lambda^{ev/odd})/\$} = -\frac{1}{4} \sum_{i,j,k,l} R^{M}_{ijkl}(dx^{i} \wedge dx^{j}) \otimes \left(\varepsilon(e^{k}) + \iota(e^{k})\right) \left(\varepsilon(e^{l}) + \iota(e^{l})\right).$$

(2) The relative Chern form of the twisted curvature  $F^{(\Lambda^{ev/odd})/\$}$  is given by

(8.5) 
$$\operatorname{Ch}\left(F^{(\Lambda^{ev/odd})/\$}\right) = i^{\frac{n}{2}}\operatorname{Eul}(R^{M}).$$

As the zeroth-degree component of the even form  $\hat{A}(R^M)$  is equal to 1, from (8.5) we get

$$\left[\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$})\right]^{(n)} = \left[\hat{A}(R^M)\right]^{(0)} \operatorname{Eul}(R^M) = \operatorname{Eul}(R^M).$$

Therefore, by using (8.4) and applying the local index formula of Atiyah-Singer we obtain

THEOREM 8.11.8 (Chern-Gauss-Bonnet Theorem). We have

$$\chi(M) = \text{ind}(d + d^*) = (2\pi)^{-n} \int_M \text{Eul}(R^M).$$

In the above form, the Chern-Gauss-Bonnet theorem is originally due to S.S. Chern. This is the most classical index theorem. When n=2 we recover the well-known Gauss-Bonnet theorem,

$$\chi(M) = \frac{1}{2\pi} \int_{M} \kappa_{M}(x) v_{g}(x),$$

where  $v_q(x)$  is the Riemannian density.

**8.11.3. Hirzebruch Signature Theorem.** Recall that the chirality operator is defined by

(8.6) 
$$\Gamma := i^{\frac{n}{2}} c(\omega_M) \in C^{\infty}(M, \operatorname{Cl}_{\mathbb{C}}(M)) \subset C^{\infty}(M, \operatorname{End}(\Lambda_{\mathbb{C}} T^*M)),$$

where  $\omega_M = \sqrt{\det g(x)} dx^1 \wedge \cdots \wedge dx^n$  is the volume form of M. As  $\Gamma^2 = 1$  it defines an alternative  $\mathbb{Z}_2$ -grading on  $\Lambda^*_{\mathbb{C}} T^* M$  given by

(8.7) 
$$\Lambda_{\mathbb{C}}^* T^* M = \Lambda^+ T_{\mathbb{C}}^* M \otimes \Lambda^- T_{\mathbb{C}}^* M, \qquad \Lambda^{\pm} T_{\mathbb{C}}^* M := \ker(\Gamma \mp 1).$$

This  $\mathbb{Z}_2$ -grading is preserved by the action of  $\mathrm{Cl}_{\mathbb{C}}(M)$ , and so we get a new Clifford-module bundle structure on  $\Lambda_{\mathbb{C}}^*T^*M$ . We shall denote by  $\Lambda_{\mathbb{C}}^{+/-}T^*M$  the  $\Lambda_{\mathbb{C}}^*T^*M$  equipped with this Clifford-module bundle structure.

The Levi-Civita connection on  $\Lambda_{\mathbb{C}}^{+/-}T^*M$  is a Clifford connection (cf. [**BGV**]), and so the associated Dirac operator is again the de Rham operator  $d+d^*$ . Notice that, as we are using a different  $\mathbb{Z}_2$ -grading, the index differs from that in (8.4). As we shall now explain the index that we get is intimately related to the signature of the manifold.

Let  $\star \in C^{\infty}(M, \operatorname{End} \Lambda^*T^*M)$  be the Hodge operator, i.e.,

$$\alpha \wedge \star \beta = \langle \alpha, \beta \rangle \omega_M \qquad \forall \alpha, \beta \in C^{\infty}(M, \Lambda^* T^* M),$$

Lemma 8.11.9 (See [BGV]). We have

$$\Gamma = (-1)^{\frac{k(k-1)}{2}} i^{\frac{n}{2}} \star \qquad on \ \Lambda_{\mathbb{C}}^{k} T^{*} M,$$
$$d^{*} = - \star d \star$$

It follows from this lemma that  $d + d^* = d - \star d\star$ . In the sequel, in order to distinguish with the index in (8.4), we shall denote by  $\operatorname{ind}(d - \star d\star)$  the index of  $d + d^*$  when using the  $\mathbb{Z}_2$ -grading defined by the chirality operator. That is,

$$\operatorname{ind}(d - \star d\star) := \dim \ker(d + d^*)_{|\Lambda^+ T^*_{\mathbb{C}} M} - \dim \ker(d + d^*)_{|\Lambda^- T^*_{\mathbb{C}} M}$$
$$= \dim \ker(d - \star d\star)_{|\Lambda^+ T^*_{\mathbb{C}} M} - \dim \ker(d - \star d\star)_{|\Lambda^- T^*_{\mathbb{C}} M}.$$

It also follows from Lemma 8.11.9 that the Hodge  $\star$ -operator descends to a linear map,

(8.8) 
$$\star : H^k(M) \longrightarrow H^{n-k}(M),$$

which is an isomorphism since  $\star^2 = (-1)^k$  on  $\Lambda^k T^* M$ . Using this isomorphism we can prove the nondegeneracy of the bilinear pairing,

$$H^k(M) \times H^{n-k}(M) \ni (\alpha, \beta) \longrightarrow \int_M \alpha \wedge \beta.$$

When n is divisible by 4 this pairing is symmetric on  $H^{\frac{n}{2}}(M) \times H^{\frac{n}{2}}(M)$ , and hence it gives rise to a nondegenerate quadratic form on  $H^{\frac{n}{2}}(M)$ . The signature of this quadratic form is called the *signature* of M and is denoted  $\sigma(M)$ . This is a topological invariant of M.

Lemma 8.11.10 (See  $[\mathbf{BGV}]$ ). If n is divisible by 4, then

$$\sigma(M) = \operatorname{ind}(d - \star d \star).$$

Because of this lemma, the operator  $d - \star d\star$  is often called *signature operator*. Next, the twisted curvature of the Levi-Civita on  $\Lambda^{+/-}T_{\mathbb{C}}^*M$  is continue to be given by (1). As we shall now see, since we are using a different  $\mathbb{Z}_2$ -grading, we obtain a different relative Chern form, which we shall denote by  $\operatorname{Ch}(F^{(\Lambda^{+/-})/\$})$ .

Recall that L-form  $L(R^M)$  is the of the curvature  $R^M$  is defined by

$$L(R^M) := \det^{\frac{1}{2}} \left( \frac{R^M/2}{\tanh(R^M/2)} \right).$$

Lemma 8.11.11 (See [BGV]). We have

$$\operatorname{Ch}(F^{E/\$}) = 2^{\frac{n}{2}} \mathrm{det}^{\frac{1}{2}} \left( \frac{\sinh(R^M/2)}{R^M/2} \right) \wedge L(R^M).$$

Observe that  $\hat{A}(R^M) \wedge \det^{\frac{1}{2}} \left( \frac{\sinh(R^M/2)}{R^M/2} \right) = 1$ . Thus,

$$\hat{A}(R^M) \wedge \operatorname{Ch}(F^{E/\$}) = 2^{\frac{n}{2}} L(R^M).$$

Therefore, using the Atiyah-Singer index formula we obtain:

THEOREM 8.11.12 (Atiyah-Singer). We have

$$\operatorname{ind}(d - \star d\star) = (i\pi)^{-\frac{n}{2}} \int_{M} \left[ L(R^{M}) \right]^{(n)}.$$

Combining this with Lemma 8.11.10, we immediately get

Theorem 8.11.13 (Hirzebruch Signature Theorem). If  $\dim M$  is divisible by 4, then

$$\sigma(M) = \operatorname{ind}(d - \star d \star) = (i\pi)^{-\frac{n}{2}} \int_{M} \left[ L(R^{M}) \right]^{(n)}.$$

Since  $\left[L(R^M)\right]^{(4)}=-\frac{1}{96\pi^2}\operatorname{Tr}[(R^M)^2],$  we see that in dimension 4,

$$\sigma(M) = -\frac{1}{24\pi^2} \int_M \text{Tr}[(R^M)^2].$$

More generally, let W be a Hermitian vector bundle equipped with a Hermitian connection  $\nabla^W$ . Then we can form the twisted Clifford-module bundle  $\Lambda^{+/-}T^*_{\mathbb{C}}M\otimes W$ , where the fiberwise action of  $\mathrm{Cl}_{\mathbb{C}}(M)$  is such that, above all  $x\in M$ ,

$$a(\xi \otimes w) = (a\xi) \otimes w, \qquad \forall a \in \operatorname{Cl}_x(M) \quad \forall (\xi, w) \in \Lambda^* T_x^* M \times W_x.$$

The twisted connection  $\nabla^{\Lambda^{+/-}}T_{\mathbb{C}}^{*}M\otimes W=\nabla^{\Lambda^{*}}T^{*}M\otimes 1+1\otimes\nabla^{W}$  is a Clifford connection on  $\Lambda^{+/-}T_{\mathbb{C}}^{*}M\otimes W$ . We shall denote by  $(d-\star d\star)_{W}$  the associated Dirac operator. Its index is denoted  $\sigma(M,W)$  and is called the twisted signature with coefficients in W. It depends only on the topology of the manifold M and the vector bundle W.

The twisting curvature  $F^{(\Lambda^{+/-}T^*_{\mathbb{C}}M\otimes W)/\$}$  is equal to

$$F^{(\Lambda^{+/-})/\$} \otimes 1 + 1 \otimes F^W$$
.

where  $F^W$  is the curvature of  $\nabla^W$ . Thus, its relative Chern form is given by

$$\operatorname{Ch}(F^{(\Lambda^{+/-}T_{\mathbb{C}}^*M\otimes W)/\$}) = \operatorname{Ch}(F^{E/\$}) \wedge \operatorname{Ch}(F^W).$$

Therefore, we obtain:

Theorem 8.11.14 (Twisted Signature Theorem). We have

$$\sigma(M,W) = (d - \star d\star)_W = (i\pi)^{-\frac{n}{2}} \int_M \left[ L(R^M) \wedge \operatorname{Ch}(F^W) \right]^{(n)}.$$

The twisted signature theorem was originally proved by Atiyah-Singer. It is important because it can be used to prove the *full* Atiyah-Singer index theorem, i.e., the index theorem for *general* elliptic  $\Psi$ DOs on compact manifolds.

**8.11.4.** The Hirzebruch-Riemann-Roch Formula. In this subsection, we assume that M is a *complex* manifold of complex dimension n. This means that the tangent bundle TM is thus endowed with a complex structure  $J \in C^{\infty}(M, \operatorname{End}_{\mathbb{R}} TM)$ ,  $J^2 = -1$ , so that the holomorphic tangent bundle  $T_{1,0}M := \ker(J-i)$  is integrable in Fróbenius' sense, i.e.,  $[Z,W] \in C^{\infty}(M,T_{1,0}M)$  for all  $Z,W \in C^{\infty}(M,T_{1,0}M)$ .

In addition, we assume that the orientation of M is the orientation defined by its complex structure and the Riemannian metric on g TM is the real part of a Hermitian metric h on  $T_{\mathbb{C}}M$  with respect to which J is unitary.

The complexified tangent bundle  $T_{\mathbb{C}}M:=TM\otimes\mathbb{C}$  admits the orthogonal decomposition,

$$T_{\mathbb{C}}M = T_{1,0}M \oplus T_{0,1}M, \qquad T_{0,1}M := \overline{T_{1,0}M}.$$

By duality this gives rise to orthogonal splittings,

$$T_{\mathbb{C}}^*M=\Lambda^{1,0}T^*M\oplus\Lambda^{0,1}T^*M,\qquad \Lambda_{\mathbb{C}}^*T^*M=\bigoplus_{q=0}^n\Lambda^{p,q}T^*M,$$

where  $\Lambda^{p,q}T^*M := (\Lambda^{1,0}T^*M)^p \wedge (\Lambda^{0,1}T^*M)^q$  is the bundle of (p,q)-covectors. If  $\alpha \in C^{\infty}(M,\Lambda^{0,q})$ , then

$$d\alpha = \partial \alpha + \overline{\partial} \alpha,$$

where  $\partial \alpha$  is a (1,q)-form and  $\overline{\partial} \alpha$  is a (0,q+1)-form. Furthermore, the integrability of  $T_{1,0}$  implies that  $\overline{\partial}^2 = 0$ . We thus get a chain complex,

$$\overline{\partial}: C^{\infty}(M, \Lambda^{0, \bullet}T^*M) \longrightarrow C^{\infty}(M, \Lambda^{0, \bullet+1}T^*M).$$

This complex is called the *Dolbeault complex* of M and its cohomology groups are denoted  $H^{0,q}(M)$ ,  $q = 0, \ldots, n$ .

Let W be a holomorphic vector bundle over M. This means that M can be covered by open subsets  $U_{\alpha}$  over which there are local trivializations  $\tau_{\alpha}: W_{|U_{\alpha}} \to U_{\alpha} \times \mathbb{C}^r$  in such way that the transition maps are holomorphic maps. Then, for  $q = 0, \ldots, n$ , there exists a unique operator

$$\overline{\partial}_W: C^{\infty}(M, (\Lambda^{0,q}T^*M) \otimes W) \longrightarrow C^{\infty}(M, (\Lambda^{0,q+1}T^*M) \otimes W)$$

such that, for all local holomorphic frames  $e_1, \ldots, e_r$  of W and sections  $\omega = \sum \omega_j \otimes e_j$  of  $(\Lambda^{0,q}T^*M) \otimes W$ ,

$$\overline{\partial}_W(\sum \omega_j \otimes e_j) = \sum_j (\overline{\partial}\omega_j) \otimes e_j.$$

Since  $(\overline{\partial}_W)^2 = 0$  this gives rise to a chain complex, called the *Dolbeault complex* of M with coefficients in W. Its cohomology groups are denoted  $H^{0,q}(M,W)$ ,  $q = 0, \ldots, n$ .

Let us also endow W with a Hermitian metric  $\langle ., . \rangle_W$ . Together with the Hermitian metric on  $\Lambda^{0,*}T^*M$  this defines a Hermitian metric on  $(\Lambda^{0,*}T^*M) \otimes W$  and

an inner inner products on  $C^{\infty}(M,(\Lambda^{0,*}T^*M)\otimes W)$ , using which we define the formal adjoint  $\overline{\partial}_W^*:C^{\infty}(M,(\Lambda^{0,*}T^*M)\otimes W)\to C^{\infty}(M,(\Lambda^{0,*}T^*M)\otimes W)$ . The Dolbeault Laplacian then is

$$\square_W := (\overline{\partial}_W + \overline{\partial}_W^*)^2 = \overline{\partial}_W^* \overline{\partial}_W + \overline{\partial}_W \overline{\partial}_W^*.$$

Notice that  $\square_W$  maps  $C^{\infty}(M, (\Lambda^{0,q}T^*M) \otimes W)$  to itself. We then denote by  $\square_W^{0,q}$  its restriction to  $(\Lambda^{0,q}T^*M) \otimes W$ .

In addition, by looking at its expression in local coordinates, it is not difficult to check that  $\square_W$  has same principal symbol as a Laplacian and hence is elliptic. As by assumption M is compact, it follows that  $\ker \square_W$  is a finite dimensional subspace of  $C^{\infty}(M, (\Lambda^{0,*}T^*M) \otimes W)$ . We then define the holomorphic Euler-characteristic of M with coefficients in W to be

$$\chi(M,W) := \sum_{q=0}^{n} (-1)^{q} \dim H^{0,q}(M,W).$$

This is a holomorphic invariant of M and W.

PROPOSITION 8.11.15 (Hodge, see [**BGV**]). For q = 0, ..., n,

$$H^{0,q}(M,W) \simeq \ker \square_W^{0,q}$$
.

Since the operators  $\overline{\partial}_W + \overline{\partial}_W^*$  and  $\square_W$  have same kernel, it follows from Proposition 8.11.15 that  $H^{0,q}(M,W)$  is isomorphic to the kernel of the Dolbeault operator  $\overline{\partial}_W + \overline{\partial}_W^*$  on (0,q)-forms. Thus,

(8.9) 
$$\chi(M,W) = \operatorname{ind}\left(\overline{\partial}_W + \overline{\partial}_W^*\right),\,$$

where the index of  $\overline{\partial}_W + \overline{\partial}_W^*$  is defined to be

$$\dim \ker \left( \overline{\partial}_W + \overline{\partial}_W^* \right)_{|(\Lambda^{0,\operatorname{ev}}T^*M) \otimes W} - \dim \ker \left( \overline{\partial}_W + \overline{\partial}_W^* \right)_{|(\Lambda^{0,\operatorname{odd}}T^*M) \otimes W}.$$

Let  $\nabla: C^\infty(M,W) \to C^\infty(M,T^*_{\mathbb C}M\otimes W)$  be a connection on W. Using the splitting  $T^*_{\mathbb C}M=\Lambda^{1,0}T^*M\oplus\Lambda^{0,1}T^*M$  we can split  $\nabla$  as

$$\nabla = \nabla^{1,0} + \nabla^{0,1}.$$

where  $\nabla^{1,0}$  (resp.,  $\nabla^{0,1}$ ) maps to sections of  $(\Lambda^{1,0}T^*M)\otimes W$  (resp., sections of  $(\Lambda^{0,1}T^*M)\otimes W$ ). We then say that  $\nabla$  is a holomorphic connection when  $\nabla^{0,1}=\overline{\partial}_W$ . If  $\nabla$  is a holomorphic connection, then its curvature  $F^W$  is actually an (1,1)-form with values in End W.

Proposition 8.11.16. There is a unique holomorphic connection  $\nabla^W$  on W that preserves its Hermitian metric, i.e.,

$$\langle \nabla^W_X w_1, w_2 \rangle_W + \langle w_1, \nabla^W_X w_2 \rangle = X(\langle w_1, w_2 \rangle_W) \quad \forall w_j \in C^\infty(M, W) \quad \forall X \in C^\infty(M, TM).$$

The connection  $\nabla^W$  in Proposition 8.11.16 is called the *Chern connection* of W.

Let  $\Omega$  be the imaginary part of the Hermitian metric h of TM. This is a 2-form on M since, for all  $X, Y \in C^{\infty}(M, TM)$ ,

$$\Omega(Y,X) = \Im h(Y,X) = \Im \overline{h(X,Y)} = -\Im h(X,Y) = -\Omega(X,Y).$$

We then say that M is a Kähler manifold when  $\Omega$  is a closed form, i.e.,  $d\Omega = 0$ .

Recall that by assumption the Riemannian metric of M is the real part of the Hermitian metric h.

PROPOSITION 8.11.17 (See [**BGV**]). The manifold M is Kähler if and only if the Levi-Civita connection  $\nabla^{TM}$  on TM is a holomorphic connection, i.e., it agrees with the Chern connection of TM.

From now on, we assume that M is a compact Kähler manifold. Then  $\nabla^{TM}$  preserves the complex structure, and hence it can be lifted to a Hermitian connection on  $\Lambda^{0,q}T^*M$  for each  $q=0,1,\ldots,n$ .

In addition, in the very same way as we constructed the spinor representation, we can endow the bundle  $\Lambda^{0,*}T^*M$  with a Clifford-module bundle structure as follows.

The  $\mathbb{Z}_2$ -grading of  $\Lambda^{0,*}T^*M$  is simply given by the splitting,

$$\Lambda^{0,*}T^*M = \Lambda^{0,\text{ev}}T^*M \oplus \Lambda^{0,\text{odd}}T^*M.$$

The action of  $Cl_{\mathbb{C}}(M)$  is such that, above all  $x \in M$ ,

$$c(\xi).\omega = \sqrt{2} \left( \varepsilon(\xi^{1,0}) - i(\xi^{1,0}) \right) \omega \qquad \forall \xi \in T_{\mathbb{C}}^* M \quad \forall \omega \in \Lambda^{0,*} T^* M,$$

where  $\varepsilon(\xi^{1,0})$  is the exterior by the component  $\xi^{0,1}$  of  $\xi$  in  $\Lambda^{0,1}T_x^*M$  and  $\iota(\xi^{1,0})$  is the interior product by its component  $\xi^{1,0}$  in  $\Lambda^{1,0}T_x^*M$ .

This allows us to realize the spinor bundle of M as the bundle  $\Lambda^{0,*}T^*M$  equipped with the above Clifford-module structure. We then endow the bundle  $(\Lambda^{0,*}T^*M)\otimes W$  with its twisted Clifford-module bundle structure. Namely, its  $Z_2$ -grading is given by

$$(\Lambda^{0,*}T^*M)\otimes W=\bigg((\Lambda^{0,\operatorname{ev}}T^*M)\otimes W\bigg)\oplus\bigg((\Lambda^{0,*}T^*M)\otimes W\bigg),$$

and the action of  $Cl_{\mathbb{C}}(M)$  is such that, for all  $x \in M$ ,

$$a.(\omega \otimes w) = (a.\omega) \otimes w \qquad \forall a \in \operatorname{Cl}_x(M) \quad \forall (\omega, w) \in \Lambda^{0,*}T_x^*M \times W_x.$$

In addition, we equipped  $(\Lambda^{0,*}T^*M)\otimes W$  with the twisted connection,

$$\nabla^{(\Lambda^{0,*}T^*M)\otimes W} := \nabla^{\Lambda^{0,*}T^*M} \otimes 1 + 1 \otimes \nabla^W,$$

where  $\nabla^W$  is the Chern connection of W.

LEMMA 8.11.18 (See [**BGV**]). If M is Kähler, then  $\nabla^{(\Lambda^{0,*}T^*M)\otimes W}$  is a Clifford connection on the Clifford-module bundle  $(\Lambda^{0,*}T^*M)\otimes W$  and its associated Dirac operator is equal to  $\sqrt{2}(\overline{\partial}_W + \overline{\partial}_W^*)$ .

In the sequel, we denote by  $R^{T^{1,0}M}$  the curvature of the Levi-Civita connection of  $T^{1,0}M$ .

LEMMA 8.11.19 (See [**BGV**]).

(1) The twisted curvature of  $\nabla^{(\Lambda^{0,*}T^*M)\otimes W}$  is given by

(8.10) 
$$F^{(\Lambda^{0,*}T^*M)\otimes W/\$} = \frac{1}{2}\operatorname{Tr}\left[R^{T^{1,0}M}\right] + 1\otimes F^W.$$

(2) We have

(8.11) 
$$\hat{A}(R^M) = \det \left[ \frac{R^{T^{1,0}M}}{e^{\frac{1}{2}R^{T^{1,0}M}} - e^{-\frac{1}{2}R^{T^{1,0}M}}} \right].$$

It follows from (8.10) that the relative Chern form of  $F^{(\Lambda^{0,*}T^*M)\otimes W/\$}$  equals

$$\exp\left(-\frac{1}{2}\operatorname{Tr}\left[R^{T^{1,0}M}\right]\right)\wedge\operatorname{Ch}(F^W)=\det\left[e^{-\frac{1}{2}R^{T^{1,0}M}}\right]\wedge\operatorname{Ch}(F^W).$$

Combining this with (8.11) shows that  $\hat{A}(R^M) \wedge \text{Ch}(F^{(\Lambda^{0,*}T^*M)\otimes W/\$})$  is equal to

$$\det \left[ \frac{R^{T^{1,0}M}}{e^{\frac{1}{2}R^{T^{1,0}M}} - e^{-\frac{1}{2}R^{T^{1,0}M}}} \right] \wedge \det \left[ e^{-\frac{1}{2}R^{T^{1,0}M}} \right] \wedge \operatorname{Ch}(F^W) = \operatorname{Td}(R^{T^{1,0}M}) \wedge \operatorname{Ch}(F^W),$$

where  $\operatorname{Td}(R^{T^{1,0}M})$  is the Todd form of  $R^{T^{1,0}M}$ , i.e.,

$$\mathrm{Td}(R^{T^{1,0}M}) = \det \left[ \frac{R^{T^{1,0}M}}{e^{R^{T^{1,0}M}} - 1} \right].$$

Therefore, by applying the local index formula for Atiyah-Singer and using (8.9) we obtain

THEOREM 8.11.20 (Hirzebruch-Riemann-Roch). If M is Kähler, then

$$\chi(M,W) = \operatorname{ind}(\overline{\partial}_W + \overline{\partial}_W^*) = (2i\pi)^{-\frac{n}{2}} \int_M \left[ \operatorname{Td}(R^{T^{1,0}M}) \wedge \operatorname{Ch}(F^W) \right]^{(n)}.$$

#### 8.12. Proof of the Local Index Theorem

In this section we reproduce the proof of the local index theorem (i.e., Theorem 8.10.5) given in [Po]. As mentioned in Section 8.10, this theorem yields the local index formula of Atiyah-Singer. The proof here is given for twisted Dirac operators on spin manifolds, but the argument can be extended to more general Dirac operators on Clifford-module bundle coming from a Clifford connection.

The argument is based on combining the rescaling of Getzler [Ge2] with the approach to the heat kernel asymptotics of Greiner [Gr].

**8.12.1.** Greiner's approach of the heat kernel asymptotics. In this section we recall Greiner's approach of the heat kernel asymptotics as in [Gr] and [BGS].

Let E be a Hermitian vector bundle over M and let  $\Delta: C^{\infty}(M, E) \to C^{\infty}(M, E)$  be a selfadjoint elliptic second order differential operator with positive principal symbol.

If we regard  $\Delta$  as an unbounded operator on  $L^2(M, E)$  with domain the Sobolev space  $L^2(M, E)$ , then  $\Delta$  is selfadjoint and bounded from below. Therefore, by standard Borel functional calculus we can define the heat semi-group  $e^{-t\Delta}$ ,  $t \geq 0$ , as a family of selfadjoint bounded operators on  $L^2(M, E)$ .

Furthermore, the ellipticity of  $\Delta$  implies that, for all t > 0, the operator  $e^{-t\Delta}$  is smoothing. That is, its Schwartz kernel  $k_t(x,y)$  is contained in  $C^{\infty}(M,E) \hat{\otimes} C^{\infty}(M,E^* \otimes |\Lambda|(M))$ , where  $|\Lambda|(M)$  is the density bundle of M. The kernel  $k_t(x,y)$  is called the heat kernel of  $\Delta$ .

Recall that the heat semigroup allows us to invert the heat equation. Namely, consider the operator  $Q_0: C_c^{\infty}(M \times \mathbb{R}, E) \to \mathcal{D}'(M \times \mathbb{R}, E)$  defined by

(8.12) 
$$Q_0 u(x,s) := \int_0^\infty e^{-s\Delta} u(x,t-s) dt \qquad \forall u \in C_c^\infty(M \times \mathbb{R}, E).$$

Then  $Q_0$  maps continuously  $C_c^{\infty}(M \times \mathbb{R}, E)$  to  $C^0(\mathbb{R}, L^2(M, E)) \subset \mathcal{D}'(M \times \mathbb{R}, E)$  and satisfies

$$(\Delta + \partial_t)Q_0 u = Q_0(\Delta + \partial_t)u = u \qquad \forall u \in C_c^{\infty}(M \times \mathbb{R}, E).$$

Notice that the operator  $Q_0$  has the *Volterra property* in the sense of  $[\mathbf{Pi}]$ , i.e., it has a Schwartz kernel of the form  $K_{Q_0}(x, y, t - s)$ , where  $K_{Q_0}(x, y, t)$  vanishes on the region t < 0. In fact,

$$K_{Q_0}(x,y,t) = \begin{cases} k_t(x,y) & \text{if } t > 0, \\ 0 & \text{if } t < 0. \end{cases}$$

These equalities are the main motivation for using pseudodifferential techniques to study the heat kernel  $k_t(x, y)$ . The idea is to consider a class of  $\Psi$ DOs, the Volterra  $\Psi$ DOs([**Gr**], [**Pi**], [**BGS**]), taking into account

- (i) The aforementioned Volterra property.
- (ii) The parabolic homogeneity of the heat operator  $\Delta + \partial_t$ , i.e., the homogeneity with respect to the dilations  $\lambda.(\xi,\tau) = (\lambda \xi, \lambda^2 \tau), \ (\xi,\tau) \in \mathbb{R}^{n+1}, \lambda \neq 0.$

In the sequel, for  $g \in \mathcal{S}'(\mathbb{R}^{n+1})$  and  $\lambda \neq 0$ , we denote by  $g_{\lambda}$  the element of  $\mathcal{S}'(\mathbb{R}^{n+1})$  defined by

$$(8.13) \qquad \langle g_{\lambda}(\xi,\tau), u(\xi,\tau) \rangle := |\lambda|^{-(n+2)} \langle g(\xi,\tau), u(\lambda^{-1}\xi, \lambda^{-2}\tau) \rangle \quad \forall u \in \mathcal{S}(\mathbb{R}^{n+1}).$$

DEFINITION 8.12.1. A distribution  $g \in \mathcal{S}'(\mathbb{R}^{n+1})$  is (parabolic) homogeneous of degree  $m, m \in \mathbb{Z}$ , when

$$g_{\lambda} = \lambda^m g \qquad \forall \lambda \in \mathbb{R} \setminus 0.$$

In the sequel, we denote by  $\mathbb{C}_-$  the complex halfplane  $\{\Im \tau > 0\}$  with closure  $\overline{\mathbb{C}_-}$ . Then:

LEMMA 8.12.2 ([BGS, Prop. 1.9]). Let  $q(\xi, \tau) \in C^{\infty}((\mathbb{R}^n \times \mathbb{R}) \setminus 0)$  be a parabolic homogeneous symbol of degree m such that:

- (i) q(x, ξ, τ) extends to a continuous function on (R<sup>n</sup> × C̄\_) \ 0 in such way to be holomorphic in the last variable when the latter is restricted to C̄\_.
   Then there is a unique q ∈ S'(R<sup>n+1</sup>) agreeing with q on R<sup>n+1</sup> \ 0 so that:
  - (ii) g is homogeneous of degree m.
  - (iii) The inverse Fourier transform  $\check{g}(x,t)$  vanishes for t < 0.

Remark 8.12.3. If  $m \le -(n+2)$ , then (ii) need not hold for symbols that do not satisfy (i).

Let U be an open subset of  $\mathbb{R}^n$ . We define Volterra symbols and Volterra  $\Psi$ DOson  $U \times \mathbb{R}^{n+1} \setminus 0$  as follows.

DEFINITION 8.12.4.  $S_{\mathbf{v}}^m(U \times \mathbb{R}^{n+1}), m \in \mathbb{Z}$ , consists of smooth functions  $q(x, \xi, \tau)$  on  $U \times \mathbb{R}^n \times \mathbb{R}$  with an asymptotic expansion  $q(x, \xi, \tau) \sim \sum_{j \geq 0} q_{m-j}(x, \xi, \tau)$ , where

-  $q_l(x,\xi,\tau) \in C^{\infty}(U \times [(\mathbb{R}^n \times \mathbb{R}) \setminus 0])$  is a homogeneous Volterra symbol of degree l, i.e.  $q_l$  is parabolic homogeneous of degree l and satisfies the property (i) in Lemma 8.12.2 with respect to the last n+1 variables.

- The sign  $\sim$  means that, for all compacts  $K \subset U$ , integers N and k and multi-orders  $\alpha$  and  $\beta$ , there is a constant  $C_{NK\alpha\beta k} > 0$  such that

$$(8.14) \quad |\partial_x^{\alpha} \partial_{\xi}^{\beta} \partial_{\tau}^{k} (q - \sum_{j \le N} q_{m-j})(x, \xi, \tau)| \le C_{NK\alpha\beta k} (|\xi| + |\tau|^{1/2})^{m-N-|\beta|-2k},$$

for all 
$$(x, \xi, \tau) \in K \times \mathbb{R}^n \times \mathbb{R}$$
 with  $|\xi| + |\tau|^{\frac{1}{2}} \ge 1$ .

In the sequel, for any symbol  $q(x, \xi, \tau) \in S_{\mathbf{v}}^m(U \times \mathbb{R})$  we denote by  $q(x, D_x, D_t)$  the operator from  $C_c^{\infty}(U \times \mathbb{R})$  to  $C^{\infty}(U \times \mathbb{R})$  defined by

$$q(x, D_x, D_t)u(x, t) := (2\pi)^{-(n+1)} \iint e^{i(x \cdot \xi + t\tau)} q(x, \xi, \tau) \hat{u}(\xi, \tau) d\xi d\tau \quad \forall u \in C_c^{\infty}(U \times \mathbb{R}).$$

DEFINITION 8.12.5.  $\Psi^m_{\mathbf{v}}(U \times \mathbb{R}), m \in \mathbb{Z}$ , consists of continuous linear operators  $Q: C_c^{\infty}(U_x \times \mathbb{R}_t) \to C^{\infty}(U_x \times \mathbb{R}_t)$  such that

- (i) Q has the Volterra property.
- (ii) Q can be put in the form,

$$Q = q(x, D_x, D_t) + R,$$

for some symbol  $q(x, \xi, \tau) \in S_{\mathbf{v}}^m(U \times \mathbb{R})$  and some smoothing operator R.

If Q is a Volterra  $\Psi$ DO we shall denote by  $K_Q(x, y, t-s)$  its distribution kernel, so that the distribution  $K_Q(x, y, t)$  vanishes for t < 0.

EXAMPLE 8.12.6. Let P be a differential operator of order 2 on U and let  $p_2(x,\xi)$  denote the principal symbol of P. Then the heat operator  $P + \partial_t$  is a Volterra  $\Psi$ DO of order 2 with principal symbol  $p_2(x,\xi) + i\tau$ .

Other examples of Volterra  $\Psi$ DOs are given by the homogeneous operators defined below.

DEFINITION 8.12.7. Let  $q_m(x,\xi,\tau) \in C^{\infty}(U \times (\mathbb{R}^{n+1} \setminus 0))$  be a homogeneous Volterra symbol of order m and let  $g_m \in C^{\infty}(U) \hat{\otimes} \mathcal{S}'(\mathbb{R}^{n+1})$  denote its unique homogeneous extension given by Lemma 8.12.2. Then:

- $\check{q}_m(x,y,t)$  is the inverse Fourier transform of  $g_m(x,\xi,\tau)$  in the last n+1 variables.
- $q_m(x, D_x, D_t)$  is the operator with kernel  $\check{q}_m(x, y x, t)$ .

PROPOSITION 8.12.8 ([Gr], [Pi], [BGS]). The following hold.

- (1) Composition. Let  $Q_j \in \Psi_{\mathbf{v}}^{m_j}(U \times \mathbb{R})$ , j = 1, 2, have symbol  $q_j$  and suppose that  $Q_1$  or  $Q_2$  is properly supported. Then  $Q_1Q_2$  lies in  $\Psi_{\mathbf{v}}^{m_1+m_2}(U \times \mathbb{R})$  and has symbol  $q_1 \# q_2 \sim \sum \frac{1}{\alpha!} \partial_{\xi}^{\alpha} q_1 D_{\xi}^{\alpha} q_2$ .
- (2) Parametrices. An operator  $Q \in \Psi^m_{\mathbf{v}}(U \times \mathbb{R})$  admits a parametrix in  $\Psi^{-m}_{\mathbf{v}}(U \times \mathbb{R})$  if, and only if, its principal symbol is nowhere vanishing on  $U \times [(\mathbb{R}^n \times \overline{\mathbb{C}_-} \setminus 0)]$ .
- (3) Invariance. Let  $\phi: U \to V$  be a diffeomorphism onto another open subset V of  $\mathbb{R}^n$  and let Q be a Volterra  $\Psi$  on  $U \times \mathbb{R}$  of order m. Then  $Q = (\phi \oplus \mathrm{id}_{\mathbb{R}})_* Q$  is a Volterra  $\Psi$  on  $V \times \mathbb{R}$  of order m.

In addition, the following property shows the relevance of Volterra  $\Psi DOs$  for deriving small times asymptotics.

LEMMA 8.12.9 ([**Gr**, Chap. I], [**BGS**, Thm. 4.5]). Let  $Q \in \Psi_{\mathbf{v}}^m(U \times \mathbb{R})$  have symbol  $q \sim \sum q_{m-j}$ . Then ,in  $C^{\infty}(U)$ ,

(8.15) 
$$K_Q(x,y,t) \sim t^{-(\frac{n}{2} + [\frac{m}{2}] + 1)} \sum_{l \ge 0} t^l \check{q}_{2[\frac{m}{2}] - 2l}(x,0,1) \quad as \ t \to 0^+,$$

where the notation  $\check{q}_k$  has the same meaning as in Definition 8.12.7.

PROOF. As the Fourier transform relates the decay at infinity to the behavior at the origin of the Fourier transform the distribution  $\check{q} - \sum_{j \leq J} \check{q}_{m-j}$  lies in  $C^N(U_x \times \mathbb{R}^n_y \times \mathbb{R}_t)$  as soon as J is large enough. Since  $Q - q(x, D_x, D_t)$  is smoothing it follows that  $R_J(x,t) := K_Q(x,x,t) - \sum_{j \leq J} \check{q}_{m-j}(x,0,t)$  is of class  $C^N$ . As  $R_J(x,y,t) = 0$  for t < 0 we see that

$$\partial_t^l R_J(x,0) = 0$$
 for  $l = 0, 1, \dots, N$ .

Thus  $R_J(.,t) = O(t^N)$  in  $C^N(U)$  as  $t \to 0^+$ . This proves that, in  $C^\infty(U)$ ,

(8.16) 
$$K_Q(x, x, t) \sim \sum \check{q}_{m-j}(x, 0, t) \text{ as } t \to 0^+.$$

Let  $j \in \mathbb{N}_0$ . Observe that, for all  $\lambda \neq 0$ ,

$$(\check{q}_{m-j})_{\lambda} = |\lambda|^{-(n+2)} (q_{m-j,\lambda^{-1}})^{\vee} = |\lambda|^{-(n+2)} \lambda^{j-m} \check{q}_{m-j}.$$

Thus, setting  $\lambda = \sqrt{t}$  with t > 0, we get

$$\check{q}_{m-j}(x,0,t) = t^{\frac{j-n-m}{2}-1} \check{q}_{m-j}(x,0,1).$$

Furthermore, if we take  $\lambda = -1$  and m - j is odd, then

$$\check{q}_{m-j}(x,0,1) = (-1)^{m-j} q_{m-j}(x,0,1) = -q_{m-j}(x,0,1) = 0.$$

Combining all this with (8.16) shows that, in  $C^{\infty}(U)$ ,

(8.17) 
$$K_Q(x,x,t) \sim_{t\to 0^+} \sum_{m-j \text{ even}} t^{\frac{j-n-m}{2}-1} \check{q}_{m-j}(x,0,1),$$

that is, (8.15) holds. The lemma is thus proved.

The invariance property in Proposition 8.12.8 allows us to define Volterra  $\Psi$ DOs on  $M \times \mathbb{R}$  acting on the sections of the vector bundle E. Then all the preceding properties hold *verbatim* in this context. In particular, the heat operator  $\Delta + \partial_t$  has a parametrix Q in  $\Psi_v^{-2}(M, \times \mathbb{R}, E)$ .

In fact, comparing the operator (8.12) with any Volterra parametrix for  $\Delta + \partial_t$  allows us to prove

THEOREM 8.12.10 ([Gr], [Pi], [BGS, pp. 363-362]). The differential operator  $\Delta + \partial_t$  is invertible and its inverse  $(\Delta + \partial_t)^{-1}$  is a Volterra  $\Psi DO$  of order -2.

Combining this with Lemma 8.12.9 gives the heat kernel asymptotics below.

THEOREM 8.12.11 ([Gr, Thm. 1.6.1]). In  $C^{\infty}(M, |\Lambda|(M) \otimes \operatorname{End} E)$  we have

(8.18) 
$$k_t(x,x) \sim_{t\to 0^+} t^{\frac{-n}{2}} \sum_{l\geq 0} t^l a_l(\Delta)(x), \qquad a_l(\Delta)(x) = \check{q}_{-2-2l}(x,0,1),$$

where the equality on the right-hand side shows how to compute the densities  $a_l(\Delta)(x)$  in local trivializing coordinates by means of the symbol  $q(x, \xi, \tau) \sim \sum q_{-2-j}(x, \xi, \tau)$  of any Volterra parametrix for  $\Delta + \partial_t$ .

This approach to the heat kernel asymptotics present several advantages. First, as Theorem 8.12.11 is a purely local statement we can easily localize the heat kernel asymptotics. In fact, given a Volterra parametrix Q for  $\Delta + \partial_t$  in some local trivializing coordinates around  $x_0 \in M$ , comparing the asymptotics (8.15) and (8.18) we get

(8.19) 
$$k_t(x_0, x_0) = K_Q(x_0, x_0, t) + O(t^{\infty}) \quad \text{as } t \to 0^+.$$

Therefore in order to determine the heat kernel asymptotics (8.18) at  $x_0$  we only need a Volterra parametrix for  $\Delta + \partial_t$  near  $x_0$ .

Second, we have a genuine asymptotics with respect to the  $C^{\infty}$ -topology and it can be differentiated as follows.

PROPOSITION 8.12.12. Let  $P: C^{\infty}(M, E) \to C^{\infty}(M, E)$  be a differential operator of order m and let us denote by  $h_t(x,y)$  the Schwartz kernel of  $Pe^{-t\Delta}$ . Then, in  $C^{\infty}(M, |\Lambda| \otimes \operatorname{End} E)$ ,

(8.20) 
$$h_t(x,x) \sim_{t\to 0^+} t^{\left[\frac{m}{2}\right] - \frac{n}{2}} \sum_{l>0} t^l b_l(x), \qquad b_l(x) = \check{r}_{2\left[\frac{m}{2}\right] - 2 - 2l}(x,0,1),$$

where the equality on the right-hand side gives a formula for computing the densities  $b_l(x)$  in local trivializing coordinates by means of the symbol  $r \sim \sum r_{m-2-j}$  of  $R = P(\Delta + \partial_t)^{-1}$  (or of R = PQ, where Q is any Volterra parametrix for  $\Delta + \partial_t$ ).

PROOF. Observe that

$$h_t(x,y) = P_x k_t(x,y) = P_x K_{(\Delta+\partial_t)^{-1}}(x,y,t) = K_{P(\Delta+\partial_t)^{-1}}(x,y,t).$$

Therefore, the result follows by applying Lemma 8.12.9 to  $P(\Delta + \partial_t)^{-1}$  (or to PQ, where Q is any Volterra parametrix for  $\Delta + \partial_t$ ).

Finally, in local trivializing coordinates the densities  $a_j(\Delta)(x)$ 's can be explicitly computed in terms of the symbol  $p=p_2+p_1+p_0$  of  $\Delta$ . To see this let  $q\sim\sum q_{-2-j}$  be the symbol of a Volterra parametrix Q for  $\Delta+\partial_t$ . As  $q\#p\sim q(p+i\tau)+\sum\frac{1}{\alpha!}\partial_\xi^\alpha qD_x^\alpha p\sim 1$  we get  $q_{-2}=(p_2+i\tau)^{-1}$  and

(8.21) 
$$q_{-2-j} = -\left(\sum_{k+l+|\alpha|=j} \frac{1}{\alpha!} \partial_{\xi}^{\alpha} q_{-2-k} D_x^{\alpha} p_{2-l}\right) (p_2 + i\tau)^{-1}, \quad j \ge 1.$$

Therefore, combining with (8.18) we deduce that, as in [**Gi**], the densities  $a_j(\Delta)(x)$ 's are universal polynomials in the pets at  $x_0$  of the symbol of  $\Delta$  with coefficients depending smoothly on its principal symbol.

Similarly, in local trivializing coordinates the densities  $b_l(x)$ 's in (8.20) can be expressed universal polynomials in the the jets at  $x_0$  of the symbols of  $\Delta$  and P with coefficients depending smoothly on the principal symbol of  $\Delta$ .

**8.12.2.** Proof of the Local Index Theorem. In this subsection, we shall prove the local index theorem. We shall use the same notation as in Section 8.10. In particular, E is a Clifford-module bundle with Clifford connection  $\nabla^E$  and associated Dirac operator  $D_E$ . In addition, we denote by  $k_t(x,y)$  the heat kernel of  $D_E^2$ .

It follows from Theorem 8.12.11 that, in  $C^{\infty}(M, |\Lambda|(M))$ ,

$$\operatorname{Str}_{E}[k_{t}(x,x)] \sim t^{-\frac{n}{2}} \sum_{j \ge 0} t^{j} \operatorname{Str}_{E}[a_{j}(D_{E}^{2})(x)]$$
 as  $t \to 0^{+}$ .

We thus have an asymptotics in  $C^{\infty}(M, |\Lambda|(M))$ . Therefore, in order to prove the local index theorem it is enough to show that, near any given point  $x_0$ , in some local coordinates centered at  $x_0$  and local trivialization of E near  $x_0E$ ,

(8.22) 
$$\operatorname{Str}_{E}[k_{t}(0,0)] = (2i\pi)^{-\frac{n}{2}} \left[ \hat{A}(R^{M}) \wedge \operatorname{Ch}(F^{E/\sharp}) \right]^{(n)} (0) + \operatorname{O}(t).$$

Furthermore, if follows from (8.19) that in order to so we may replace the heat kernel  $k_t(x,x)$  by the kernel  $K_Q(x,y,t)$  of any Volterra- $\Psi$ DO parametrix for  $D_E^2 + \partial_t$  defined near x=0. Indeed, we then have

(8.23) 
$$k_t(0,0) = K_O(0,0,t) + O(t^{\infty})$$
 as  $t \to 0^+$ .

This enables use to reduce to the case where M is a neighborhood of the origin in  $\mathbb{R}^n$  and E is a trivial vector bundle.

As we have total freedom on the choice of the local coordinates and the trivializations of E, we may choose to use local Riemannian coordinates and introduce a synchronous orthonormal tangent frame  $\{e_1,\cdots,e_n\}$  such that  $e_j=\partial_j$ . Using this frame to construct the spinor representation, near x=0 we may realize E as a trivial twisted Clifford bundle with fiber  $\S_n \otimes W$ , where  $\S_n$  is the spinor space of  $\mathbb{R}^n$  and W is a  $\mathbb{Z}_2$ -graded vector space. Notice that the Clifford connection  $\nabla^E$  then is a twisted Clifford connection  $\nabla^{\S_n} \otimes 1 + 1 \otimes \nabla^W$ , where  $\nabla^{\S_n}$  is the spin connection on the trivial vector bundle  $M \times \S_n$  and  $\nabla^W$  is a Hermitian connection on  $M \times W$ . Incidentally, if we denote by  $F^W$  the curvature of  $\nabla^W$ , then the twisted curvature  $F^{E/\$}$  and its relative Chern form are given by

$$F^{E/\$} = 1 \otimes F^W$$
 and  $\operatorname{Ch}(F^{E/\$}) = \operatorname{Ch}(F^W)$ .

In addition, as we are using normal coordinates, near x=0 the coefficients  $g_{ij}=g(\partial_i,\partial_j)$  of the metric g are such that

(8.24) 
$$g_{ij}(x) = \delta_{ij} + O(|x|^2), \qquad \omega_{ikl}(x) = -\frac{1}{2} R^M_{ijkl}(0) x^j + O(|x|^2),$$

A proof of this asymptotics can be found for instance in  $[\mathbf{BGV}]$ .

Let us also introduce the coefficients of the Levi-Civita connection  $\nabla^{TM}$  and the curvature tensor  $R^M$  defined in terms of the synchronous orthonormal tangent frame  $\{e_i\}$  by

$$\omega_{ikl} = \langle \nabla_i^{TM} e_k, e_l \rangle \quad \text{and} \quad R_{ijkl}^M = \langle R^M(\partial_i, \partial_j) e_k, e_l \rangle.$$

Then, near x = 0,

(8.25) 
$$\omega_{ikl}(x) = -\frac{1}{2} R_{ijkl}^{M}(0) x^{j} + O(|x|^{2}).$$

See, e.g., [BGV] for a proof of this asymptotics.

Recall that the quantification and symbol maps are linear isomorphisms,

(8.26) 
$$\Lambda_{\mathbb{C}}^* \mathbb{R}^n \xrightarrow{c} \operatorname{Cl}_{\mathbb{C}}(\mathbb{R}^n) \text{ and } \operatorname{Cl}_{\mathbb{C}}(\mathbb{R}^n) \xrightarrow{\sigma} \Lambda_{\mathbb{C}}^* \mathbb{R}^n.$$

As n is even the spinor representation  $\rho: \mathrm{Cl}_{\mathbb{C}}(\mathbb{R}^n) \to \mathrm{End}\,\mathcal{S}_n$  is an algebra isomorphism which allows us to identify  $\mathrm{Cl}_{\mathbb{C}}(\mathbb{R}^n)$  with  $\mathrm{End}\,\mathcal{S}_n$ . Bearing in mind this identification, we shall also denote by c and  $\sigma$  the linear isomorphisms,

$$(8.27) \Lambda_{\mathbb{C}}^* \mathbb{R}^n \stackrel{c}{\longrightarrow} \operatorname{End} \mathcal{S}_n \quad \text{and} \quad \operatorname{End} \mathcal{S}_n \stackrel{\sigma}{\longrightarrow} \Lambda_{\mathbb{C}}^* \mathbb{R}^n,$$

which are obtained by composing the linear isomorphisms (8.26) with  $\rho$  or its inverse.

As E is a trivial bundle with fiber  $\mathcal{S}_n \otimes W$ , we can regard the Volterra  $\Psi$ DOs on  $M \times \mathbb{R}$  acting on the sections of E as elements of  $\Psi^*_{\mathbf{v}}(M \times \mathbb{R}) \otimes (\operatorname{End} \mathcal{S}_n) \otimes (\operatorname{End} W)$ . Using the linear maps (8.27), we then get linear maps,

$$\Psi_{\mathbf{v}}^{*}(M \times \mathbb{R}) \otimes (\Lambda_{\mathbb{C}}^{*}\mathbb{R}^{n}) \otimes (\operatorname{End} W) \xrightarrow{c} \Psi_{\mathbf{v}}^{*}(M \times \mathbb{R}, E),$$

$$\Psi_{\mathbf{v}}^{*}(M \times \mathbb{R}, E) \xrightarrow{\sigma} \Psi_{\mathbf{v}}^{*}(M \times \mathbb{R}) \otimes (\Lambda_{\mathbb{C}}^{*}\mathbb{R}^{n}) \otimes (\operatorname{End} W).$$

We get similar linear maps at the level of symbols and Schwartz kernels.

As E is the trivial  $\mathbb{Z}_2$ -graded bundle with fiber  $\mathcal{S}_n \otimes W$ , its supertrace is just  $\operatorname{Str}_{\mathcal{S}_n} \otimes \operatorname{Str}_W$ , where  $\operatorname{Str}_{\mathcal{S}_n}$  is the supertrace on the trivial bundle  $M \times \mathcal{S}_n$  and  $\operatorname{Str}_W$  is the supertrace on W. Notice that, as the metric on M varies on the fibers of TM, so does the supertrace on the fibers of  $M \times \operatorname{End}(\mathcal{S}_n)$ . However, thanks (8.24) at x=0 the metric g(x) agrees with the standard Euclidean metric, and hence we may use Proposition 4.8 to get

$$\operatorname{Str}_{\mathscr{S}_{n}}[T(0)] = (-2i)^{\frac{n}{2}} \sigma[T(0)]^{(n)} \qquad \forall T \in C^{\infty}(M, \operatorname{End} \mathscr{S}_{n}),$$

where  $\sigma[T(0)]^{(n)}$  is the *n*-the degree component of  $\sigma[T(0)] \in \Lambda_{\mathbb{C}}^* \mathbb{R}^n$ .

It follows from all this that, for all  $P \in \Psi^*_{\mathbf{v}}(M \times \mathbb{R}, E)$ ,

$$\operatorname{Str}_{E}[K_{P}(0,0,t)] = (-2i)^{\frac{n}{2}} (\sigma \otimes \operatorname{Str}_{W})[K_{P}(0,0,t)]^{(n)} \quad \forall t > 0.$$

Combining this with (8.23) then shows that

$$(8.28) k_t(0,0) = (-2i)^{\frac{n}{2}} (\sigma \otimes \operatorname{Str}_W) [K_P(0,0,t)]^{(n)} + O(t^{\infty}) \text{as } t \to 0^+.$$

Therefore, the proof of (8.22) boils down to showing the existence of a small-time limit of  $(\sigma \otimes \operatorname{Str}_W)[K_P(0,0,t)]^{(n)}$  and identifying it.

In order to study the small-time behavior of  $(\sigma \otimes \operatorname{Str}_W)[K_P(0,0,t)]^{(n)}$  we shall implement the rescaling of Getzler [**Ge2**] as a filtration on  $\Psi^*_{\mathbf{v}}(M \times \mathbb{R}, E)$ . Roughly speaking this rescaling aims at assigning the following degrees:

(8.29) 
$$\deg \partial_j = \frac{1}{2} \deg \partial_t = \deg c(dx^j) = -\deg x^j = 1,$$

while deg B=0 for any  $B\in \operatorname{End} W$ . We then obtain a filtration on  $\Psi_{\mathbf{v}}^*(M\times\mathbb{R},E)$  as follows.

Let  $Q \in \Psi^*_{\mathbf{v}}(M \times \mathbb{R}, E)$  have symbol  $q(x, \xi, \tau) \sim \sum_{k \leq m'} q_k(x, \xi, \tau)$ . Then taking components in each subspace  $\Lambda^j T^*_{\mathbb{C}} \mathbb{R}^n(n)$  and then using Taylor expansions at x = 0 gives formal expansions

(8.30) 
$$\sigma[q(x,\xi,\tau)] \sim \sum_{j,k} \sigma[q_k(x,\xi,\tau)]^{(j)} \sim \sum_{j,k,\alpha} \frac{x^{\alpha}}{\alpha!} \sigma[\partial_x^{\alpha} q_k(0,\xi,\tau)]^{(j)}.$$

According to (8.29) the symbol  $\frac{x^{\alpha}}{\alpha!}\partial_x^{\alpha}\sigma[q_k(0,\xi,\tau)]^{(j)}$  is Getzler homogeneous of degree  $k+j-|\alpha|$ . Therefore, we can expand  $\sigma[q(x,\xi,\tau)]$  as

(8.31) 
$$\sigma[q(x,\xi,\tau)] \sim \sum_{j>0} q_{(m-j)}(x,\xi,\tau), \qquad q_{(m)} \neq 0,$$

where  $q_{(m-j)}$  is a Getzler homogeneous symbol of degree m-j.

Definition 8.12.13. Using (8.31) we make the following definitions:

- The integer m is the Getzler order of Q.
- The symbol  $q_{(m)}$  on  $\mathbb{R}^n \times \mathbb{R}^n \times \mathbb{R}$  is the principal Getzler homogeneous symbol of Q.

- The operator  $Q_{(m)} = q_{(m)}(x, D_x, D_t)$  on  $\mathbb{R}^n \times \mathbb{R}$  is the model operator of Q.

Remark 8.12.14. The model operator  $Q_{(m)}$  is well defined according to definition 8.12.7.

Remark 8.12.15. By construction we always have Getzler order  $\leq$  order + n, but this is not an equality in general.

Example 8.12.16. Let  $A=A_idx^i$  be the connection one-form on W. Then by (8.25) the covariant derivative  $\nabla^E_i=\partial_i+\frac{1}{4}\omega_{ikl}(x)c(e^k)c(e^l)+A_i$  on  $\mathcal{S}_n\otimes W$  has Getzler order 1 and model operator

(8.32) 
$$\nabla_{i(1)}^{E} = \partial_{i} - \frac{1}{4} R_{ij}^{M}(0) x^{j}, \qquad R_{ij}^{M}(0) := R_{ijkl}^{M}(0) dx^{k} \wedge dx^{l}.$$

The interest to introduce Getzler orders stems from the following.

LEMMA 8.12.17. Let  $Q \in \Psi_{v}^{*}(M \times \mathbb{R}, M)$  have Getzler order m and model operator  $Q_{(m)}$ . Then, as  $t \to 0^{+}$ ,

$$\sigma[K_Q(0,0,t)]^{(j)} = \begin{cases} O(t^{\frac{j-m-n-1}{2}}) & \text{if } m-j \text{ is odd,} \\ t^{\frac{j-m-n}{2}-1} K_{Q_{(m)}}(0,0,1)^{(j)} + O(t^{\frac{j-m-n}{2}}) & \text{if } m-j \text{ is even.} \end{cases}$$

In particular, if m = -2, then

(8.33) 
$$\sigma[K_Q(0,0,t)]^{(n)} = K_{Q_{(-2)}}(0,0,1)^{(n)} + O(t).$$

PROOF. Let  $q(x,\xi,\tau) \sim \sum q_k(x,\xi,\tau)$  be the symbol of Q and let  $q_{(m)}(x,\xi,\tau)$  be its principal Getzler-homogeneous symbol. By Lemma 8.12.9 we have

$$(8.34) \qquad \qquad \sigma[K_Q(0,0,t)]^{(j)} \sim_{t \to 0^+} \sum t^{-\frac{n+2+m-j}{2}} \sigma[\check{q}_k(0,0,1)]^{(j)},$$

and we know that  $\check{q}_k(0,0,1) = 0$  if k is odd. Moreover, the symbol  $\sigma[q_k(0,\xi,\tau)]^{(j)}$  is Getzler homogeneous of degree k+j, and so it must be zero if k+j>m since otherwise Q would not have Getzler order m. Therefore, (8.35)

$$\sigma[K_Q(0,0,t)]^{(j)} = \begin{cases} O(t^{\frac{j-m-n-1}{2}}) & \text{if } m-j \text{ is odd,} \\ t^{\frac{j-m-n}{2}-1}\sigma[\check{q}_{m-j}(0,0,1)]^{(j)} + O(t^{\frac{j-m-n}{2}}) & \text{if } m-j \text{ is even.} \end{cases}$$

In addition, the symbol  $\sigma[q_{(m)}(0,\xi,\tau)]^{(j)}$  is equal to

(8.36) 
$$\sum_{k+j-|\alpha|=m} \left( \frac{x^{\alpha}}{\alpha!} \partial_x^{\alpha} \sigma[q_k(0,\xi,\tau)]^{(j)} \right)_{x=0} = \sigma[q_{m-j}(0,\xi,\tau)]^{(j)}.$$

As  $K_{Q_{(m)}}(x, y, t) = (q_{(m)})^{\vee}(x, y, t)$ , we deduce that

$$\sigma[\check{q}_{m-j}(0,0,1)]^{(j)} = (q_{(m)})^{\vee}(0,0,1)^{(j)} = K_{Q_{(m)}}(0,0,1)^{(j)}.$$

Combining this with (8.35) proves the lemma.

In the sequel, we say that a symbol or a  $\Psi DO$  is  $O_G(m)$  if it has Getzler order  $\leq m$ .

LEMMA 8.12.18. For j=1,2 let  $Q_j \in \Psi^*_{\mathbf{v}}(M \times \mathbb{R}, E)$  have Getzler order  $m_j$  and model operator  $Q_{(m_j)}$ . In addition, assume that  $Q_1$  or  $Q_2$  is properly supported. Then

(8.37) 
$$Q_1Q_2 = c \left[ Q_{(m_1)}Q_{(m_2)} \right] + O_G(m_1 + m_2 - 1).$$

PROOF. Let  $q_j$  be the symbol of  $Q_j$  and let  $q_{(m_j)}$  be its principal Getzler homogeneous symbol. By Proposition 8.12.8 the operator  $Q_1Q_2$  has symbol  $q_1\#q_2$ .

Moreover, for N large enough, the symbol  $q_1 \# q_2 - \sum_{|\alpha| \leq N} \frac{1}{\alpha!} \partial_{\xi}^{\alpha} q_1 D_x^{\alpha} q_2$  has order  $< m_1 + m_2 - n$ , and hence has Getzler order  $< m_1 + m_2$ . As  $\partial_{\xi}^{\alpha} q_1 . D_x^{\alpha} q_2 - c[\partial_{\xi}^{\alpha} q_{(m_1)} \wedge D_x^{\alpha} f_{(m_2)}]$  has Getzler order  $\le m_1 + m_2 - |\alpha| - 1$  it follows that, for N large enough,

(8.38) 
$$q_1 \# q_2 = \sum_{|\alpha| < N} \frac{1}{\alpha!} c(\partial_{\xi}^{\alpha} q_{m_1} \wedge D_x^{\alpha} q_{m_2}) + \mathcal{O}_G(m_1 + m_2 - 1).$$

Observe that  $\sum \frac{1}{\alpha!} \partial_{\xi}^{\alpha} q_{(m_1)} \wedge D_x^{\alpha} q_{(m_2)}$  is exactly the symbol of  $Q_{(m_1)} Q_{(m_2)}$  since  $q_{(m_2)}(x,\xi,\tau)$  is polynomial in x and thus the sum is finite. Therefore, taking N large enough in (8.38) shows that the symbols of  $\sigma[Q_1Q_2]$  and  $Q_{(m_1)}Q_{(m_2)}$  coincide modulo a symbol of Getzler order  $\leq m_1 + m_2 - 1$ . This proves the lemma.  $\square$ 

Recall that by the Lichnerowicz formula,

$$D_E^2 = (\nabla^E)^* \nabla^E + \mathcal{F}^{E/\$} + \frac{1}{4} \kappa_M,$$

where

$$\mathcal{F}^{E/\$} = \frac{1}{2} \sum F^{E/\$}(e_i, e_j) c(e^i) c(e^j) = \frac{1}{2} \sum F^W(e_i, e_j) c(e^i) c(e^j).$$

Moreover, as explained in [BGV, p. 66],

$$(\nabla^E)^* \nabla^E = -g^{ij} (\nabla^E_i \nabla^E_j - \Gamma^k_{ij} \nabla_k),$$

where the  $g^{ij}$  are the coefficients of the inverse metric  $g^{-1}$  and the  $\Gamma^k_{ij}$  are the Christoffel symbols of g so that  $\nabla^{TM}_i \partial_k = \Gamma^l_{ik} \partial_l$ .

Combining this with Lemma 8.12.18 and the equations (8.24), (8.25) and (8.32) shows that  $D_E^2$  has Getzler order 2 and model operator,

(8.39) 
$$D_{E(2)}^{2} = -\delta_{ij} \nabla_{i(1)} \nabla_{j(1)} + \frac{1}{2} F^{W}(\partial_{i}, \partial_{j})(0) dx^{i} \wedge dx^{j}$$

$$= H_{R} + F^{W}(0), \qquad H_{R} := -\sum_{i=1}^{n} \left(\partial_{i} - \frac{1}{4} R_{ij}^{M}(0) x^{j}\right)^{2}.$$

LEMMA 8.12.19. Let Q be a Volterra parametrix for  $D_E^2 + \partial_t$ . Then

- (1) Q has Getzler order -2 and its model operator is  $(H_R + F^W(0) + \partial_t)^{-1}$ .
- (2) We have

(8.40) 
$$K_{(H_R+F^W(0)+\partial_t)^{-1}}(0,x,t) = G_R(x,t) \wedge e^{-tF^W(0)},$$

where  $G_R(x,t)$  is the fundamental solution of  $H_R + \partial_t$ , i.e., the solution of the equation,  $(H_R + F^W(0) + \partial_t)G_R(x,t) = \delta(x,t)$ , where  $\delta(x,t)$  is the Dirac function on  $\mathbb{R}^n \times \mathbb{R}$ .

(3)  $As t \rightarrow 0^+$ ,

(8.41) 
$$\sigma[K_Q(0,0,t)]^{(2j)} = t^{j-\frac{n}{2}} [G_R(0,1) \wedge e^{-F^W(0)}]^{(2j)} + O(t^{j-\frac{n}{2}+1}).$$

PROOF. Note that (3) follows by combining (1) and (2) with Lemma 8.12.17. Therefore, we only need to prove the first two assertions.

Let  $p(x,\xi) = \sum p_j(x,\xi)$  be the symbol of  $\mathbb{Z}^2$  and let  $q \sim \sum q_{-2-j}$  be the symbol of Q. As  $\mathbb{Z}^2$  is elliptic and has Getzler order 2 we have  $p_{(2)}(0,\xi)^{(0)} = p_2(0,\xi) \neq 0$ .

Hence  $q_{-2} = (p_2 + i\tau)^{-1}$  has Getzler order -2. It then follows from (8.21) that each symbol  $q_{-2-j}$  has Getzler order  $\leq -2$ , and hence Q has Getzler order -2.

Notice also that, as  $(p^2 + \partial_t)Q - 1$  is smoothing, it follows from Lemma 8.12.18 the operator  $(H_R + F^W(0) + \partial_t)Q_{(-2)} - 1$  has Getzler order  $\leq -1$ . As it is Getzler-homogeneous of degree 0 it must be zero, and hence  $Q_{(-2)} = (H_R + F^W(0) + \partial_t)^{-1}$ . It is not difficult to check that, at the level of Schwartz kernels, the equality  $Q_{(-2)}(H_R + F^W(0) + \partial_t) = 1$  means that

$$(H_{R,y} + F^W(0) + \partial_s)^T [K_{Q_{(-2)}}(x, y, t - s)] = \delta(x - y, t - s),$$

where  $(H_{R,y} + F^W(0) + \partial_s)^T = H_{R,y} + F^W(0) - \partial_s$  is the transpose of  $H_{R,y} + F^W(0) + \partial_s$ . This implies that

$$(H_{R,y} + F^{W}(0) + \partial_{s}) [K_{Q_{(-2)}}(0, y, s)] = \delta(y, s),$$

that is,  $K_{Q_{(-2)}}(0,x,t)$  is the fundamental solution of  $H_R + F^W(0) + \partial_t$ .

Observe that if we denote by  $G_R(x,t)$  be the fundamental solution of  $H_R + \partial_t$ , then  $G_{R,F}(x,t) := G_R(x,t) \wedge e^{-tF^W(0)}$  too is the fundamental solution of  $H_R + F^W(0) + \partial_t$ . Thus  $K_{Q_{(-2)}}(0,x,t) = G_R(x,t) \wedge e^{-tF^W(0)}$ . The proof is complete.  $\square$ 

At this stage observe that  $H_R$  is the harmonic oscillator associated to the skew-symmetric matrix  $R^M(0) = (R^M_{ij}(0))$ . We shall now make use of a version of the Melher's formula to determine the fundamental solution of  $H_R + \partial_t$  (compare [**BGV**]).

LEMMA 8.12.20 (Melher Formula). Let a > 0 and consider the harmonic oscillator  $H_a := -\frac{d}{dx^2} + \frac{1}{4}a^2x^2$  on  $\mathbb{R}$ . Then the fundamental solution of  $H_a + \partial_t$ 

$$G_a(x,t) := \chi(t)(4\pi t)^{-\frac{1}{2}} \left(\frac{at}{\sinh at}\right)^{\frac{1}{2}} \exp\left(-\frac{1}{4t}x^2 \frac{at}{\tanh at}\right),$$

where  $\chi(t)$  is the characteristic function of the interval  $(0,\infty)$ .

PROOF. For  $(x,t) \in \mathbb{R} \times (0,\infty)$  define

(8.42) 
$$S_a(x,t) := (4\pi t)^{-\frac{1}{2}} \left(\frac{at}{\sinh at}\right)^{\frac{1}{2}} \exp\left(-\frac{1}{4t}x^2 \frac{at}{\tanh at}\right), \quad t > 0.$$

Observe that  $S_a(x,t)$  is integrable on all products  $\mathbb{R} \times (0,c)$  with c>0 and

$$(H_a + \partial_t)S_a(x,t) = 0$$
  $\forall (x,t) \in \mathbb{R} \times (0,\infty).$ 

Notice further that, as  $t \to 0^+$ ,

$$\hat{S}_{x \to \xi}(\xi, t) = \cosh^{-\frac{1}{2}}(at) \exp(-\xi^2 t \frac{\tanh at}{at}) \longrightarrow 1,$$

uniformly on compact sets of  $\mathbb{R}$ . Thus, as  $t \to 0^+$ ,

$$S(x,t) \longrightarrow \delta(x)$$
 in  $S'(\mathbb{R})$ .

It follows from this that, for all u and v in  $C_c^{\infty}(\mathbb{R})$ ,

$$\langle (H_a + \partial_t) G_a(x, t), u(x)v(t) \rangle = \int_0^\infty \int_{-\infty}^\infty S_a(x, t) (H_a - \partial_t) (u(x)v(t)) dx dt$$
$$= \int_0^\infty \int_{-\infty}^\infty ((H_a + \partial_t) S_a(x, t)) u(x)v(t) dx dt$$
$$+ \lim_{t \to 0^+} \left\{ v(t) \int_{-\infty}^\infty S_a(x, t) u(x) dx \right\} = v(0)u(0).$$

This shows that  $(H_a + \partial_t)G_a(x,t) = \delta(x,t)$ , that is,  $G_a(x,t)$  is the fundamental solution of  $H_a + \partial_t$ . The lemma is thus proved.

Lemma 8.12.21. The fundamental solution  $H_R + \partial_t$  is given by

$$G_R(x,t) = \chi(t)(4\pi t)^{-\frac{n}{2}} \det^{\frac{1}{2}} \left( \frac{tR^M(0)/2}{\sinh(tR^M(0)/2)} \right) \exp\left( -\frac{1}{4t} \left\langle \frac{tR^M(0)/2}{\tanh(tR^M(0)/2)} x, x \right\rangle \right).$$

PROOF. Let  $A \in M_n(\mathbb{R})$  be a skew-symmetric matrix and set  $B = -A^2$ . Then an elaboration of Lemma 8.12.20 shows that the fundamental solution of the heat operator  $-\sum_i \partial_i^2 + \frac{1}{4} \langle Bx, x \rangle + \partial_t$  on  $\mathbb{R}^n \times \mathbb{R}$  is given by

$$(8.43) G_A(x,t) = \chi(t)(4\pi t)^{-\frac{n}{2}} \det^{\frac{1}{2}} \left(\frac{iAt}{\sinh(iAt)}\right) \exp\left(-\frac{1}{4t} \left\langle \frac{iAt}{\tanh(iAt)} x, x \right\rangle\right).$$

To see this we can notice that the O(n)-invariance of  $G_A(x,t)$  enables us to reduce to the case where A is in normal form, i.e., it is a block-diagonal matrix of  $2 \times 2$ -matrices,

$$\left(\begin{array}{cc} 0 & -a_j \\ a_j & 0 \end{array}\right), \qquad a_j > 0,$$

so that the eigenvalues of A are  $\pm ia_j$  and B is a diagonal matrix with the  $a_j^2$  as entries.

A further consequence of the O(n)-invariance of  $G_A(x,t)$  is its invariance under rotations in the  $(x^j, x^k)$ -planes j < k. As the infinitesimal generator of the 1-parameter group of rotations in the  $(x^j, x^k)$ -plane is  $x^j \partial_k - x^k \partial_j$ , it follows that

$$(8.44) (x^j \partial_k - x^k \partial_j) G_A(x, t) = 0.$$

Consider the harmonic oscillator on  $\mathbb{R}^n \times R$  associated to A, i.e.,

$$H_A := -\sum_{j} \left( \partial_j - \sum_{k} \frac{i}{2} A_{jk} x^k \right)^2,$$

and observe that

$$H_A = -\sum_j \partial_j^2 - i \sum_{j,k} A_{jk} x^k \partial_j + \frac{1}{4} \sum_{jkl} A_{jk} A_{jl} x^k x^l$$
$$= -\sum_j \partial_j^2 + i \sum_{j < k} A_{jk} (x^j \partial_k - x^k \partial_j) + \frac{1}{4} \langle Bx, x \rangle.$$

Therefore, in view of (8.44) we that  $G_A(x,t)$  is also the fundamental solution of  $H_A + \partial_t$ . This means that, for all  $u(x,t) \in C_c^{\infty}(\mathbb{R}^n \times \mathbb{R})$ ,

$$u(0,0) = \langle (H_A + \partial_t)G_A(x,t), u(x,t)\rangle \rangle = \langle G_A(x,t), (H_A - \partial_t)u(x,t)\rangle.$$

Observe that  $\langle G_A(x,t), (H_A-\partial_t)u(x,t)\rangle$  is an analytic function of A. Therefore, substituting  $-iR^M(0)/2$  for A in the formula (8.43) yields the fundamental solution of  $H_R + \partial_t$ . The proof is complete.

Combining Lemma 8.12.19 and Lemma 8.12.21 we get

$$\sigma \left[ K_Q(0,0,t) \right]^{(n)} = \left[ G_R(0,1) \wedge e^{-F^W(0)} \right]^{(n)} + \mathcal{O}(t)$$

$$= (4\pi)^{-\frac{n}{2}} \left[ \det^{\frac{1}{2}} \left( \frac{R^M(0)/2}{\sinh(R^M(0)/2)} \right) \wedge e^{-F^W(0)} \right]^{(n)} + \mathcal{O}(t)$$

$$= (4\pi)^{-\frac{n}{2}} \left[ \hat{A}(R^M(0)) \wedge e^{-F^W(0)} \right]^{(n)} + \mathcal{O}(t).$$

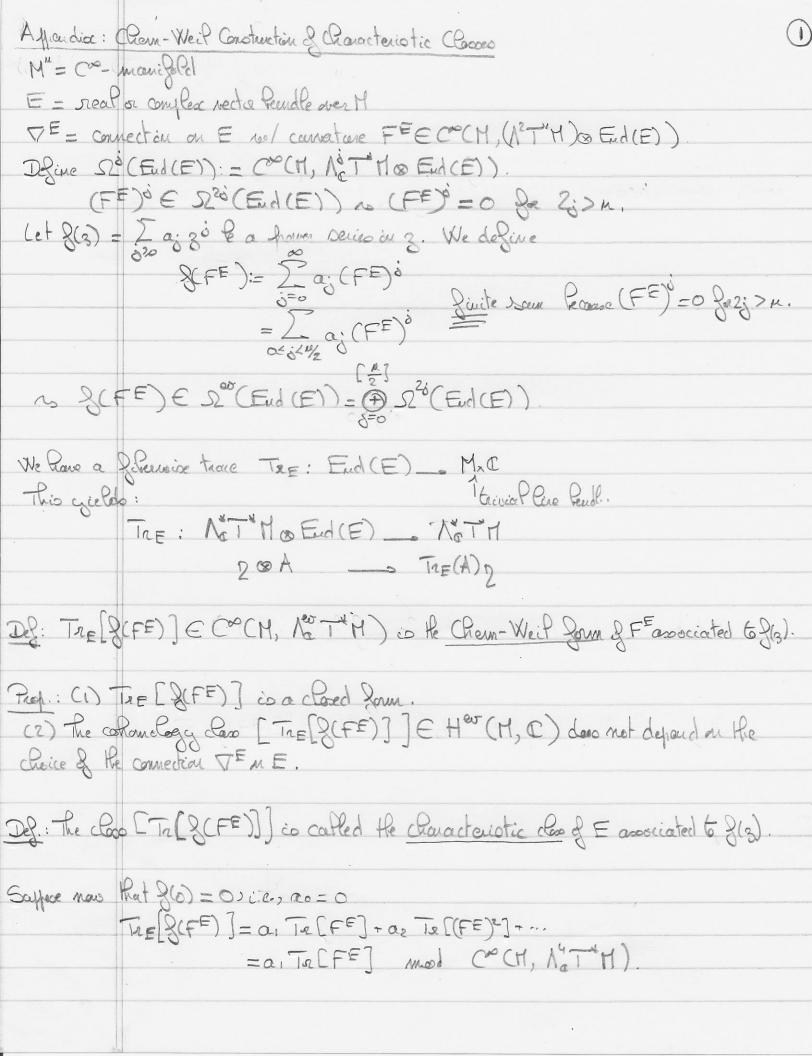
Combining this with (8.28) and observing that  $F^W(0) = F^{E/\$}(0)$ , we deduce that

$$(8.45) \quad \operatorname{Str} k_t(0,0) = (2i\pi)^{-\frac{n}{2}} [\hat{A}(R^M(0)) \wedge \operatorname{Ch}(F^{E/\$}(0))]^{(n)} + \operatorname{O}(t) \quad \text{as } t \to 0^+.$$

This proves (8.23) and completes the proofs of the local index theorem and local index formula of Atiyah-Singer.

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Thus, [TRE S(FE)] = ais[TRE (FE)] o mod COCH, Not2 Total)	2
Therefore, we can define	
exp [Tre g(FF)]:= \( \frac{1}{3!} \) [Tre g(FF)].  =\( \frac{1}{3!} \) [Tre g(FF)] \( 1	
Course the previous proposition it is not difficult to check that each form [TRE &(FE)] is is closed and its commology as	ú
Peap.: The form eap [The &(FF)] is closed and its cohomology class in HeO (M, C) does not depend in the choice of the annection $\nabla^E n E$ .	
Recall that if A is a pointie definite unu motrice, then  det A = eap { Te (Eg A) }.	
et g(z) = I fizz le a paver revies s.t. g(o) = 1. Then leg g(z) = log (1+ (g/z)-1)) is a well defined paver revies. Therefore, for any of ENG, we can define	
det g(z):= copf x TrE[log(g(FE))]]	
Example 1: Chem Character: Cet E le a complex mecta leandle.	
Def.: (1) The (total) Chem form of FE is	
(2) The characteristic class [ CR(FF)] is denoted (R(E)) and called the Chem character of E.	

Eccomple 2: The A-Genes: Let E le a real vector femille over M.

Def.: (1) The Â-form of FF is

 $A(FF) := det^2 \left[ \frac{FF_2}{SinR(FF_2)} \right].$ 

(2) The characteristic class [A(FE)] às called le A-gences of Eaud is donnéted A(E).

(3) If E=TM, Hen A (TM) is ramply denoted A (M) and caffed the A-genus of M.

Prof. (et E, and Ez be real recta lundles aver M. There  $\hat{A}(E, \oplus E_2) = \hat{A}(E_1) \wedge \hat{A}(E_2).$ 

Example 3: The Todd Gencio: Let E le a complex recta leudle aver M

Def.: (1) The Todd form of FE is

Td (FE):= det FE.

(2) The characteristic class [Td(FF)] às called the Todd genus of E and is deneted Td(E).

(3) If M is a complex manifold and E=TM, Hen Td (TM) is simply denoted Td (M) and is called the Todd genes of M.

Prof.: IS E, and Es are complex sector fearables over M, Hen

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